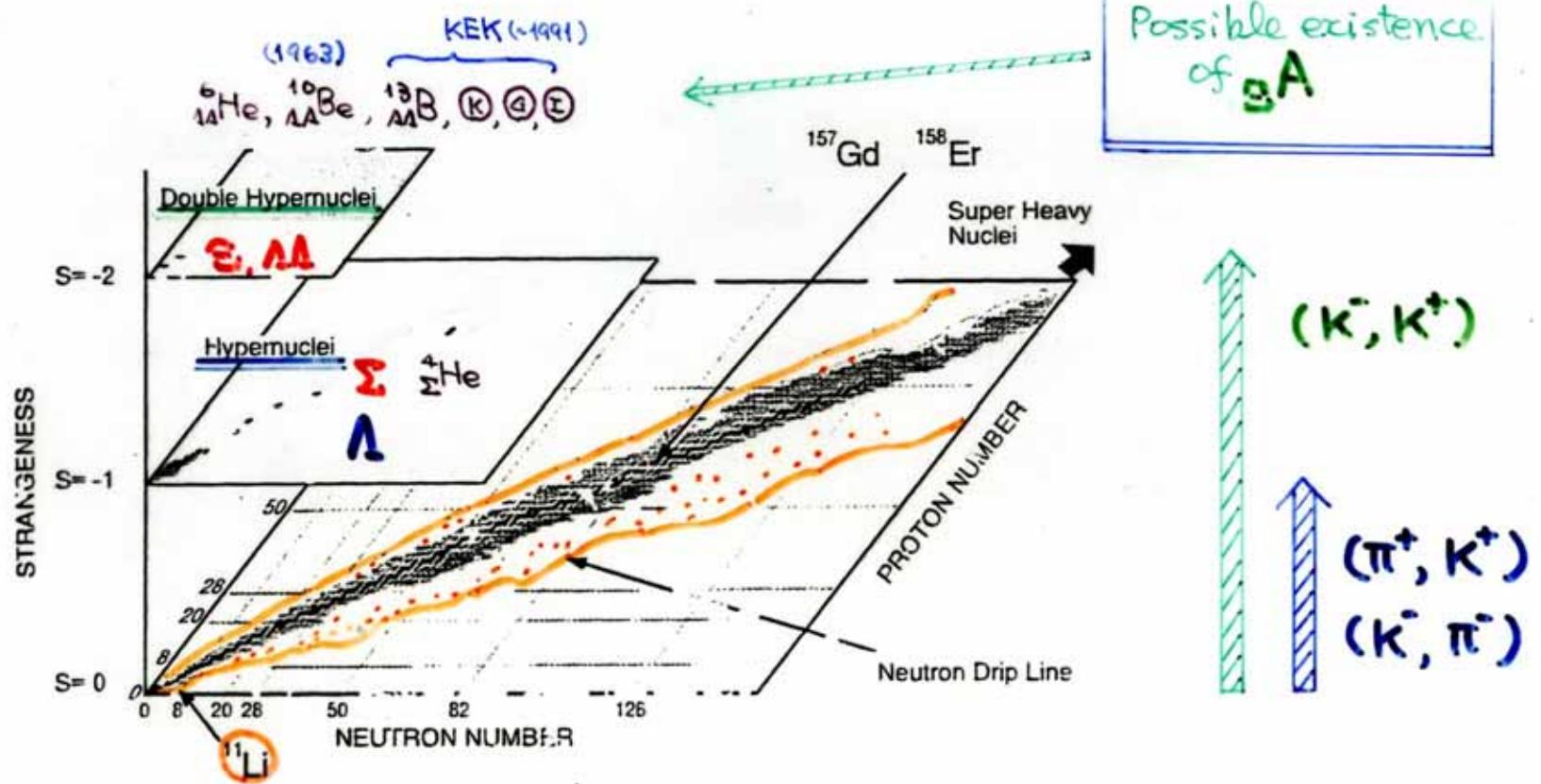


CISS03
Sept.16-20, 2003

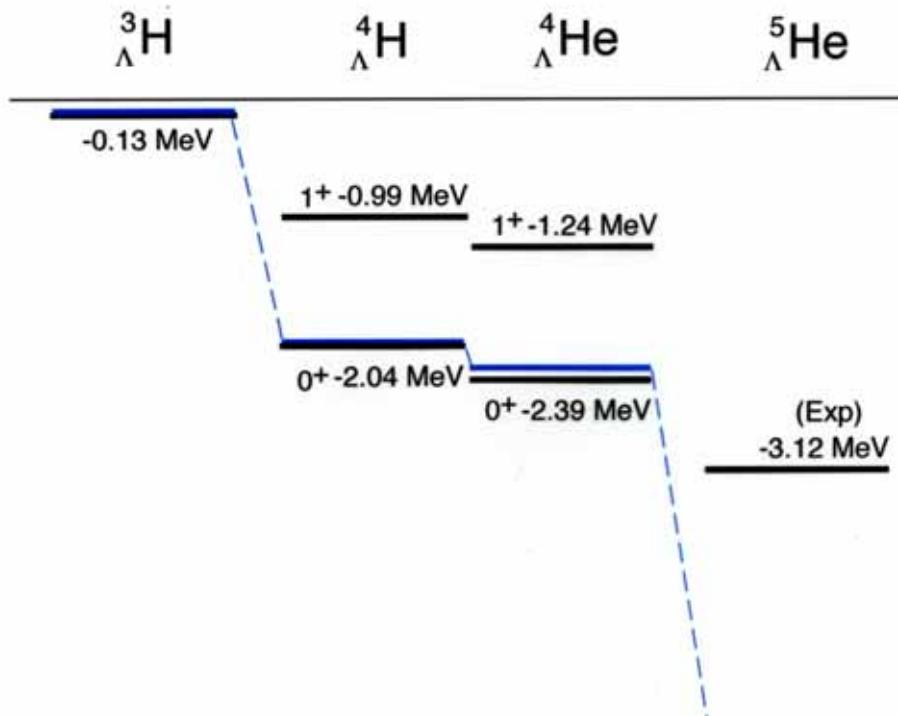
Coherent Λ - Σ Coupling in Asymmetric Hypernuclei

Yoshinori AKAISHI

Institute of Particle and Nuclear Studies, KEK



The Overbinding Problem



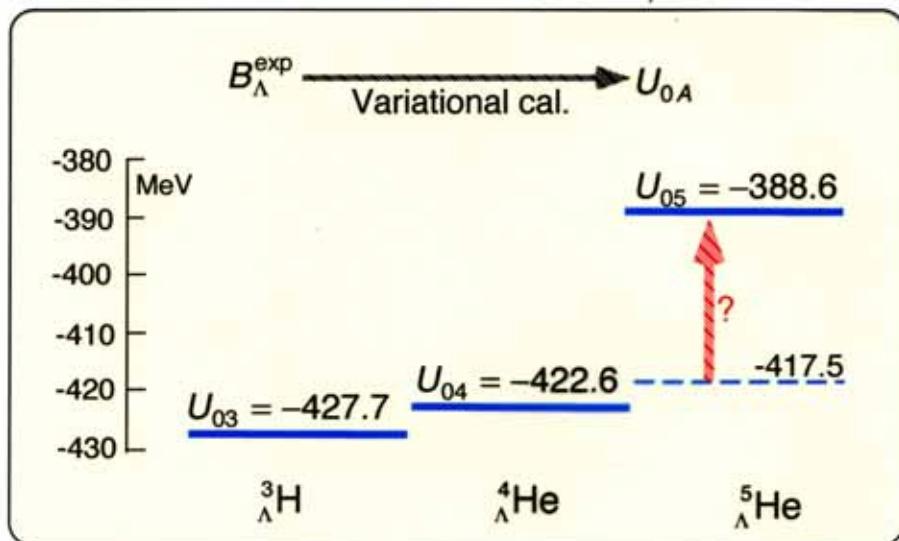
R.H. Dalitz, R.C. Herndon and Y.C. Tang,
Nucl. Phys. **B47** (1972) 109

(Exp)
-3.12 MeV
Overbound

Dalitz et al's analysis

$$V_{\Lambda N}(r) = \begin{cases} \infty, & r < d \\ U_{0A} \exp[-\lambda(r-d)], & r > d \end{cases}$$

$\lambda = 3.219 \text{ fm}^{-1}, d = 0.45 \text{ fm}$



$$\begin{aligned} U_{03} &= \frac{1}{4}U_0^t + \frac{3}{4}U_0^s \\ U_{04} &= \frac{1}{2}U_0^t + \frac{1}{2}U_0^s \\ U_{05} &= \frac{3}{4}U_0^t + \frac{1}{4}U_0^s \end{aligned}$$

$\frac{1}{2}(U_{03} + U_{05}) - U_{04} = 0 \quad .03W_3$

$W_3 \approx 480 \text{ MeV}$

$W_3 = 1.43 \text{ MeV}$

Nogami's 3BF

$$V_{\Lambda NN} = -\frac{1}{3}W_3(\vec{\sigma}_1 \vec{\sigma}_2)(\vec{\tau}_1 \vec{\tau}_2) \frac{\exp(-\mu r_{1\Lambda})}{\mu r_{1\Lambda}} \frac{\exp(-\mu r_{2\Lambda})}{\mu r_{2\Lambda}}$$

Singlet int. is more attractive than triplet int.

Central YN interaction

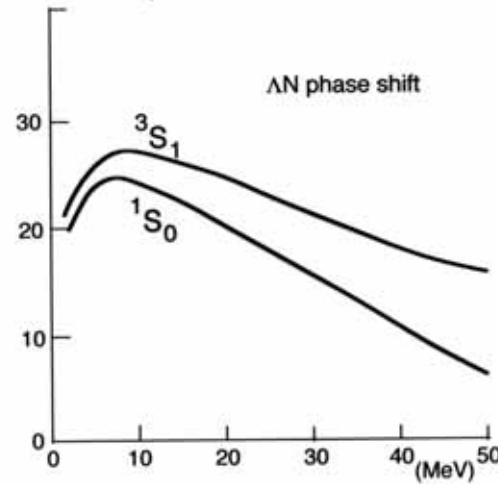
	ΛN	ΣN
D0	○	
D2	○	○

1-channel picture

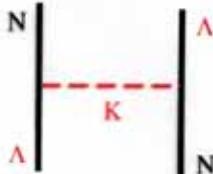
2-channel picture

A.R. Bodmer
(1966)

Phase-equivalent to Nijmegen D



Brueckner-Hartree-Fock Calculation on Gaussian Basis



Hyperon-nucleus potential

$$U_\mu^{\text{eq}}(\vec{r}_1)\varphi_\mu(\vec{r}_1) \equiv U_H(\vec{r}_1)\varphi_\mu(\vec{r}_1) + \int d\vec{r}_2 U_F(\vec{r}_1, \vec{r}_2)\varphi_\mu(\vec{r}_2)$$

g-matrix

$$U_H(\vec{r}_1) = \int d\vec{r}_2 \sum_v \varphi_v^*(\vec{r}_2) g(\vec{r}_1, \vec{r}_2) \varphi_v(\vec{r}_2)$$

$$U_F(\vec{r}_1, \vec{r}_2) = - \sum_v \varphi_v^*(\vec{r}_2) g(\vec{r}_1, \vec{r}_2) \varphi_v(\vec{r}_1)$$

$$g(\vec{r}_1 - \vec{r}_2) = \sum_j \gamma_j^\mu \exp\left\{-\left((\vec{r}_1 - \vec{r}_2)/c_j\right)^2\right\}, \quad c_1, \dots, c_{20} \text{ fixed}$$

$$U_\mu^{\text{eq}}(\vec{r}_1) = \sum_j \alpha_j^\mu \exp\left\{-\left(r_1/a_j\right)^2\right\}, \quad a_1, \dots, a_{10} \text{ fixed}$$

$$\varphi_\mu(\vec{r}_1) = \sum_j \beta_j^\mu \exp\left\{-\left(r_1/b_j\right)^2\right\} r_1^\ell Y_{\ell, s_1 j_1 m_1}(\hat{r}_1), \quad b_1, \dots, b_{20} \text{ f}$$

