

Theoretical overview of the Pentaquarks

Makoto Oka

Tokyo Institute of Technology



3rd CNS International Summer School

8/19 2004

Contents

- 1. Introduction*
- 2. QCD prediction*
- 3. PP vs NP*
- 4. Conclusion*



New Narrow Hadrons found in 2003

- ✦ Θ^+ (1540) by LEPs, SPring-8 ($\Gamma < 10 \text{ MeV}$)
- ✦ Ξ^{--} (1862) ($\Gamma < 18 \text{ MeV}$), $\Xi_c^0, \dots$
- ✦ X (3872) by Belle, KEK ($\Gamma < 2.3 \text{ MeV}$)
- ✦ D_s^* (2317) and $D_s^*(2463)$ by BaBar ($\Gamma < 10 \text{ MeV}$)
- ✦ $K^- \text{ppn } (^3\text{K}_H)$ deep bound state at 12 GeV PS, KEK

$BE = 173 \text{ MeV}$ ($\Gamma < 25 \text{ MeV}$) and more

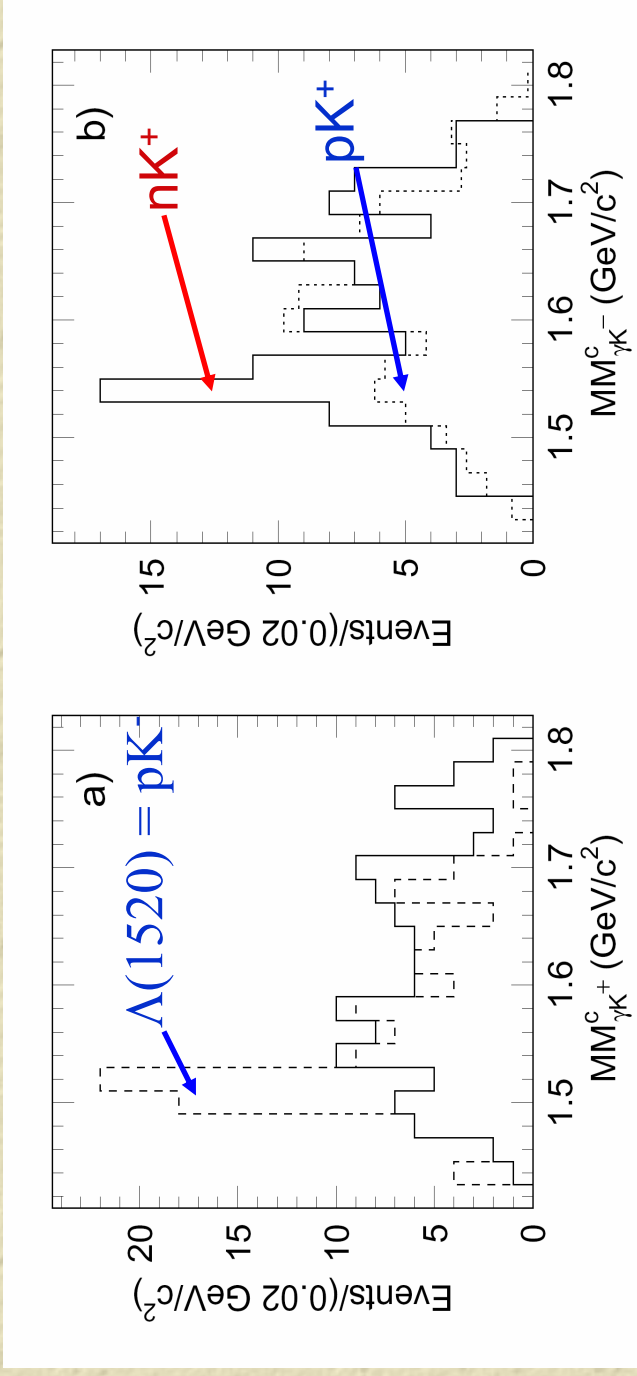
narrow

contain heavy quarks



Discovery of Θ^+ at SPring-8

T. Nakano et al. (LEPS collaboration)



Whoever believes ...

Skepticism

Why not seen before?

should have many brothers

Some bad history of narrow exotics

baryonium, dibaryon, ...

Why so narrow?

