Neutrino-Nucleus Reactions Based on Recent Progress of Shell Model Calculations

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New shell model calculations in p-shell modified shell model Hamiltonian (SFO) with improved spin-dependent transitions and moments Suzuki, Fujimoto, Otsuka Neutrino-nucleus reactions on ¹²C, ⁴He Suzuki, S. Chiba (JAEA), O. Iwamoto (JAEA), Kajino (NAO) New shell model Hamiltonian in fp-shell by Honma et al. Neutrino-nucleus reactions on Fe and Ni isotopes Suzuki, Higashiyama, Otsuka, Kajino, Balantekin

New Shel-Model interaction in p-shell: SFO Suzuki, Fujimoto & Otsuka, PR C67, 044302 (2003)



Spin-tensor decomposition of nuclear effective interactions $V = \sum_{k} V_{k}$ k=0:central k=1:spin-orbit k=2:tensor $< abJT | V_k | cdJT >= (-)^J (2k+1) \sum < ab | LSJ > < cd | L'S'J >$ LĽSS $\left\{ \frac{LSJ}{S'L'k} \left\{ \sum_{J'} (-)^{J'} (2J'+1) \right\} \frac{LSJ}{S'L'k} \right\}_{i'} \sum_{j'=i'=i'} (-)^{J'} (2J'+1) \left\{ \frac{LSJ}{S'L'k} \right\}_{i'=i'=i'} (-)^{J'} (2J'+1) \left\{ \frac{LSJ}{S'L'k} \right\}_{i'=i'=i'} (-)^{J'} (2J'+1) \left\{ \frac{LSJ}{S'L'k} \right\}_{i'=i'=i'=i'} (-)^{J'} (-)^{J'} (2J'+1) \left\{ \frac{LSJ}{S'L'k} \right\}_{i'=i'=i'=i'} (-)^{J'} (-)^{J$ *<a'b'J'T\Vc'd'J'T>* where $a = \{n_a l_a j_a\}, a' = \{n_a l_a j_a'\}$ etc.



Effects of Tensor Force on Shell Evolution



FIG. 1: (a) Schematic picture of the monopole interaction produced by the tensor force between a proton in $j_{>,<} = l \pm 1/2$ and a neutron in $j'_{>,<} = l' \pm 1/2$. (b) Exchange processes contributing to the monopole interaction of the tensor force.



FIG. 2: Intuitive picture of the tensor force acting two nucleons on orbits j and j'.





Energy levels of B and C isotopes



B(GT) values for ${}^{12}C$ -> ${}^{12}N$

Magnetic moments of p-shell nuclei



Magnetic Moments s.p.: $[p_{1/2} \times p_{3/2}]_{1}^{1}$ $(-g^{?} + 1/2g^{s}) + 1/2g^{s} = -0.12\mu_{N}$

 $CK: \mu = 0.778 \mu N$ Mixing of p1/2 and p3/2





Spin-dependent excitations

Gamow-Teller (1⁺):
$$\vec{\mathbf{s}} \cdot \vec{\mathbf{t}}_{\pm}$$

Spin-diole (0⁻, 1⁻, 2⁻): $[\vec{\mathbf{s}} \times \vec{r}]^J \cdot \vec{\mathbf{t}}_{\pm}$

Multipoles

1⁺: E_51 , M1, C_51 , L_51 2⁻: E_52 , M2, C_52 , L_52 1⁻: M_51 , E1, C1 0⁻: C_50 , L_50

Folding over neutrino spectrum $\mathbf{s}_{E_x} = \int \mathbf{s}_{E_x} (E) f(E) dE$ $\mathbf{s}_{E_x} (E) = \int \frac{d\mathbf{s}_{E_x}}{d\Omega} d\Omega$













Spin-dipole strength in ⁴He







⁴He($\overline{v}_{e} e^{+}n$)³H ⁴He($v_{e} e^{-}n$)³He

Neutrino Nucleus Reactions on Fe and Ni Isotopes

Charge-exchange reactions; $(\mathbf{n}_{e}, e^{-}), (\overline{\mathbf{n}}_{e}, e^{+})$ Gamow-Teller transitions GT strength: shell model calculation by Honma

Emissions of proton, neutron, ,

 $(\mathbf{n}_e, e^- p), (\mathbf{n}_e, e^- n), (\mathbf{n}_e, e^- a), (\mathbf{n}_e, e^- g), (\mathbf{n}_e, e^- pn),$

 $(\mathbf{n}_e, e^- pp), (\mathbf{n}_e, e^- ap), (\bar{\mathbf{n}}_e, e^+ n)$

New features in GXPF1 by Honma et al.