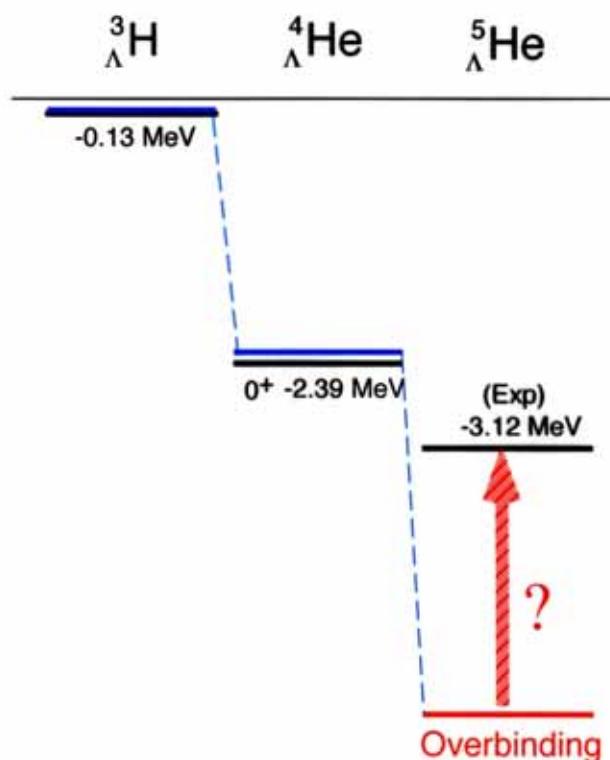
Self-consistently determined

Matrix elements for s-shell hyperon:

$$\begin{aligned}
 & \sum_{j_1, j_2} \langle (a_1'' j_1 \mu_1 \phi_1) (a_2'' j_2 \mu_2 \phi_2) | g | (a_1 j_1 \mu_1 \phi_1) (a_2 j_2 \mu_2 \phi_2) - \text{exch.} \rangle \\
 &= (2j_2+1)(2\ell_2+1)^2 \sum_{\ell'=0}^{\ell_2} \sum_{n''=0}^{\ell_2} \left(\frac{M_1}{M_1+M_2} \right)^{\ell'+\ell''} \sqrt{\frac{2\ell_2}{2\ell'} \frac{2\ell_2}{2\ell''}} \sum_n \sum_{n''} (2n+1)(2n''+1) \\
 & \times \sum_{\kappa=0}^{\infty} \int dr r^{\ell'+\ell''+2} \exp(-(A_{12}'' + A_{12} + c_\kappa^2) r^2) \int_0^R R^{2\ell_2 - \ell' - \ell'' + 2} \exp(-(a_1'' + a_2'' + a_1 + a_2) R^2) \\
 & \quad \times i^{n''} j_{n''} (ia_{12}'' r R) i^n j_n (ia_{12} r R) \\
 & \times \sum_{\ell} (\ell'' n 0 0 \bar{f} 0) (\ell''' n'' 0 0 \bar{f} 0) \sum_{\tilde{\ell}} (\ell_2 - \ell' n 0 0 \bar{f} 0) (\ell_2 - \ell''' n'' 0 0 \bar{f} 0) \left\{ \begin{matrix} n & \ell' & \ell \\ \ell_2 & \tilde{\ell} & \ell_2 - \ell \end{matrix} \right\} \left\{ \begin{matrix} n'' \ell''' & \ell \\ \ell_2 & \tilde{\ell} & \ell_2 - \ell''' \end{matrix} \right\} \\
 & \times \left[\sum_{K=\ell_2} \frac{2K+1}{2} \left\{ \begin{matrix} j_2 & K & \frac{1}{2} \\ 0 & \frac{1}{2} & \ell_2 \end{matrix} \right\} \right]^2 \sum_{J=\tilde{\ell}} (2J+1) \left\{ \begin{matrix} \tilde{\ell} & 0 & J \\ K & \tilde{\ell} & \ell_2 \end{matrix} \right\}^2 \frac{1 - (-)^{n-\ell}}{2} [V_{\tilde{\ell}, S=0}^{J, F=0}(\kappa)]_{Yn \text{ and } Yp} \\
 & + \dots [V_{\tilde{\ell}, S=0}^{J, F=1}(\kappa)]_{Yn \text{ and } Yp} + \dots [V_{\tilde{\ell}, S=1}^{J, F=0}(\kappa)]_{Yn \text{ and } Yp} + \dots [V_{\tilde{\ell}, S=1}^{J, F=1}(\kappa)]_{Yn \text{ and } Yp}
 \end{aligned}$$

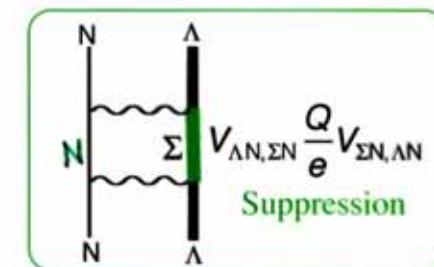
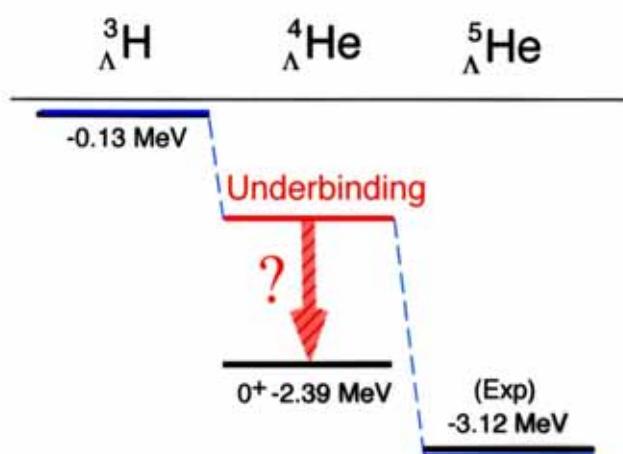
D0

The Overbinding Problem



D2

The Underbinding Problem



Λ - Σ Coupling

B.F. Gibson, A.Goldberg and M.S. Weiss, Phys. Rev. C6 (1972) 741;
 J. Dabrowski, Phys. Rev. C8 (1973) 835.

$$|{}^5_{\Lambda}\text{He}\rangle = \Phi_{\Lambda}(\vec{r}) |{}^4\text{He}\rangle \quad T=0$$

$$|{}^4_{\Lambda}\text{He}\rangle = \Phi_{\Lambda}(\vec{r}) |{}^3\text{He}\rangle + \sqrt{\frac{2}{3}}\Phi_{\Sigma^+}(\vec{r}) |{}^3\Sigma\rangle - \sqrt{\frac{1}{3}}\Phi_{\Sigma^0}(\vec{r}) |{}^3\text{He}\rangle$$

$T=1/2$

${}^4_{\Sigma}\text{He}(0^+)$

$$\frac{9}{6} {}^3g_{\Sigma N} + \frac{1}{6} {}^1g_{\Sigma N} + \frac{8}{6} {}^1g_{\Sigma N}^{T=3/2}$$

$$\frac{7}{6} {}^3g_{\Sigma N} + \frac{3}{6} {}^1g_{\Sigma N} + \frac{8}{6} {}^3g_{\Sigma N}^{T=3/2}$$

$$\frac{3}{2} {}^3g_{\Lambda N, \Sigma N} - \frac{1}{2} {}^1g_{\Lambda N, \Sigma N}$$

$$\frac{3}{2} {}^3g_{\Lambda N} + \frac{3}{2} {}^1g_{\Lambda N}$$

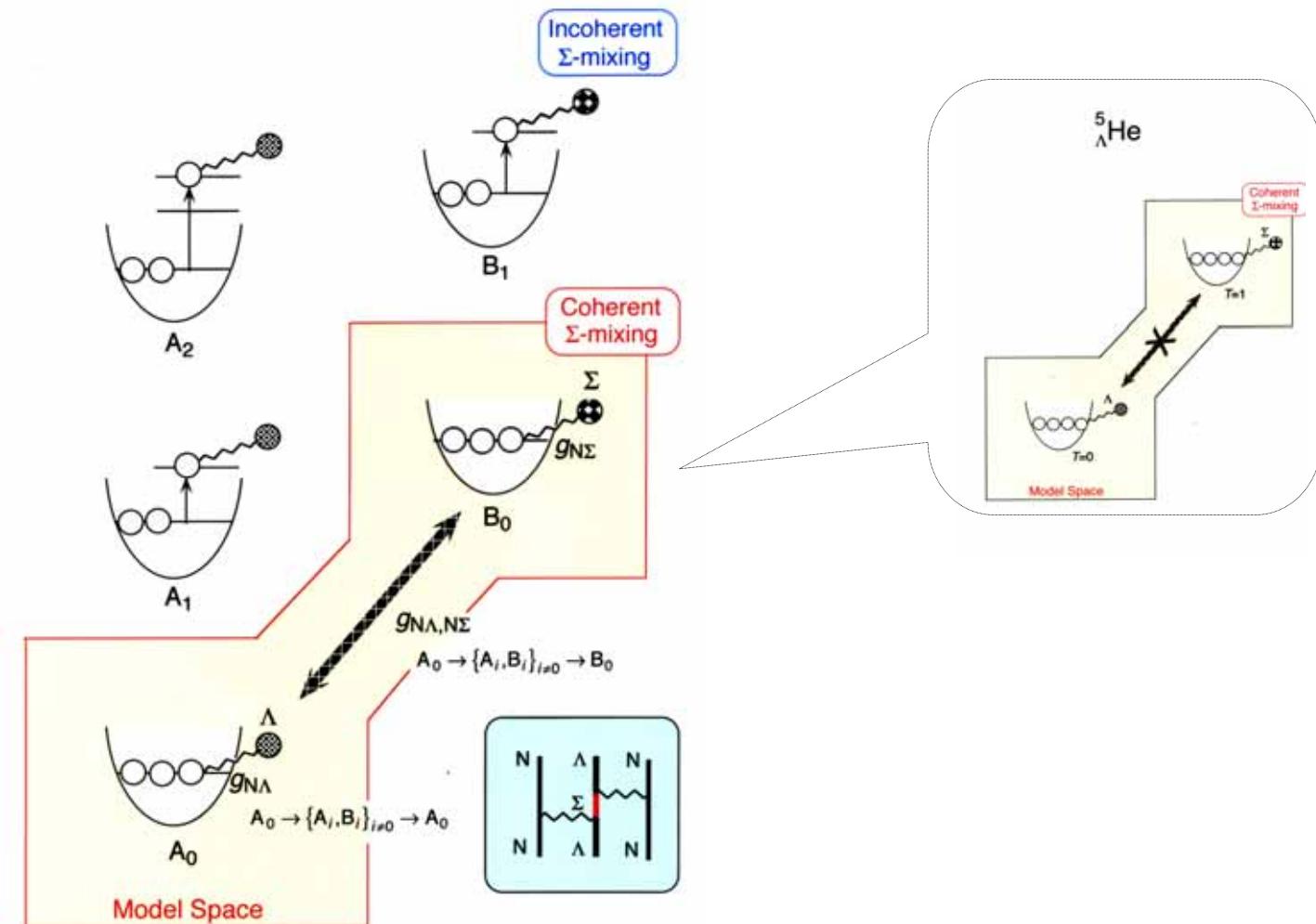
${}^4_{\Lambda}\text{He}(0^+)$

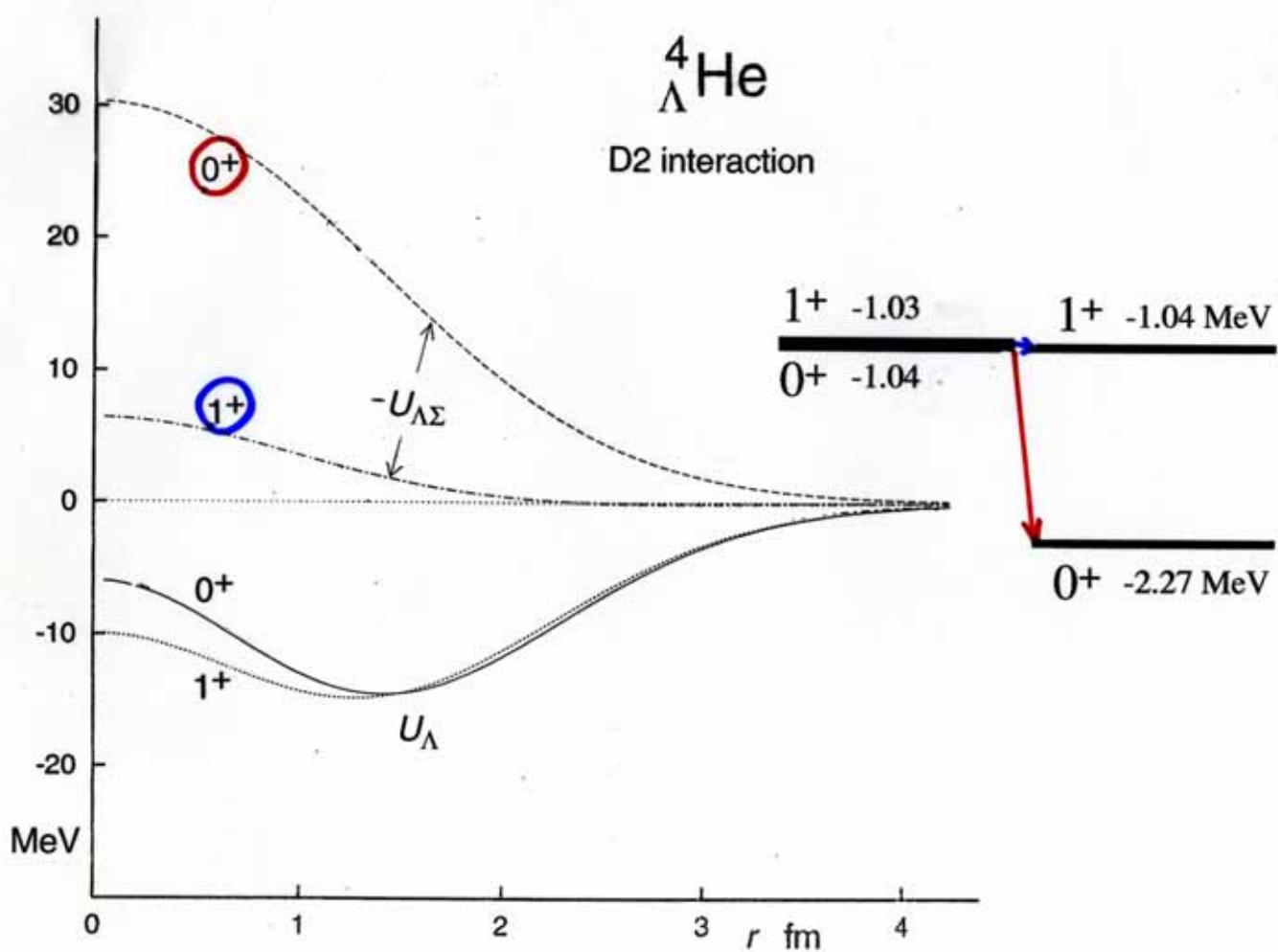
$$\frac{1}{2} {}^3g_{\Lambda N, \Sigma N} + \frac{1}{2} {}^1g_{\Lambda N, \Sigma N}$$