Fall-apart decay theory

Renaissance of Hadron Spectroscopy

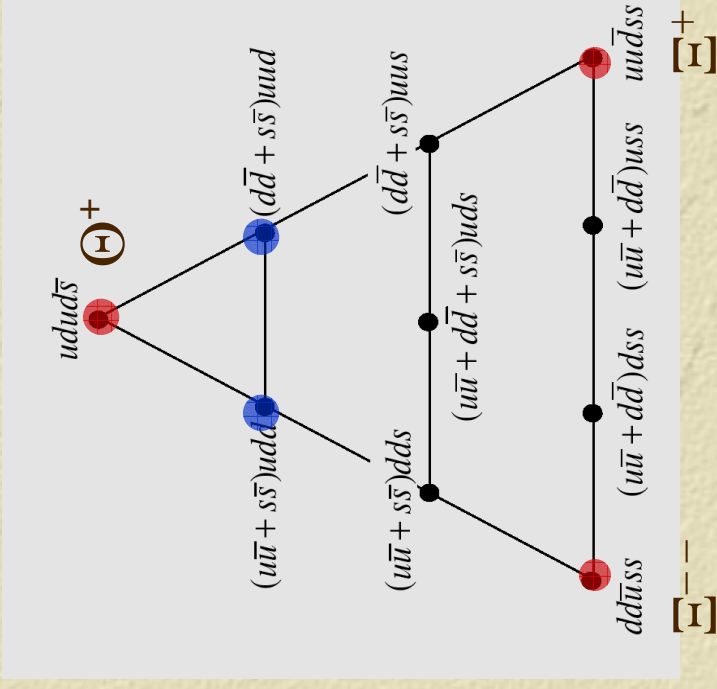


Facts of Θ^+

- ✧ $M \approx 1530\text{--}1540\text{MeV}$, $\Gamma < 9\text{ MeV}$
- ✧ $S = +1$, $B = 1$, $Q = +1$ $\longrightarrow$ $uudd\bar{s}$
 $\longrightarrow Y = B + S = 2$
- ✧ No pK^+ ($I_3 = +1$) state $\longrightarrow I = 0$
- ✧ Simplest $SU(3)$ irrep. is $F = 10^*$
 unique for ≤ 5 quarks
- ✧ No lower $S = +1$ states

This is the ground state!

NK is the only open channel



Main issues

✠ Mass spectrum

$$M(\Theta^+) = 1540 \text{ MeV}, \quad M(\Xi^{--}) = 1862 \text{ MeV}$$

✠ Width $\Gamma(\Theta^+) \sim 0.9 \text{ MeV}$

very narrow for a resonance at 100 MeV above the threshold ex. $N^*(1650)$ S_{11} $\Gamma \sim 150 \text{ MeV}$

✠ Spin, Parity, $SU(3)$ irrep., . . .undetermined yet

✠ Production mechanisms

low energy vs high energy



Main issues

✠ Parity

pentaquark $q^4 + \bar{q}$ ($\pi = -$)

$L = 0$

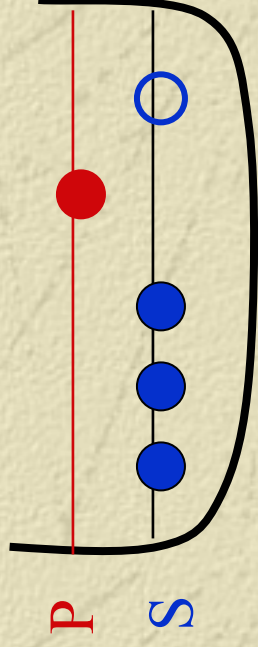
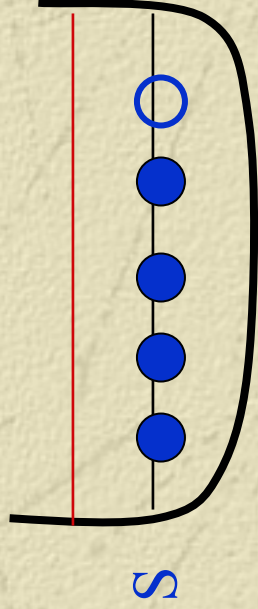
$\pi = -$

NK $L = \text{even}$ (S, D, ...)

$L = 1$

$\pi = +$

NK $L = \text{odd}$ (P, F, ...)