KB3G A = 47-52KB + monopole corrections G: improved gap, 3: fine tuning in V_{fr} (A = 50 - 52) GXPF1 A = 47-66More attraction for T=0 m.e. than G-matrix $E(1p_{3/2}) - E(0f_{7/2}) \sim 3 \text{ MeV}$ cf. ~ 2 MeV for KB3, FPD6 New magic number at N = 34 cf. N = 32 for KB3, FPD6 r.m.s deviations from observed values GXPF1 vs KB3G Z, N <28; Z<28, N>28: similar Z or N =28 (56 Ni, 55 Co, 57 Ni); Z, N>28: smaller for GXPF1 Core excitations are not well described by KB3G Monopoles GXPF1 vs KB3G 1p-1p part differ GXPF1: not a constant shift (J-dependence) Systematic description of 2_1^+ energies in Ni, Ca, Ti, Cr, Fe isotopes cf. KB3G ⁵⁶Ni : $E_{cal} - E_{exp} \sim 2 \text{ MeV}$ ⁵⁶Ni-core $(f_{7/2})^{16}$: 69% (GXPF1), 49% (FPD6)

 $B(GT_+)$: 11.3 (GXPF1), 10.1 (KB3), 13.7 (closed core)

Cross Sections for GT (J=1⁺) transitions

$$(\frac{d \, \mathbf{s}}{d \, \Omega})_{(\mathbf{n}, e^{-})/(\bar{\mathbf{n}}, e^{+})} = \frac{G_{F}^{2} \cos^{2} \, \mathbf{q}_{C}}{2 \, \mathbf{p}^{2}} E_{e} \, p_{e} F \left(Z_{f}, E_{f}\right) \frac{4 \, \mathbf{p}}{2 \, J_{i} + 1} \cos^{2} \frac{\mathbf{q}}{2}$$

$$\{ K_{T} \left(q, \mathbf{w}\right) [| < J_{f} || T_{J}^{mag} || J_{i} > |^{2} + | < J_{f} || T_{J}^{el,5} || J_{i} > |^{2}] \mp K_{TI} \left(q, \mathbf{w}\right)$$

$$2 \operatorname{R} e[< J_{f} || T_{J}^{mag} || J_{i} > < J_{f} || T_{J}^{el,5} || J_{i} >^{*}] \}$$

$$< J_{f} || T_{J}^{el,5} || J_{i} > \cong ig_{A} \sqrt{\frac{2}{3}} \frac{1}{\sqrt{4p}} < J_{f} || \mathbf{s} j_{0}(qr) || J_{i} >$$

$$< J_{f} || T_{J}^{mag} || J_{i} > \approx -\frac{q}{2M} g^{IV} \sqrt{\frac{2}{3}} \frac{1}{\sqrt{4p}} < J_{f} || \mathbf{s} j_{0}(qr) || J_{i} >$$

$$< J_{f} || \mathbf{s} |j_{0}(qr) || J_{i} >; < J_{f} || \mathbf{s} || J_{i} > j_{0}(qR)$$

$$cf. \quad B(GT) = \frac{1}{2J_{i}+1} |< J_{f} || \mathbf{s} || J_{i} > |^{2}$$

$$K_{T}(q, \mathbf{w}) = \frac{q^{2} - \mathbf{w}^{2}}{2q^{2}} + \tan^{2}\frac{q}{2}, \quad K_{TI}(q, \mathbf{w}) = \tan\frac{q}{2}\sqrt{\frac{q^{2} - \mathbf{w}^{2}}{2q^{2}}} + \tan^{2}\frac{q}{2}$$













B(GT) for 58Ni















Summary

- 1. P-shell SFO $B(GT) \uparrow \Rightarrow s(n,e) \uparrow$ $(n,n'n) \uparrow$
- 2. fp-shell GXPF1
 B(GT) ↑, more fragmented, s(n,e) ↑
 e-capture and -decay rates steller core collapse and supernovae explosion mechanism

Remaning issues:

Inclusion of spin-dipole contributions Study of effects of n from (n, en)(n,) in r-process in stars