$$\frac{5}{2} {}^3g_{\Lambda N} + \frac{1}{2} {}^1g_{\Lambda N}$$

${}^4_{\Lambda}\text{He}(1^+)$

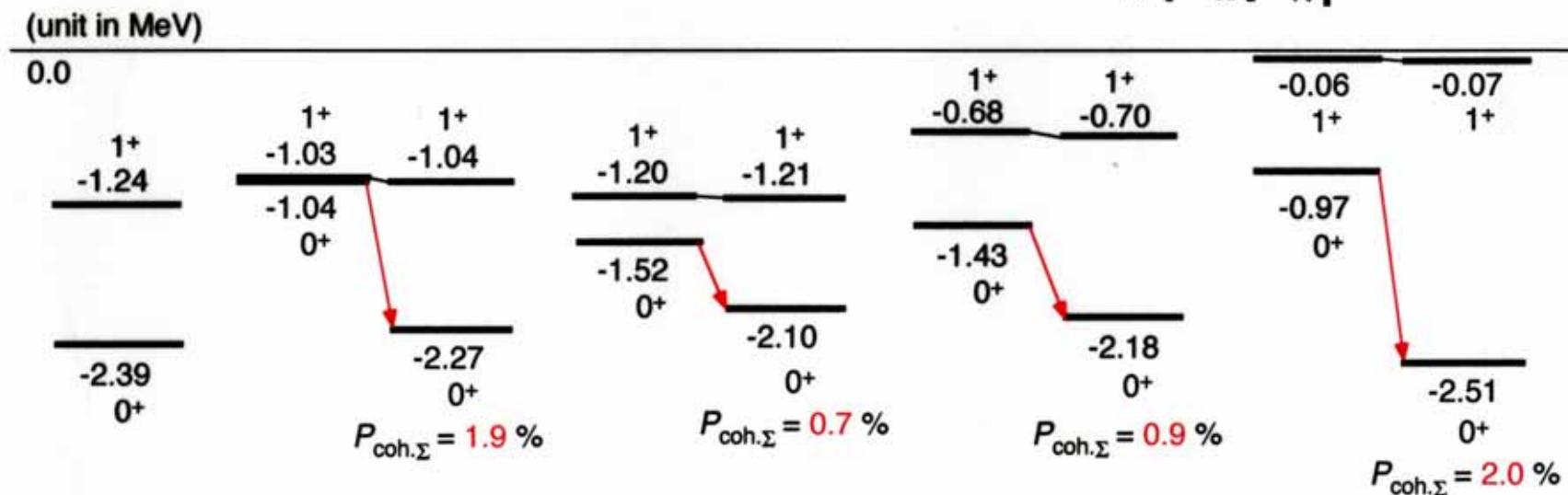
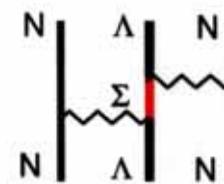
Coherent Λ - Σ Coupling





"The 0^+-1^+ difference is not
a measure of ΛN spin-spin interaction."

B. Gibson



Exp

D2

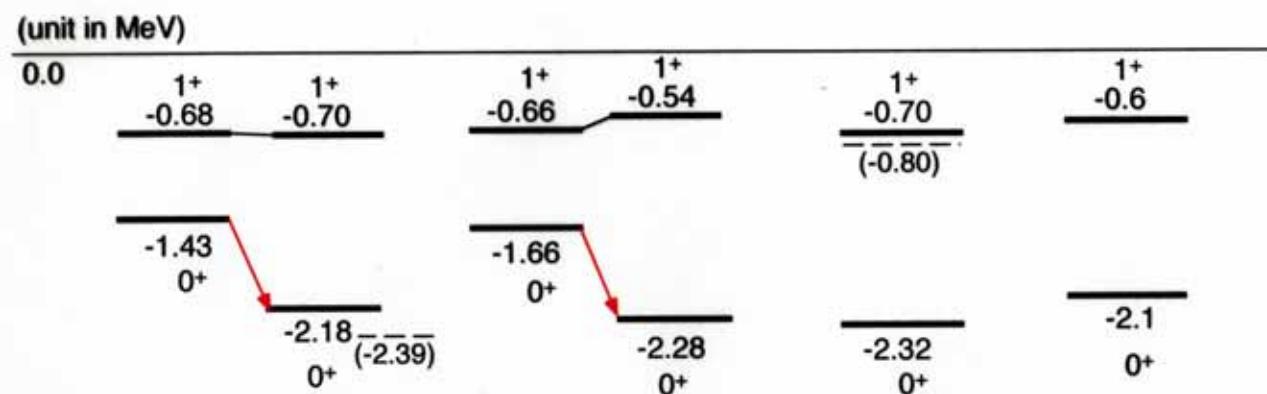
SC97e(S)

SC97f(S)

SC89(S)

T. Rijken et al., Phys. Rev. C59 (1999) 21.

$^4_{\Lambda}\text{He}$ SC97f(S)



Y. Akaishi et al.
BHF (2000)

E. Hiyama et al.
GVM (2001)
AV8

A. Nogga
F-Y (2001)
AV8

H. Nemura
SVM (2001)
G3RS

E. Hiyama, M. Kamimura, T. Motoba, T. Yamada & Y. Yamamoto, Phys. Rev. C65 (2001) 011301(R).
A. Nogga, Doctoral dissertation.

Faddeev-Yakubovsky calculations for ${}^4\Lambda$ He

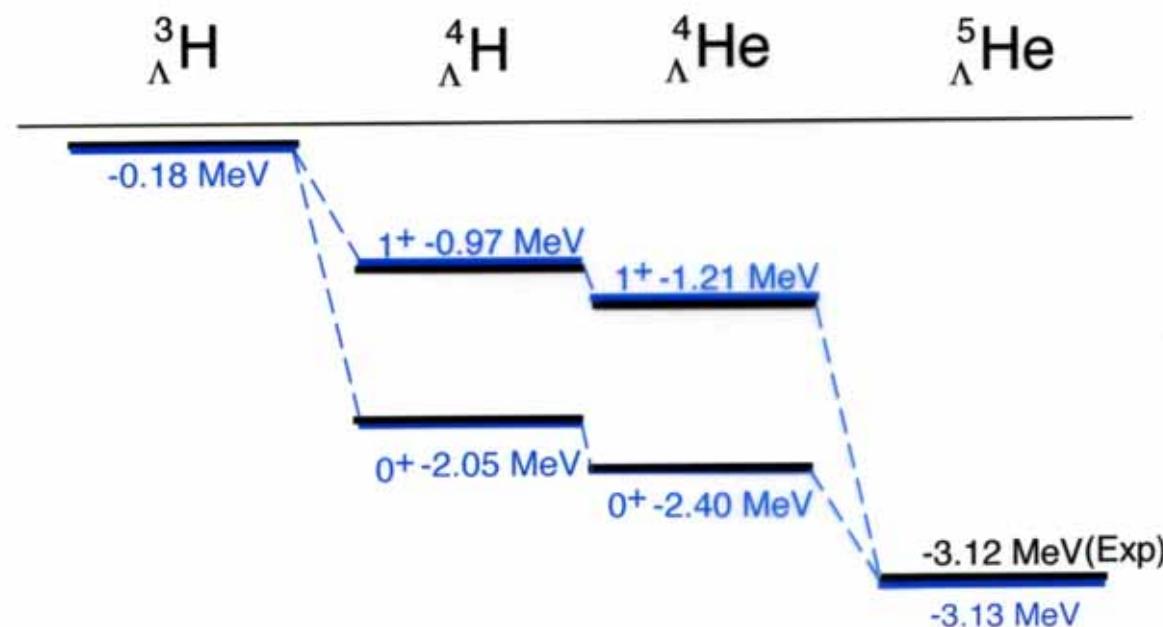
A. Nogga, H. Kamada and W. Glöckle, Phys. Rev. Lett. **88** (2002) 172501

<i>YNF</i>	$E_{\text{sep}}^{\Lambda}(0^+)$	$E_{\text{sep}}^{\Lambda}(1^+)$	Δ
SC89	2.14	0.02	2.06
SC97 <i>f</i>	1.72	0.53	1.16
SC97 <i>e</i>	1.54	0.72	0.79
SC97 <i>d</i>	1.29	0.80	0.47
Expt.	2.39(3)	1.24(5)	1.15

		$^4_{\Lambda}H$	
S=1 pairs		1+	0+
$\Lambda p \Leftrightarrow -\sqrt{\frac{1}{3}} \Sigma^0 p + \sqrt{\frac{2}{3}} \Sigma^+ n$	$\begin{cases} S_3 = 1 \\ S_3 = 0 \\ S_3 = -1 \end{cases}$	-1/3 +1/3 +1/2	+1/2 +1/2 +1/2
$\Lambda n \Leftrightarrow \sqrt{\frac{1}{3}} \Sigma^0 n - \sqrt{\frac{2}{3}} \Sigma^- p$	$\begin{cases} S_3 = 1 \\ S_3 = 0 \\ S_3 = -1 \end{cases}$		
Contribution to $U_{\Sigma\Lambda}$		1/2	3/2
Λ - Σ coupling effect ($\sim \Lambda NN$ force)		1 : 9	

The Overbinding Problem

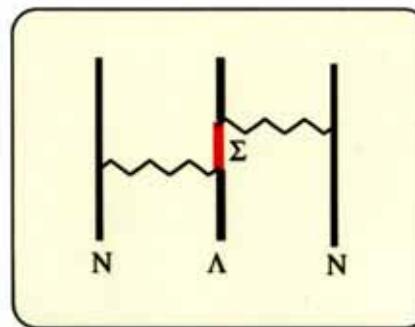
has been solved.



H. Nemura: D2'+Minnesota NN

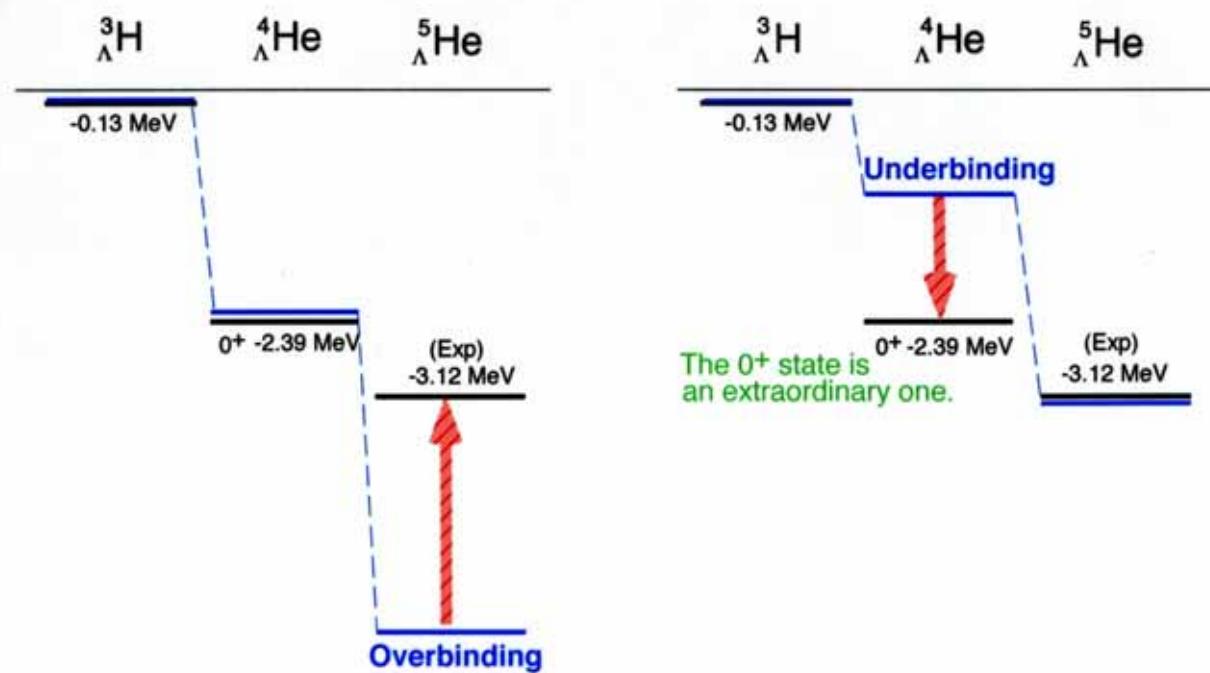
Repulsion !