Pentaquarks in QCD

Lattice QCD

Csikor, Fodor, Katz, Kovacs, JHEP 0311 (2003) 070
Sasaki, hep-lat/0310014
Chiu, Hsieh, hep-ph/0403020
Mazur et al., hep-ph/0406196
Ishii et al., hep-lat/0408030

QCD sum rule for Θ^+

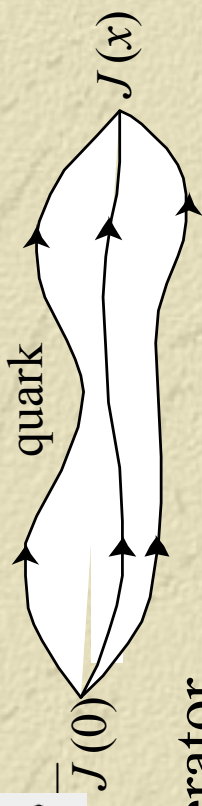
Zhu, Phys. Rev. Lett. 91 (2003) 232002
Matheus et al., Phys. Lett. B578 (2004) 323
No parity projection was made.
Sugiyama, Doi, Oka, Phys. Lett. B581 (2004) 167
determine mass and parity



Two Point Correlators

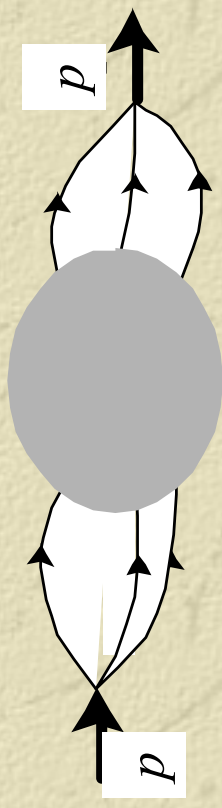
- TPC $\Pi(x)$ contains information of the spectrum

$$\Pi(x) = \langle 0 | T(J(x) \bar{J}(0)) | 0 \rangle$$



- $J(x)$: interpolating field operator
determines quantum numbers

- Fourier transform



$$\Pi(p) \equiv i \int d^4x e^{ip \cdot x} \langle 0 | T(J(x) \bar{J}(0)) | 0 \rangle$$

Two Point Correlators



interpolating field operator

Mesons

$I=1$ vector meson

$I=1$ pseudoscalar meson

Baryons

$J=1/2$ baryon

$J=1/2$ pentaquark

$$J_\rho(x) = \bar{q}(x) \gamma^\mu \frac{\vec{\tau}}{2} q(x)$$

$$J_\pi(x) = \bar{q}(x) \gamma^5 \frac{\vec{\tau}}{2} q(x)$$

$$B(x) = \epsilon_{abc} (u_a^T(x) C \gamma^5 d_b(x)) u_c(x)$$

$$\Theta^+ \equiv \epsilon_{abc} \epsilon_{ade} \epsilon_{bfg} [u_d^T C \gamma^5 d_e] [u_f^T C d_g] C \bar{s}_c^T$$

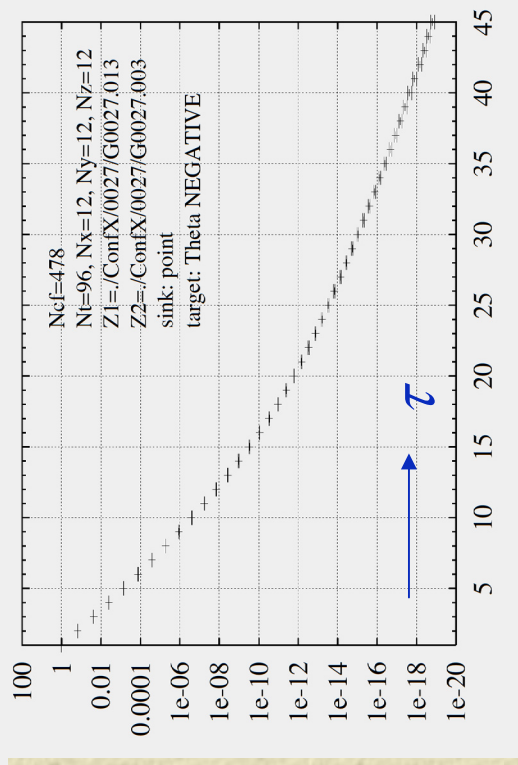


Lattice QCD

✪ Imaginary time TPC

$$\Pi(\vec{p} = 0; \tau) = \int d^3x \langle 0 | T(J(\vec{x}, \tau) \bar{J}(0, 0)) | 0 \rangle$$

$$= \sum_i \lambda_i^2 e^{-m_i \tau} \xrightarrow{\text{large } \tau} \boxed{\lambda_0^2 e^{-m_0 \tau}}$$