Y. Nogami et al.
Nucl. Phys. B19 (1970) 93

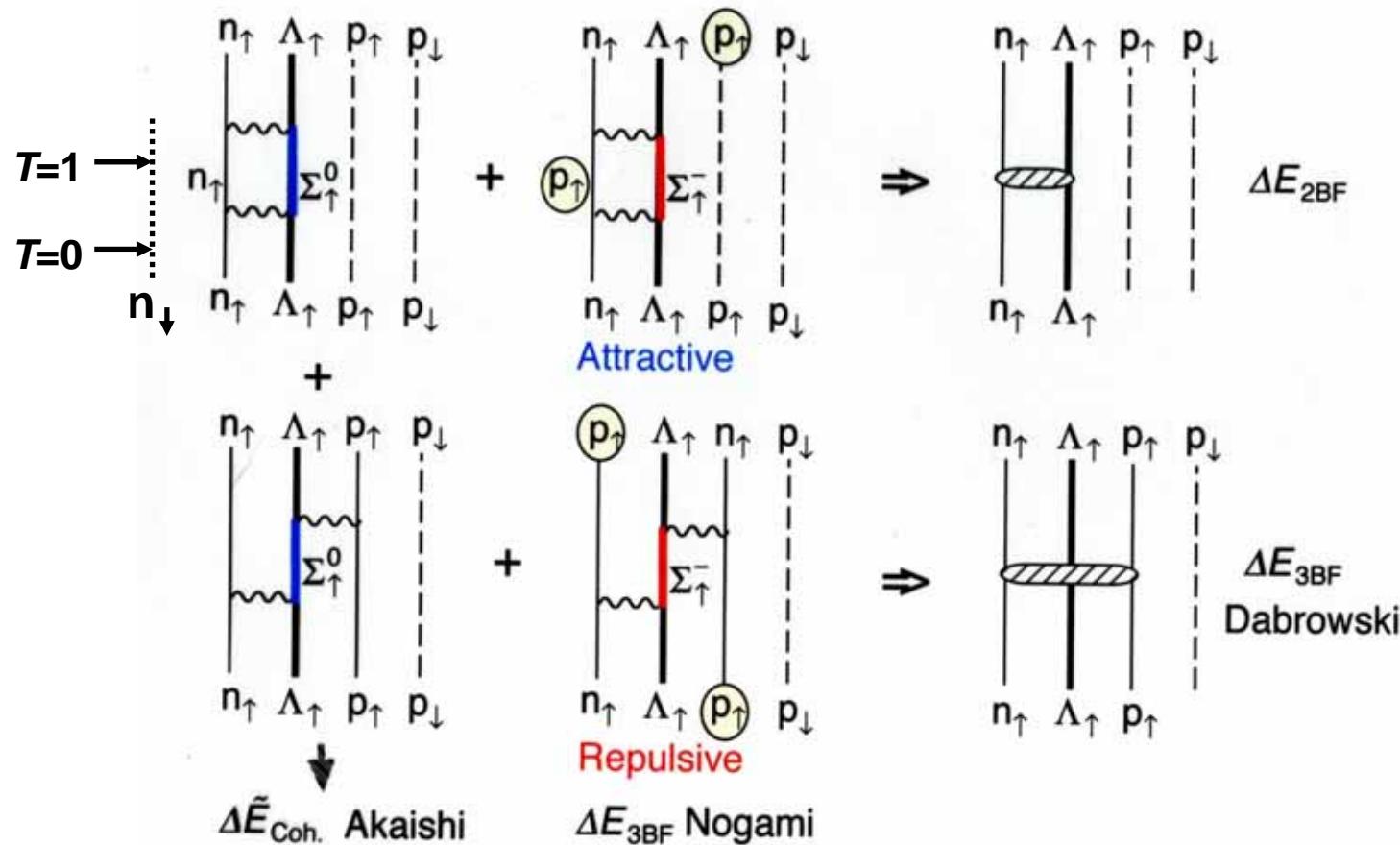


Attraction !

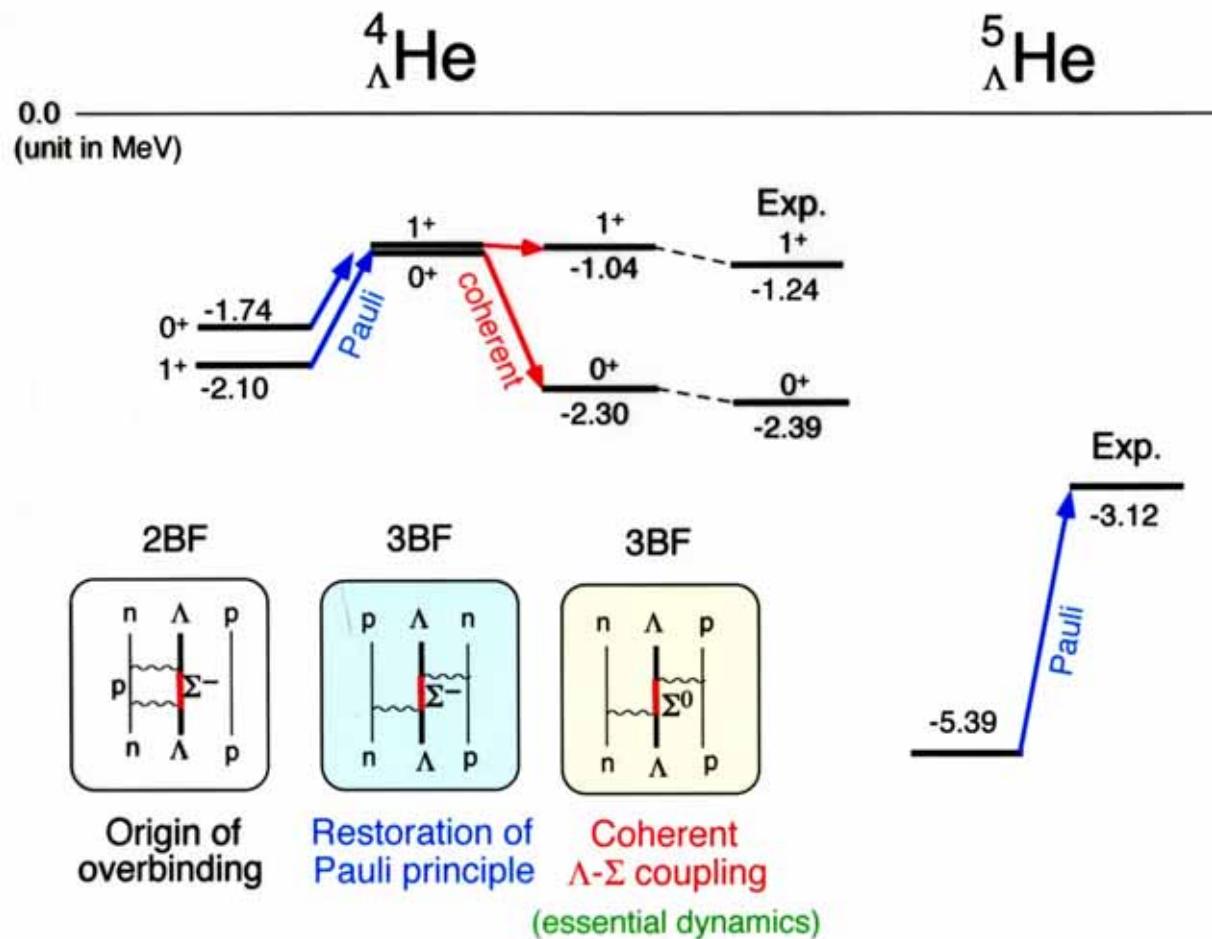
Y. Akaishi et al.
Phys. Rev. Lett. 84 (2000) 3539



$^4_{\Lambda}\text{He}$ NNN in $(0s)^3$



Effects of Λ NN Three-Body Force

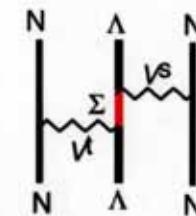


Three-Body Force due to Coherent Λ - Σ Coupling : [for D0]

$$U_{\Lambda NN} = \sum_{\alpha=tt,ts,ss} W_3^\alpha(r_{1\Lambda}, r_{\Lambda 2}) \left[a_\alpha + b_\alpha (\vec{\sigma}_1 \vec{\sigma}_2) + c_\alpha \frac{1}{2} \vec{\sigma}_\Lambda (\vec{\sigma}_1 + \vec{\sigma}_2) \right]$$

ΛN spin-spin

$$\begin{pmatrix} a_{tt} & b_{tt} & c_{tt} \\ a_{ts} & b_{ts} & c_{ts} \\ a_{ss} & b_{ss} & c_{ss} \end{pmatrix} = \begin{pmatrix} \frac{7}{16} & \frac{3}{16} & \frac{3}{8} \\ \frac{1}{8} & \frac{1}{8} & -\frac{1}{4} \\ \frac{5}{48} & \frac{1}{48} & -\frac{1}{8} \end{pmatrix}$$



$$W_3^{ts}(r_{1\Lambda}, r_{\Lambda 2}) = V_{\Lambda N, \Sigma N}^t(r_{1\Lambda}) \frac{1}{\Delta M^*} V_{\Sigma N, \Lambda N}^s(r_{\Lambda 2}) + V_{\Lambda N, \Sigma N}^s(r_{1\Lambda}) \frac{1}{\Delta M^*} V_{\Sigma N, \Lambda N}^t(r_{\Lambda 2})$$

$$\langle V_{\Sigma N, \Lambda N}^s \rangle = -\beta \langle V_{\Sigma N, \Lambda N}^t \rangle \quad \beta = 0.67, \quad \langle W_3^t \rangle_4 = 0.74 \text{ MeV}$$

$${}_{\Lambda}^5 \text{He} \quad \frac{1}{2} (3 + \beta^2) \langle W_3^t \rangle_5 \quad \approx 3 \text{ MeV}$$

$${}_{\Lambda}^4 \text{H}^*(1^+) \quad \frac{1}{8} (9 + 2\beta + \beta^2) \langle W_3^t \rangle_4 \quad \approx 1.0 \text{ MeV}$$

$${}_{\Lambda}^4 \text{H}(0^+) \quad \frac{1}{8} (-3 - 6\beta + 5\beta^2) \langle W_3^t \rangle_4 \quad \approx -0.44 \text{ MeV}$$

$${}_{\Lambda}^3 \text{H} \quad \frac{1}{8} (-1 - 6\beta + 3\beta^2) \langle W_3^t \rangle_3 \quad \approx -0.05 \text{ MeV}$$

Light Hypernuclei

A.R. Bodmer and Q.N. Usmani, Nucl. Phys. **A477** (1988) 621.

R. Sinha and Q.N. Usmani, Nucl. Phys. **A684** (2001) 586c.

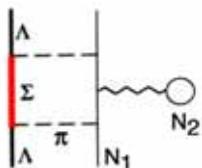
$$V_{\Lambda N} = \left[(V_{\text{core}}(r) - \bar{V})(1 - \varepsilon + \varepsilon P_x) + \frac{1}{4} V_\sigma \bar{\sigma}_\Lambda \bar{\sigma}_N \right] T_\pi^2(r)$$

spin-spin

$$V_{\Lambda NN} = W_p \{\text{Nogami}\} + V_{\Lambda NN}^{\text{DS}} \{\text{Dispersive}\}$$

$$V_{\Lambda NN}^{\text{DS}}(r_{ij\Lambda}) = W_0 T_\pi^2(r_{i\Lambda}) T_\pi^2(r_{j\Lambda}) \left[1 + \frac{1}{6} \bar{\sigma}_\Lambda (\bar{\sigma}_i + \bar{\sigma}_j) \right]$$

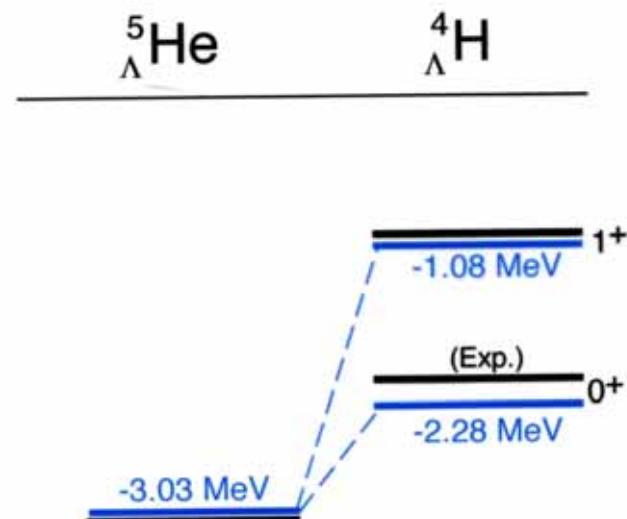
spin-spin



The 0^+ - 1^+ splitting
 0.38 + 0.86 MeV
 (70%)

The major part is attributed
 to ΛNN and not to ΛN .

0.56 + 0.92 MeV for SC97f
 (62%)
 -0.17 + 1.40 MeV for D2
 (114%)

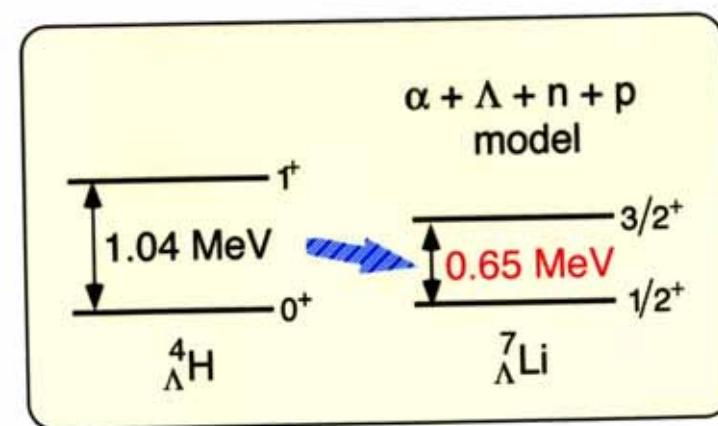
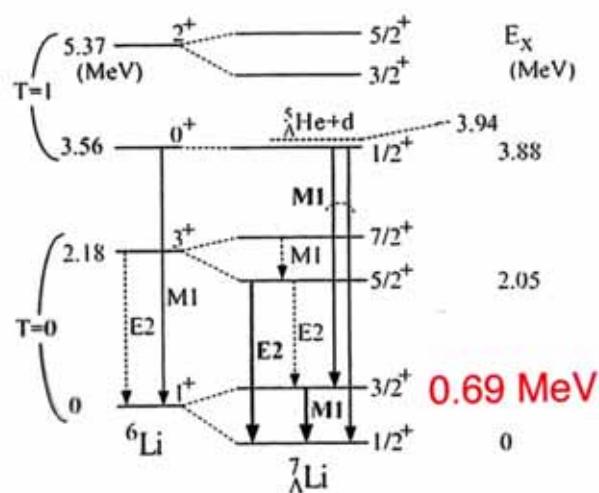
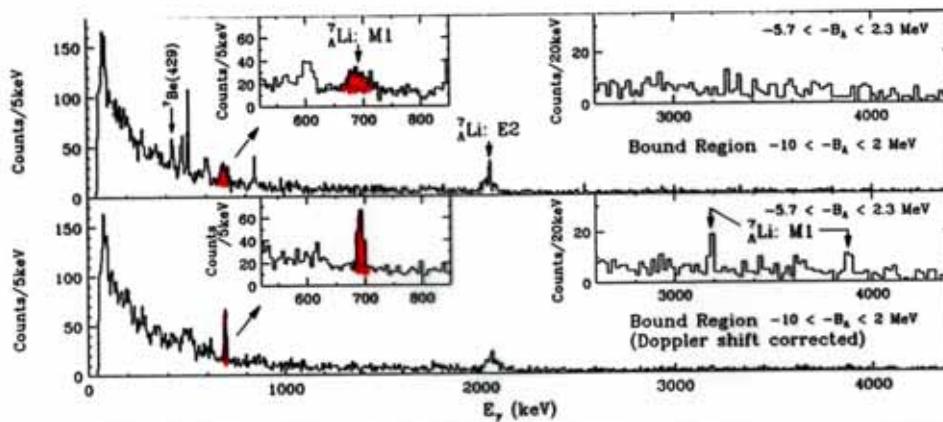


Variational Monte-Carlo

J. Lomnitz-Adler, V.R. Pandharipande & R.A. Smith,
 Nucl. Phys. **A361** (1981) 399.

ΔN spin-spin interaction

H. Tamura et al., Phys. Rev. Lett. **84** (2000) 5963



E. Hiyama et al., Nucl. Phys. **A639** (1998) 173c

$\Sigma^0 \uparrow$ p_\uparrow
 $\Lambda \uparrow$ p_\uparrow

$\Sigma^0 \uparrow$ p_\downarrow
 $\Lambda \uparrow$ p_\downarrow

$$-\sqrt{\frac{1}{3}} V^t$$

$$-\sqrt{\frac{1}{12}} \{V^t + V^s\}$$

Proton contribution

$\Sigma^0 \uparrow$ n_\uparrow
 $\Lambda \uparrow$ n_\uparrow

$\Sigma^0 \uparrow$ n_\downarrow
 $\Lambda \uparrow$ n_\downarrow

$$\sqrt{\frac{1}{3}} V^t$$

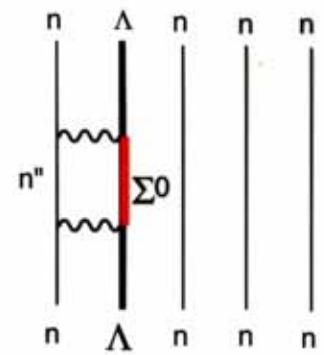
$$\sqrt{\frac{1}{12}} \{V^t + V^s\}$$

Neutron contribution

Cancellation
due to isospin selection.

$$\left\{ P_{\Sigma N}^+(^3S_1) V^t P_{\Lambda N}(^3S_1) + P_{\Sigma N}^+(^1S_0) V^s P_{\Lambda N}(^1S_0) \right\} \Lambda_\uparrow^+ p_\uparrow^+ p_\downarrow^+ n_\uparrow^+ n_\downarrow^+ |0\rangle$$

g-matrix



$$\sqrt{\frac{1}{T+1}}$$

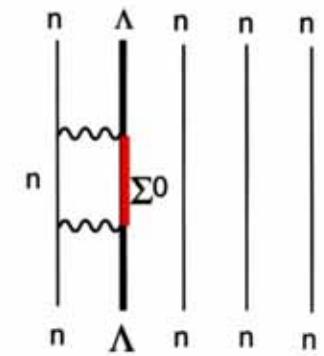
$$\sqrt{\frac{T}{T+1}} \quad \text{for } T = T_z = \frac{N}{2}$$