Lattice QCD for Θ^+

✠ Lattice QCD for Θ^+

F. Csikor, S. Sasaki, T.-W. Chiu, N. Mathur
N. Ishii

- ◆ No $1/2^+$ ($I = 0$) resonance state below 2 GeV.
(except *Chiu*)
- ◆ Existence of $1/2^-$ ($I = 0$) state is under question.



Lattice QCD for Θ^+

Sasaki, Ishii, Chiu

✠ interpolating field $J=1/2, I=0$

$$\Theta^+ \equiv \epsilon_{abc}\epsilon_{ade}\epsilon_{bfg}[u_d^T C \gamma^5 d_e][u_f^T C d_g] C \bar{s}_c^T$$

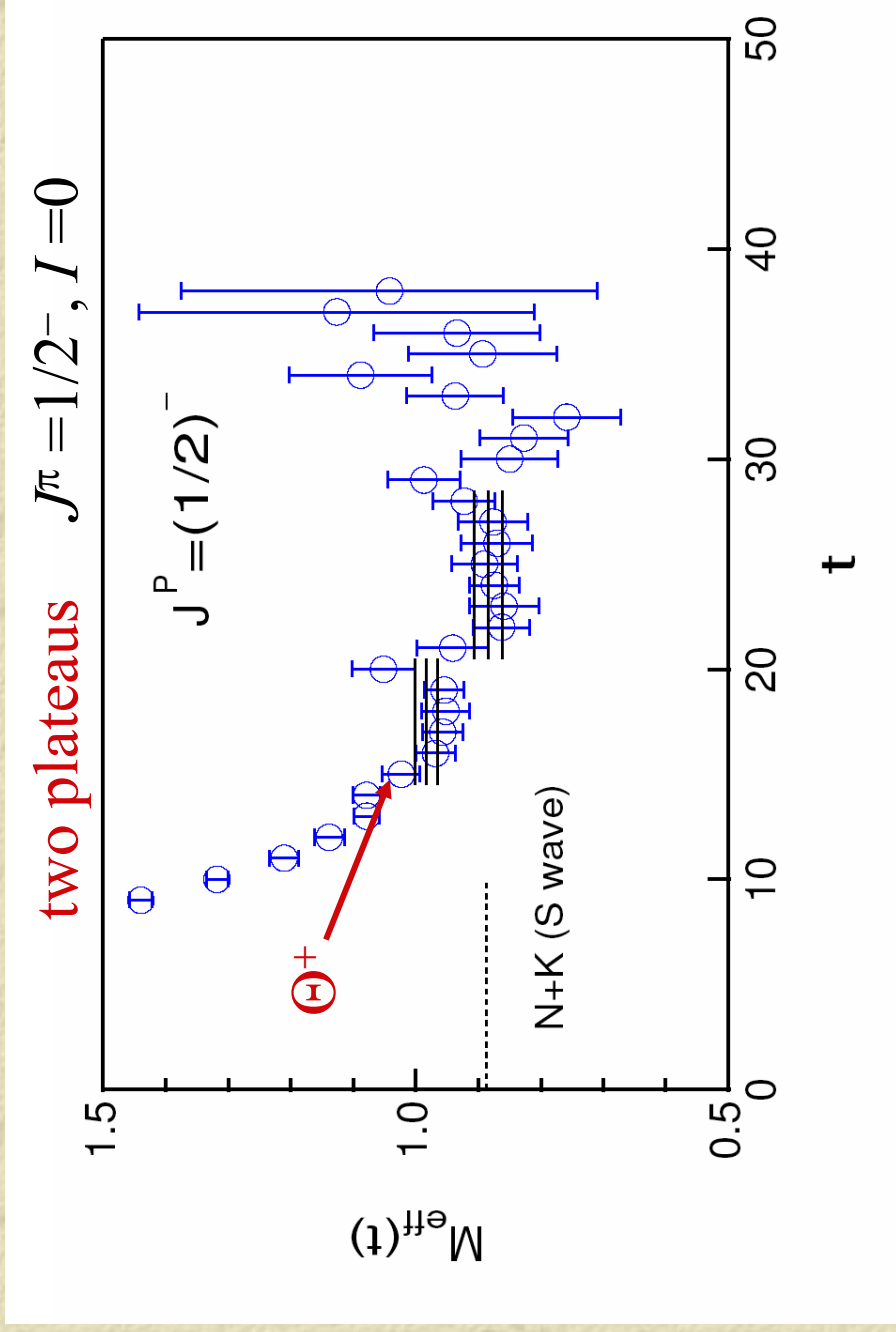
(diquark)²-s type

$$\begin{aligned} \gamma^5 \Theta^+ &= \epsilon_{abc}\epsilon_{ade}\epsilon_{bfg}[u_d^T C \gamma^5 d_e][u_f^T C d_g] \gamma^5 C \bar{s}_c^T \\ &= \left(-\frac{1}{4} \times \frac{2}{3}\right) \epsilon_{deg}[u_d^T C \gamma^5 d_e] d_g \cdot (\bar{s} \gamma^5 u) \\ &\quad + \left(-\frac{1}{4} \times \frac{2}{3}\right) \epsilon_{deg}[u_d^T C \gamma^5 d_e] \gamma_\mu \gamma^5 d_g \cdot (\bar{s} \gamma^\mu u) + \dots \\ &= \left(-\frac{1}{6}\right) J_n J_{K^+} + \left(-\frac{1}{6}\right) J_n J_{K^*} + \dots \end{aligned}$$