Σ^-

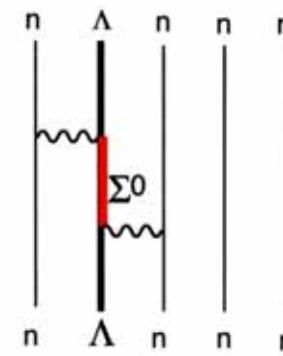
Σ^0

Σ^+

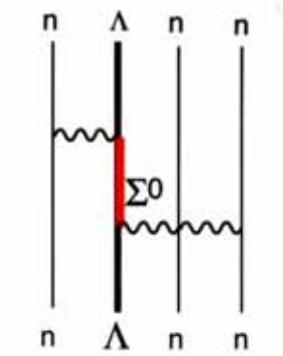
Coherently enhanced 3BF



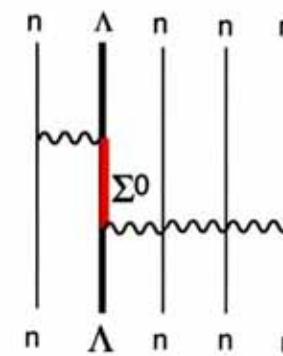
+

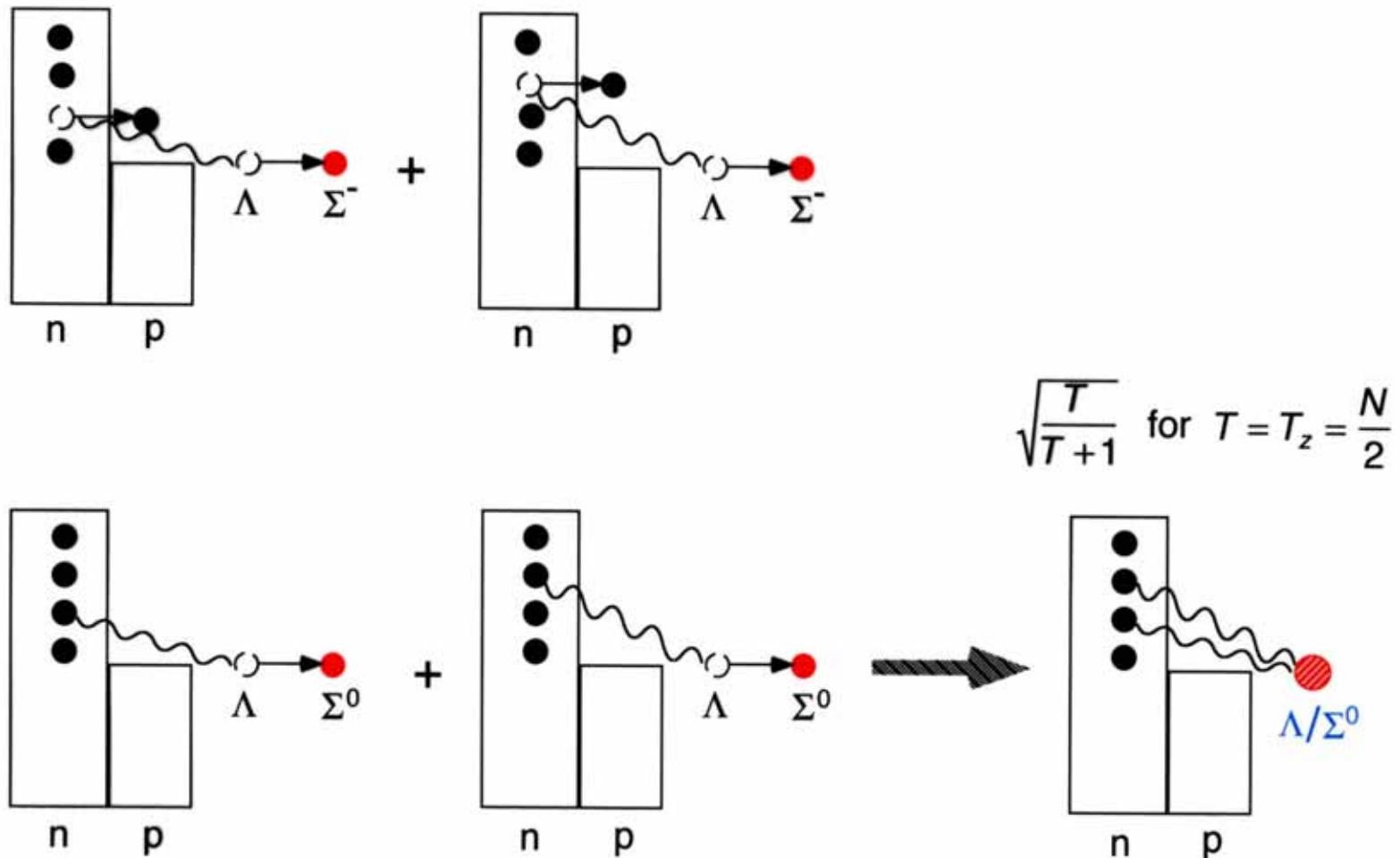


+

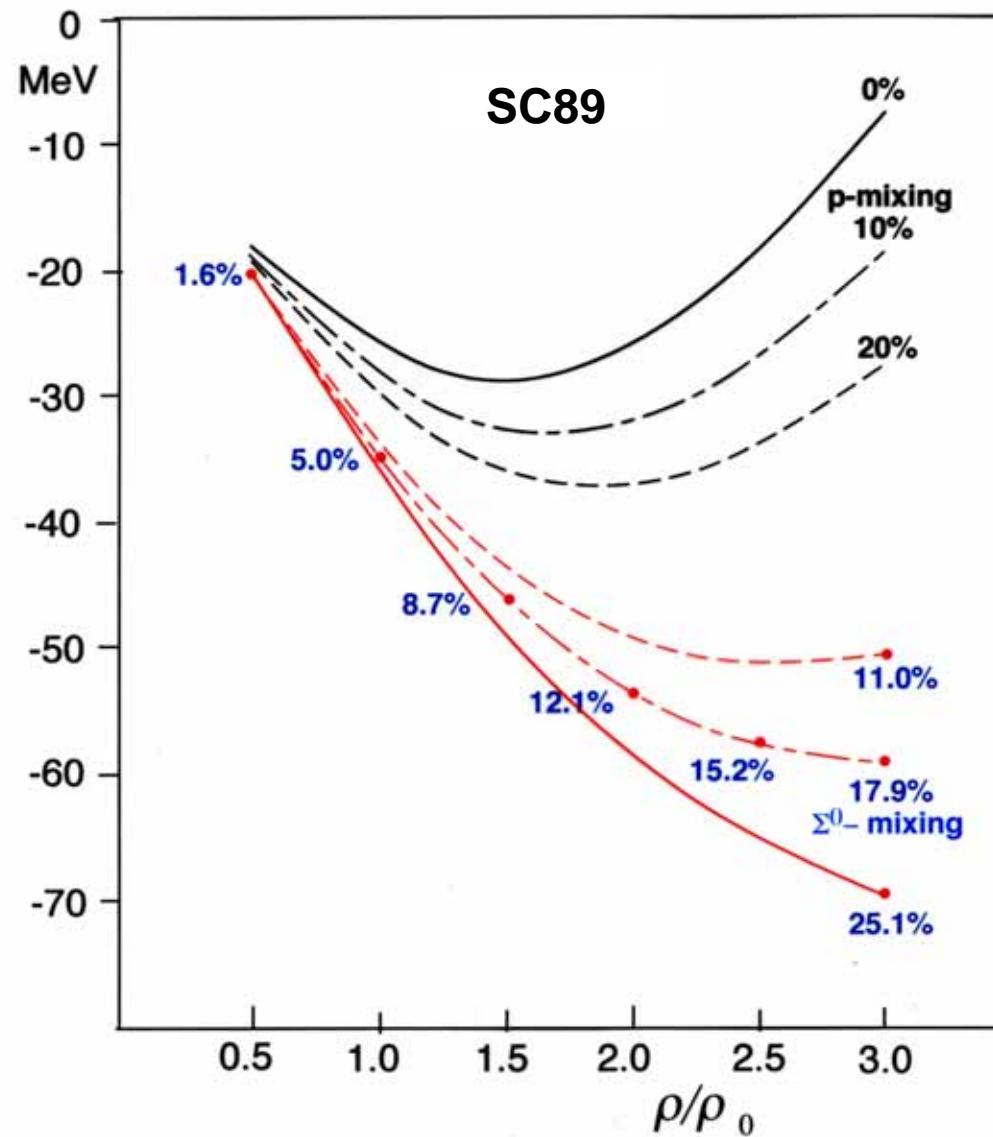


+



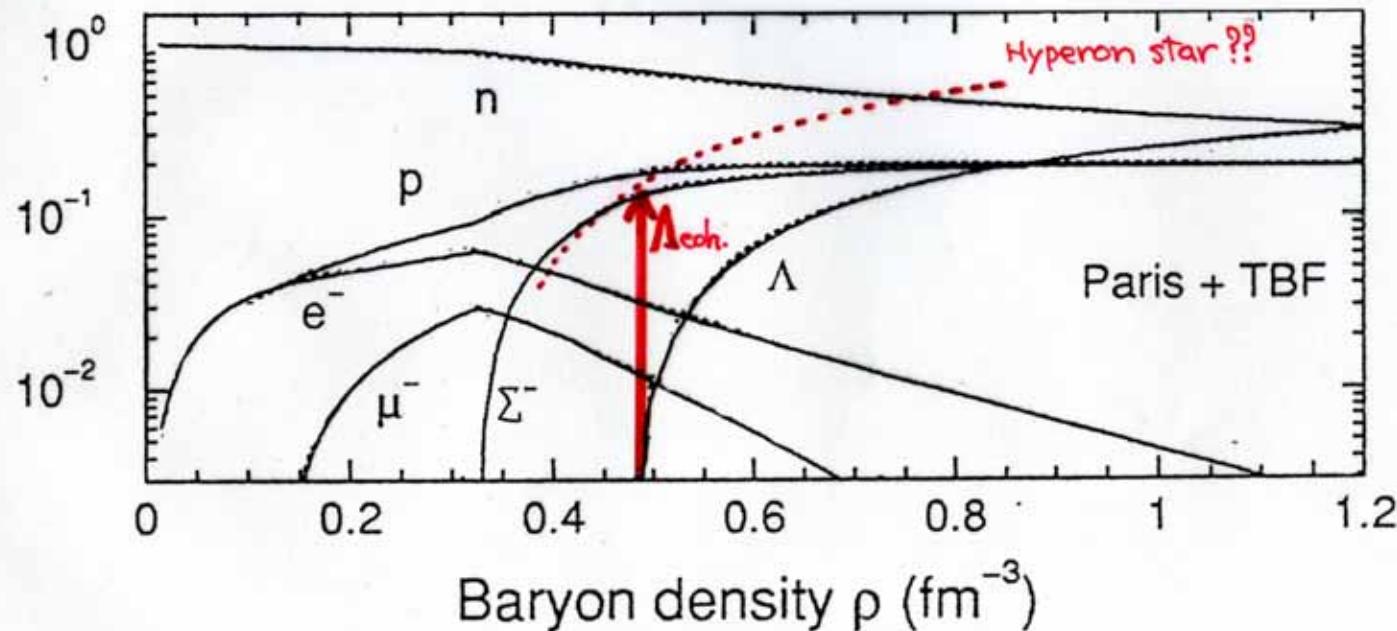


U_Λ in neutron matter



Composition of neutron star matter

M. Baldo and G.F. Burgio, Phys. Rev. C61 (2000) 055801.
(Brueckner-Bethe-Goldstone theory)



Relativistic Mean-Field Model

Lagrangian density

$$L = L_B^0 + L_M^0 + L_{\text{int}}$$

Baryons : n, p, Λ , Σ^0 , Σ^-

Mesons : σ , ρ , ω

For Λ and Σ^0

$$(\rho - \gamma^0 g_{\Lambda\Lambda\omega} \omega_0 - M_\Lambda + g_{\Lambda\Lambda\sigma} \sigma) \Lambda - [\gamma^0 g_{\Lambda\Sigma\rho} \rho_0] \Sigma^0 = 0$$

$$(\rho - \gamma^0 g_{\Sigma\Sigma\omega} \omega_0 - M_\Sigma + g_{\Sigma\Sigma\sigma} \sigma) \Sigma^0 - [\gamma^0 g_{\Sigma\Lambda\rho} \rho_0] \Lambda = 0$$

For mesons

$$m_\sigma^2 \sigma = \sum g_{BB\sigma} \langle \bar{B}B \rangle$$

$$m_\omega^2 \omega^0 = \sum g_{BB\omega} \langle \bar{B}\gamma^0 B \rangle$$

$$m_\rho^2 \rho^0 = \sum g_{BB\rho} \langle \bar{B}\gamma^0 B \rangle + [g_{\Lambda\Sigma\rho} (\langle \bar{\Lambda}\gamma^0 \Sigma \rangle + \langle \bar{\Sigma}\gamma^0 \Lambda \rangle)]$$

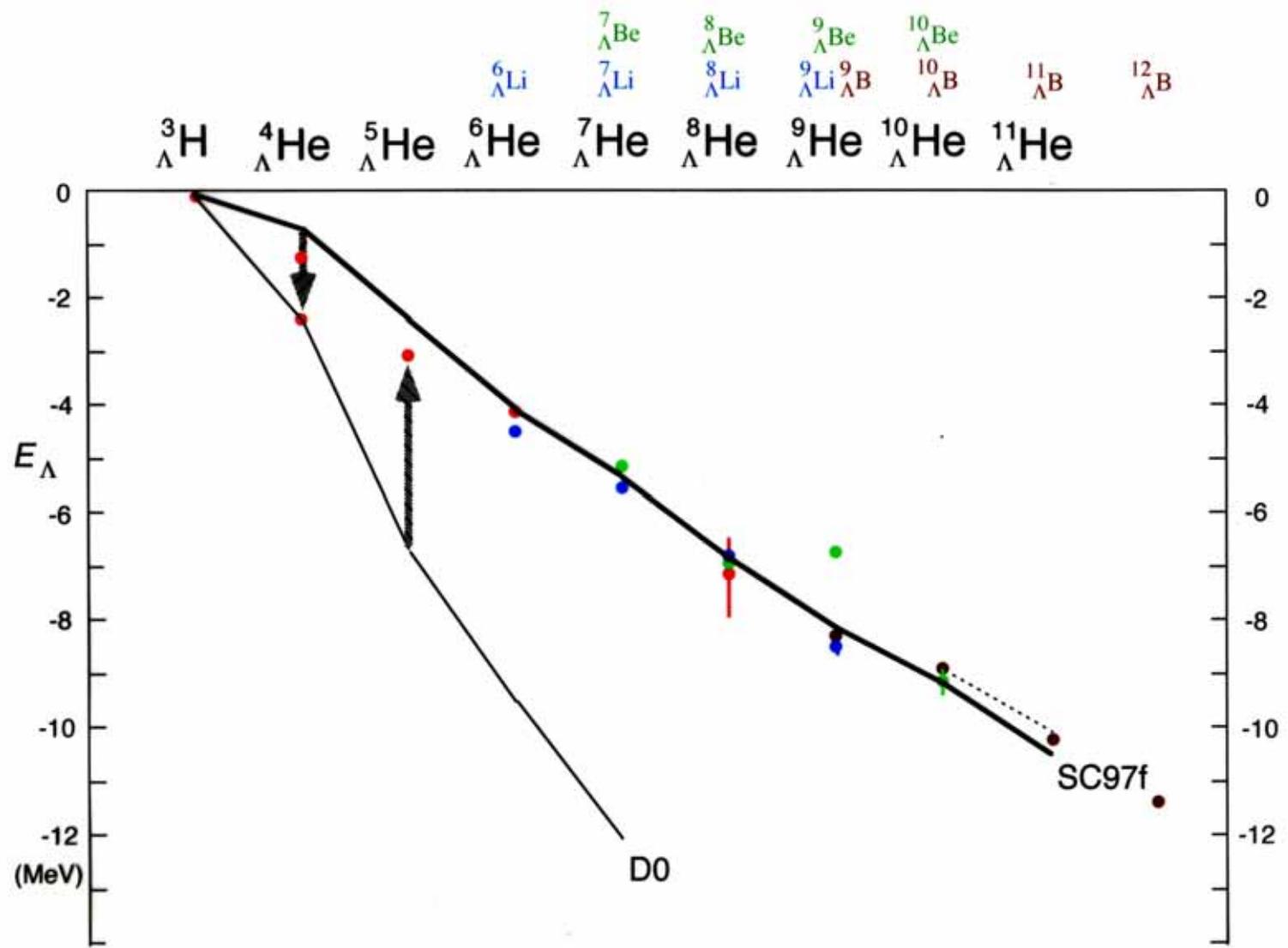
The **normal** state of infinite matter :

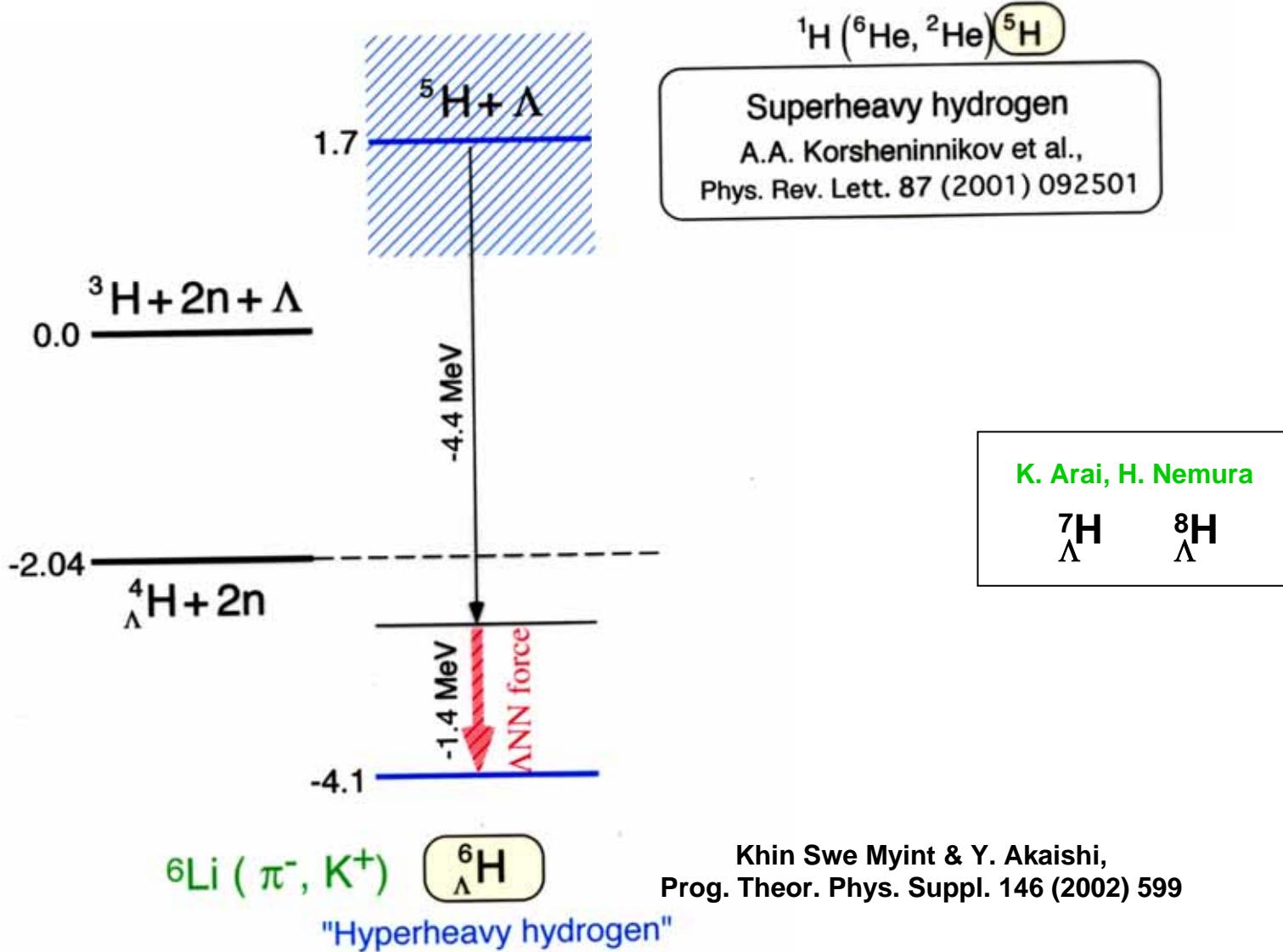
N.K. Glendenning, *Astrophys. J.* **293** (1985) 470.

Baryons in the medium **carry the same Q.N.'s as in vacuum.**

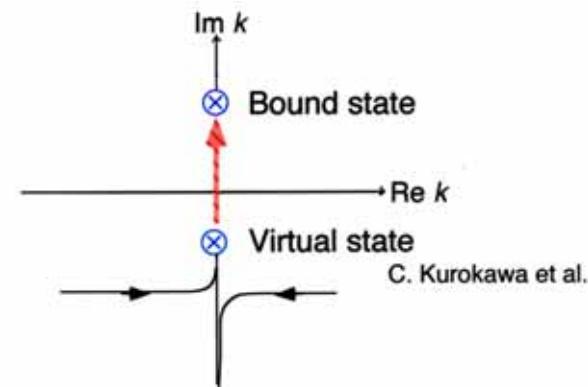
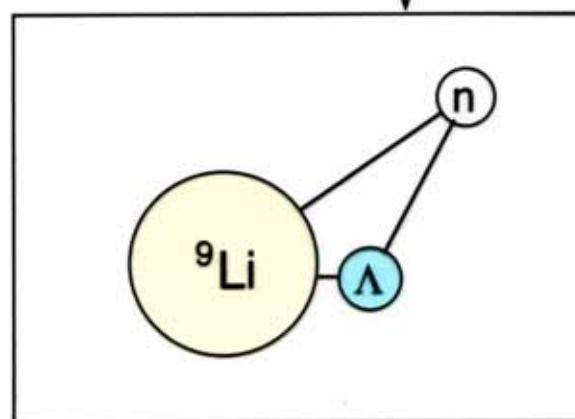
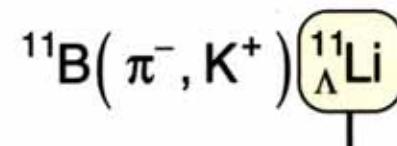
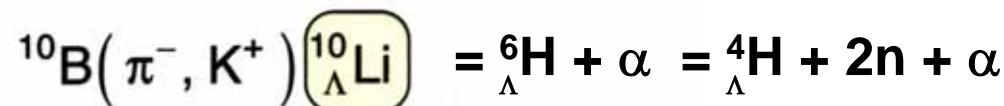
$$\boxed{\rho} = 0$$

M. Serra





Double-charge & Strangeness Exchange Reaction



YN interaction from p-shell nucleons

$$\frac{22}{27} {}^3g_{\Sigma N}^{E,O} + \frac{14}{27} {}^1g_{\Sigma N}^{E,O} : (T = \frac{3}{2})$$

$$\frac{23}{27} {}^3g_{\Sigma N}^{E,O} + \frac{17}{54} {}^1g_{\Sigma N}^{E,O}$$

$$\frac{4\sqrt{5}}{9} {}^3g_{\Sigma N,\Lambda N}^{E,O} + \frac{\sqrt{5}}{18} {}^1g_{\Sigma N,\Lambda N}^{E,O}$$

$$\frac{5}{3} {}^3g_{\Lambda N}^{E,O} + \frac{5}{6} {}^1g_{\Lambda N}^{E,O}$$

${}^{10}_{\Lambda} Li (1^-)$

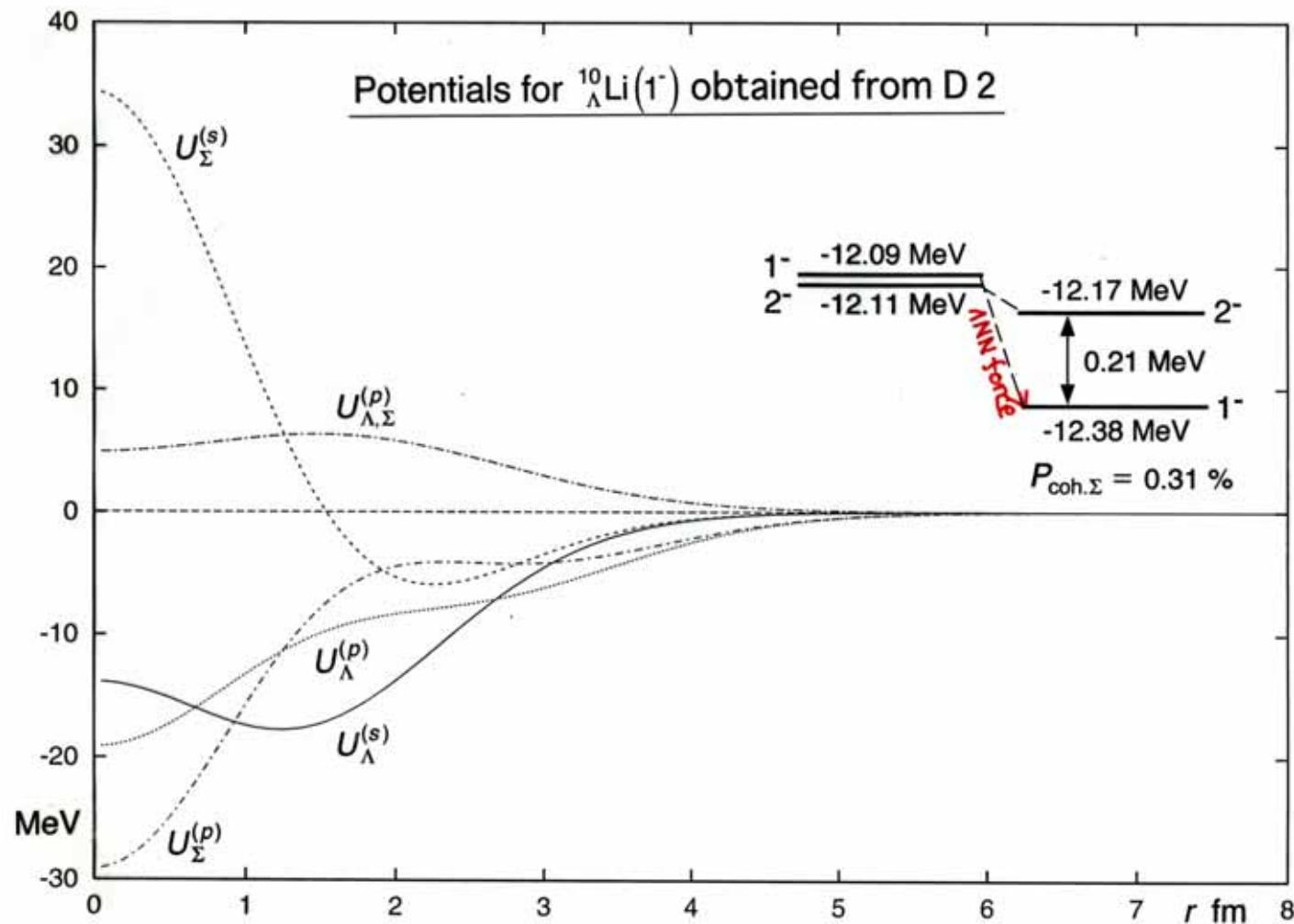
$$\frac{10}{9} {}^3g_{\Sigma N}^{E,O} + \frac{2}{9} {}^1g_{\Sigma N}^{E,O} : (T = \frac{3}{2})$$

$$\frac{8}{9} {}^3g_{\Sigma N}^{E,O} + \frac{5}{18} {}^1g_{\Sigma N}^{E,O}$$

$$\frac{\sqrt{5}}{3} {}^3g_{\Sigma N,\Lambda N}^{E,O} + \frac{\sqrt{5}}{6} {}^1g_{\Sigma N,\Lambda N}^{E,O}$$

$$2 {}^3g_{\Lambda N}^{E,O} + \frac{1}{2} {}^1g_{\Lambda N}^{E,O}$$

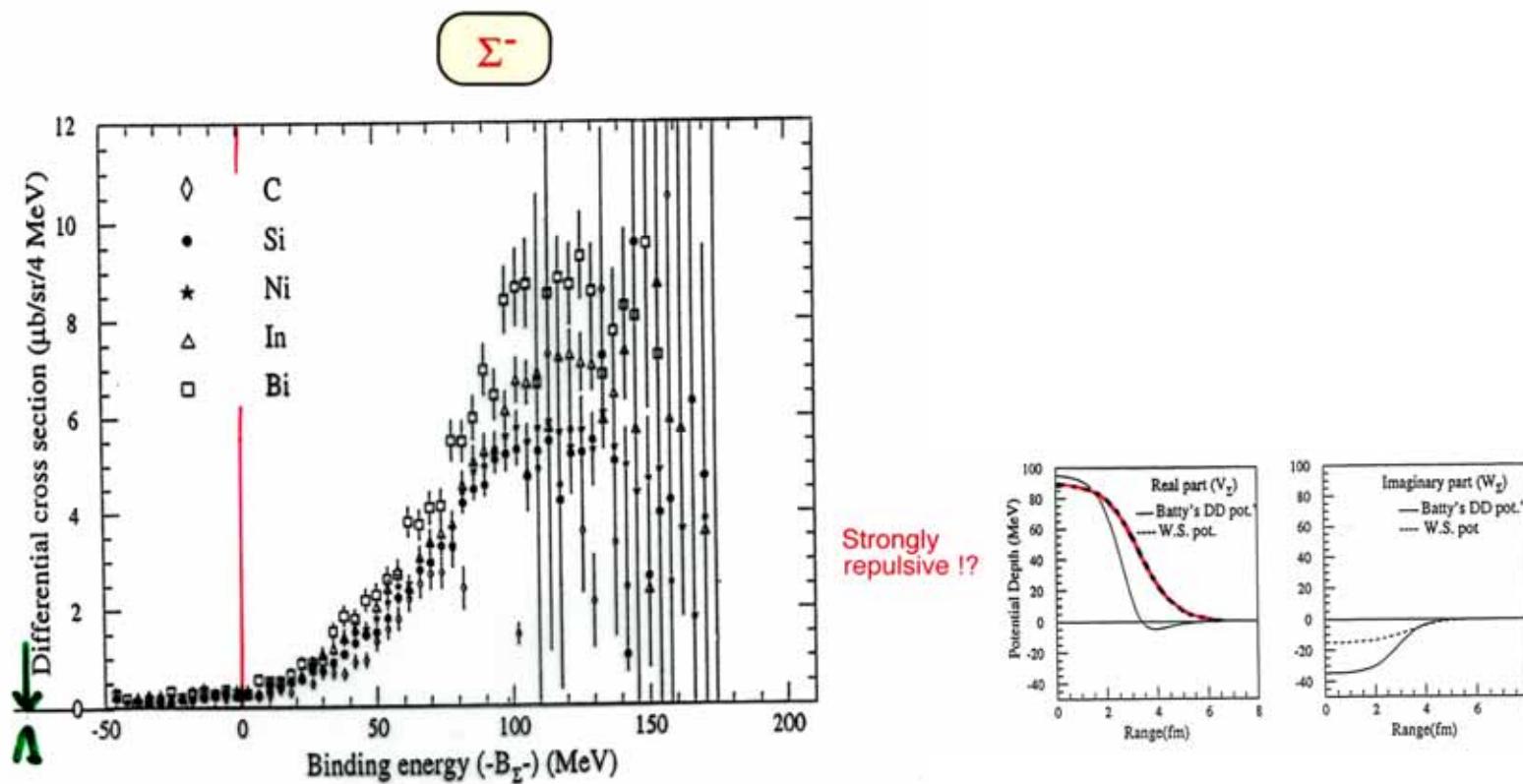
${}^{10}_{\Lambda} Li (2^-)$



AMD cal.; Y. Enyo, A. Dote et al.

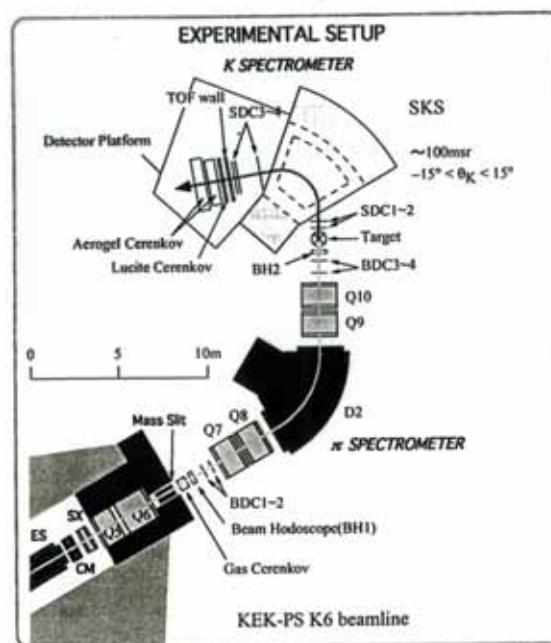
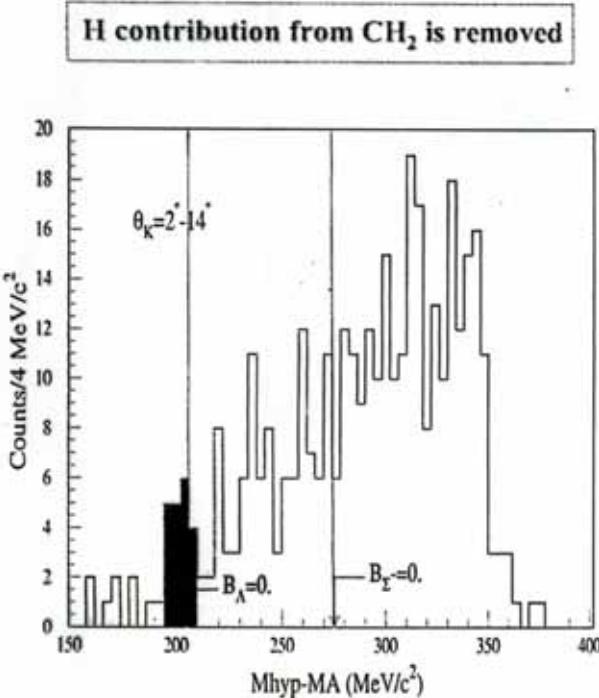
Double-Charge Exchange Experiment (π^-, K^+) at KEK

H. Noumi, P.K. Saha et al., Phys. Rev. Lett. **89** (2002) 072301

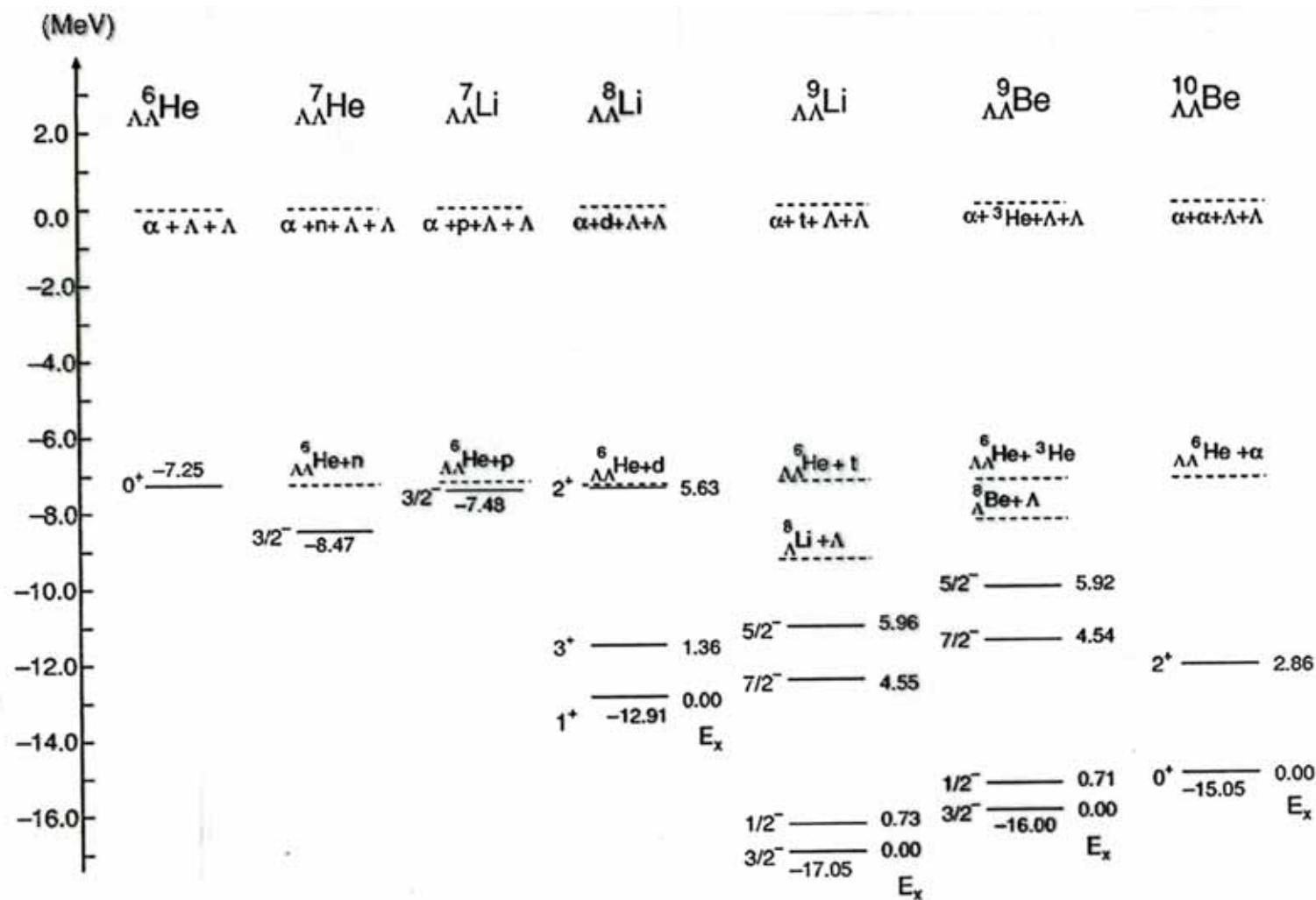


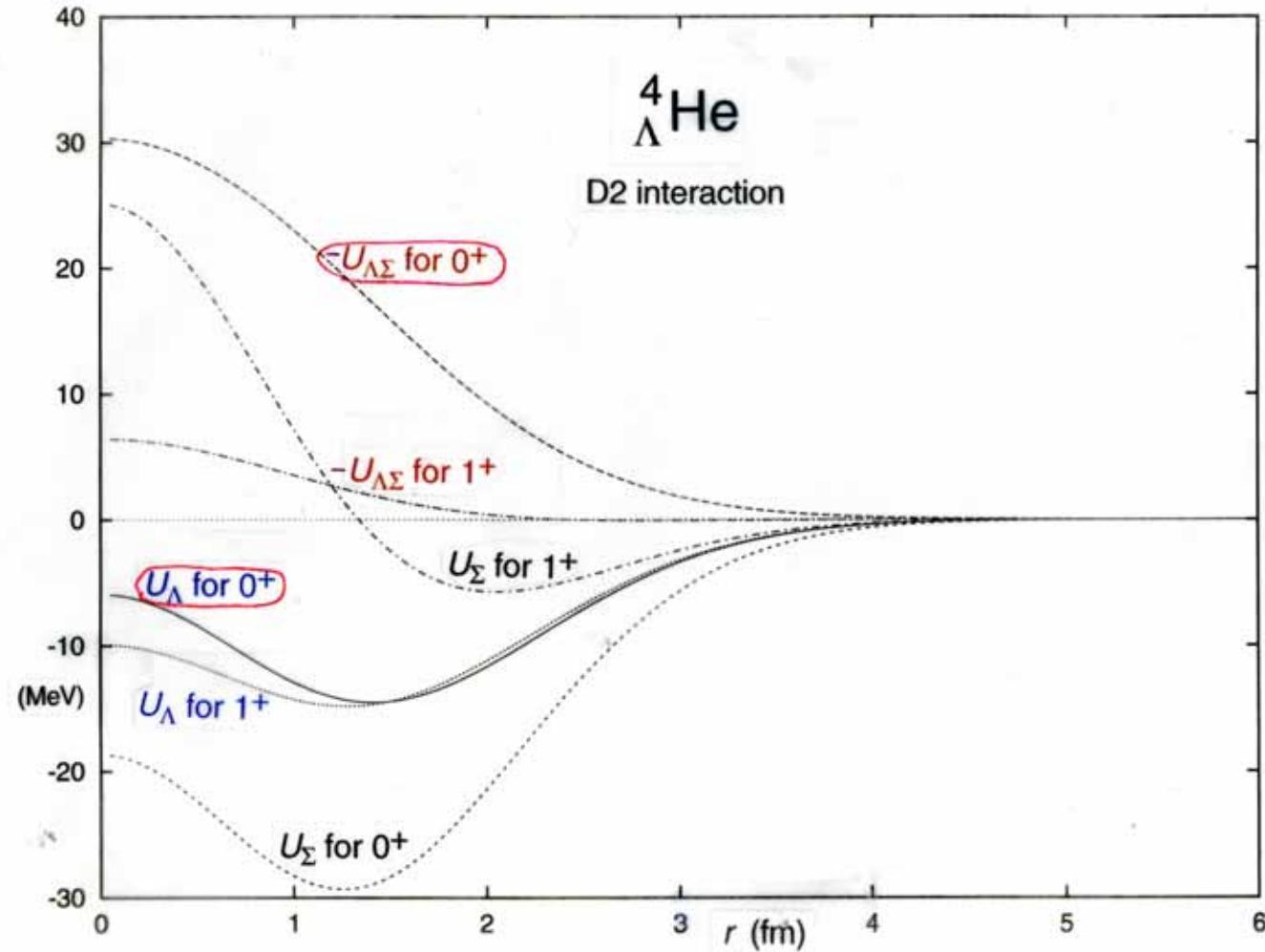
$^{12}\text{C}(\pi^-, \text{K}^+)$ spectrum from the CH_2 target

by Saha, Fukuda, Noumi

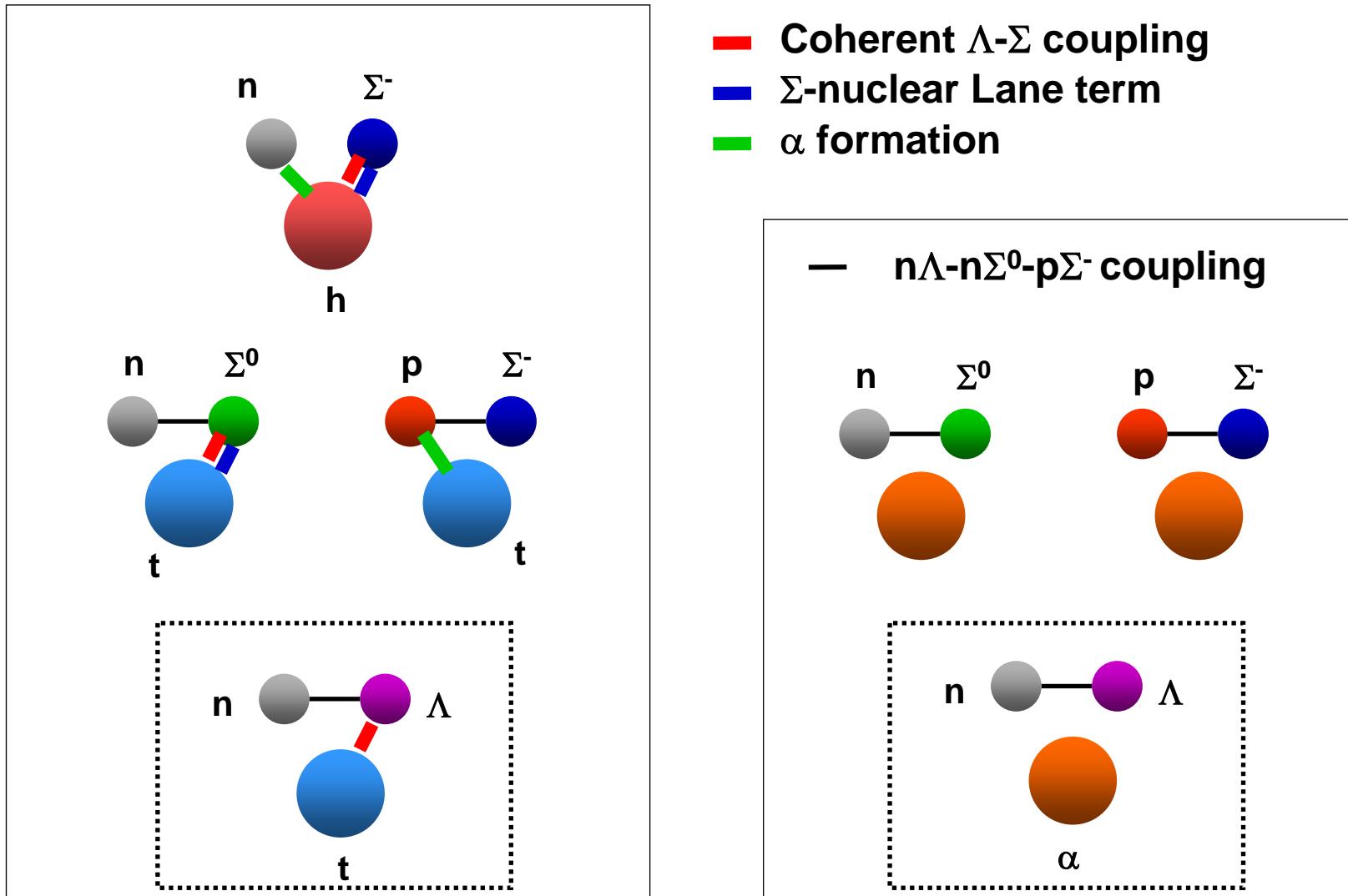


$^{10}\text{B} (\pi^-, \text{K}^+)$; W. Imoto, P.K. Saha et al.; this evening

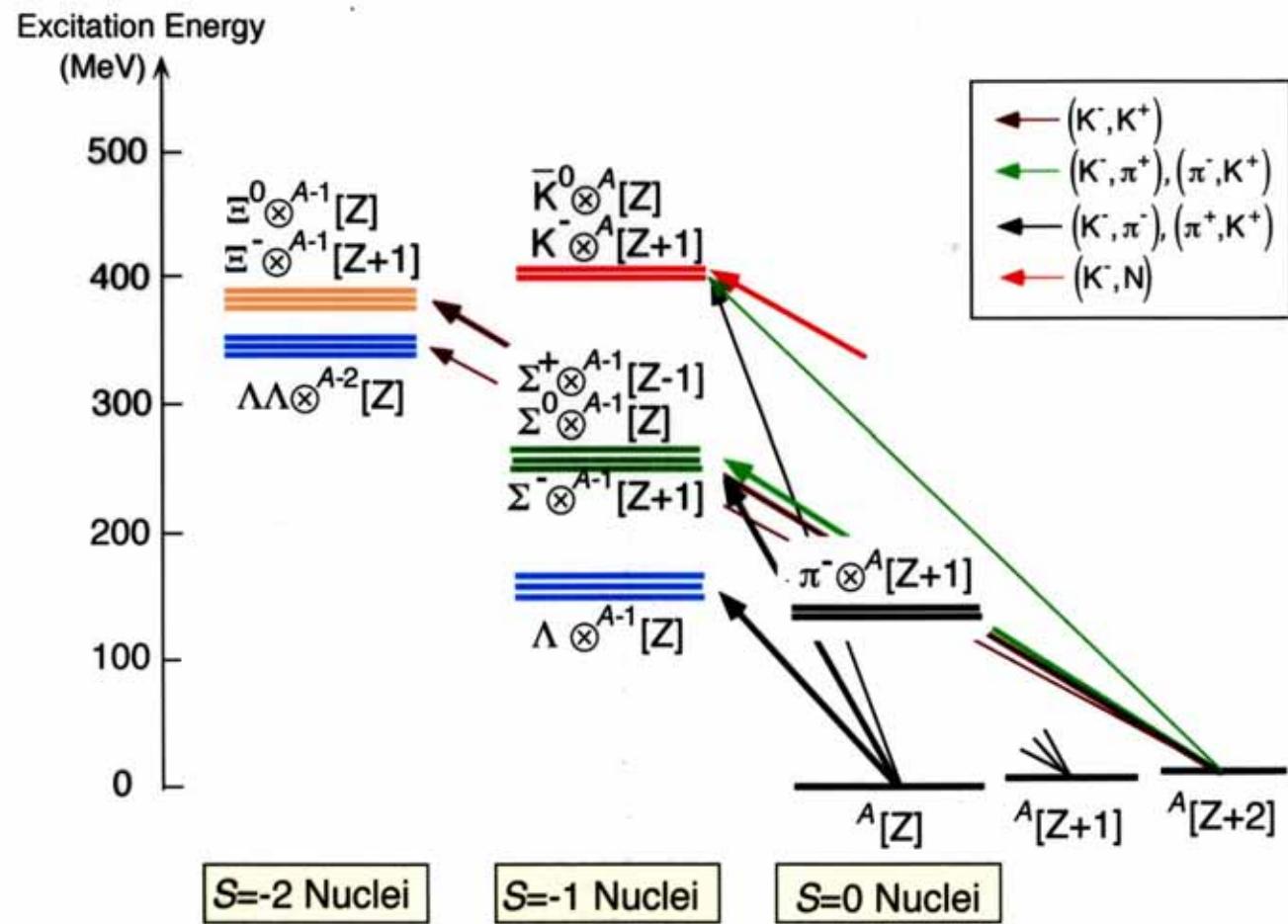




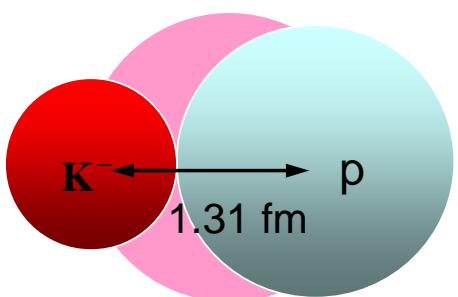
Coupling Scheme



Yamazaki's diagram



$\Lambda(1405)$



$p p K^-$