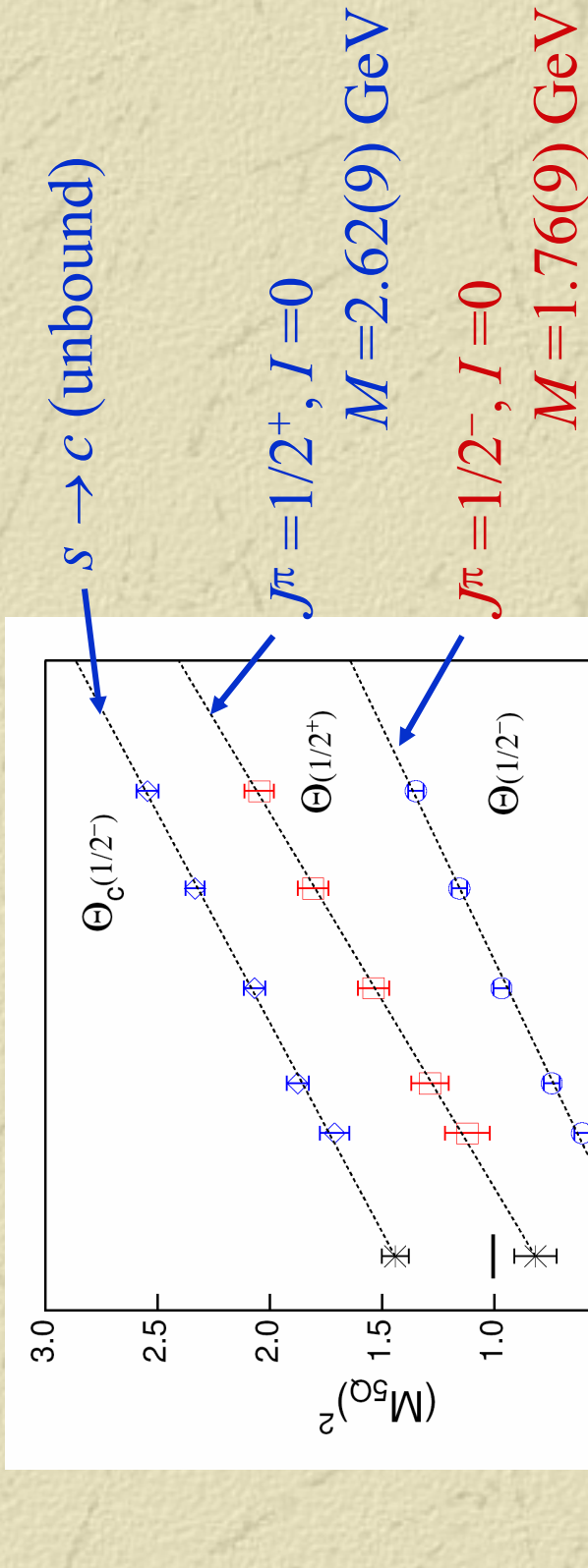
Lattice QCD for Θ^+

Sasaki, hep-lat/0310014



Lattice QCD for Θ^+

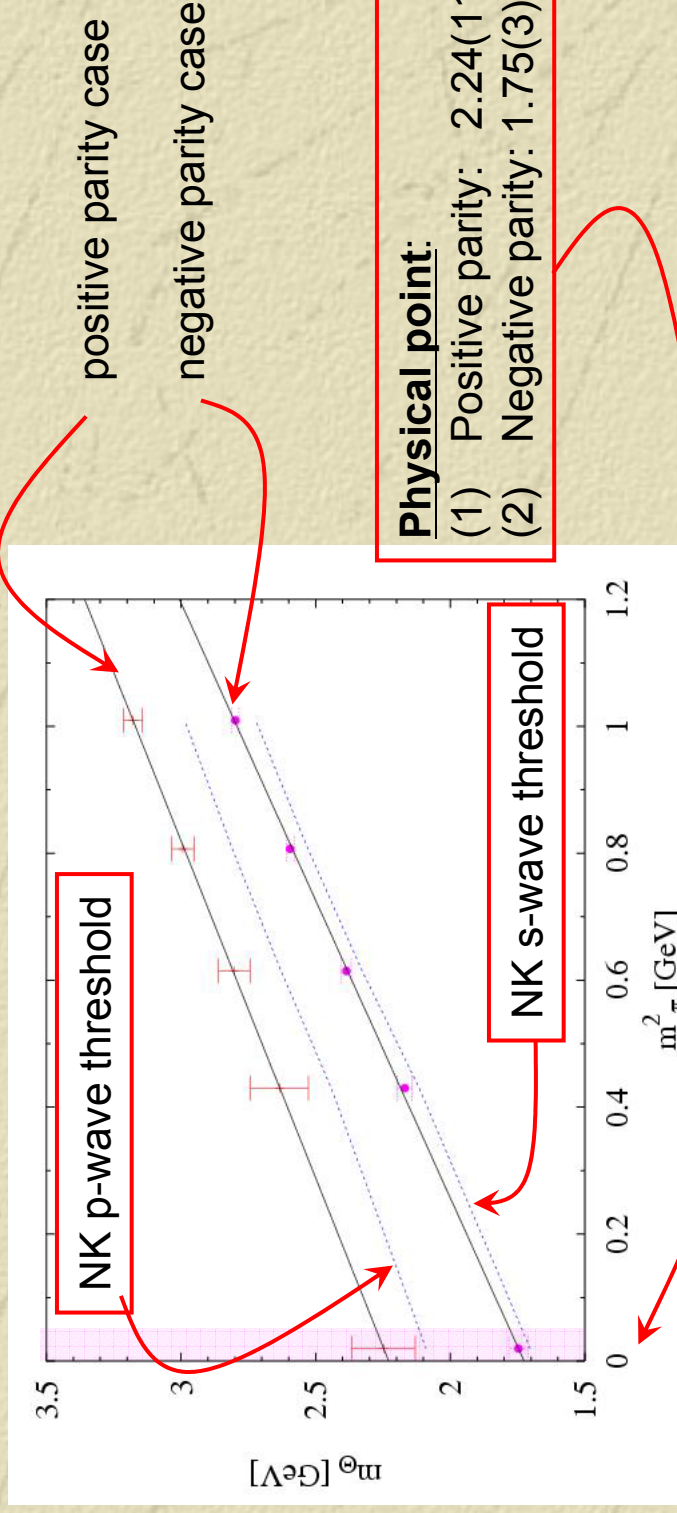
Sasaki, hep-lat/0310014



Lattice QCD for Θ^+

N. Ishii *et al.*, hep-lat/0408030

non NK operator on anisotropic lattice

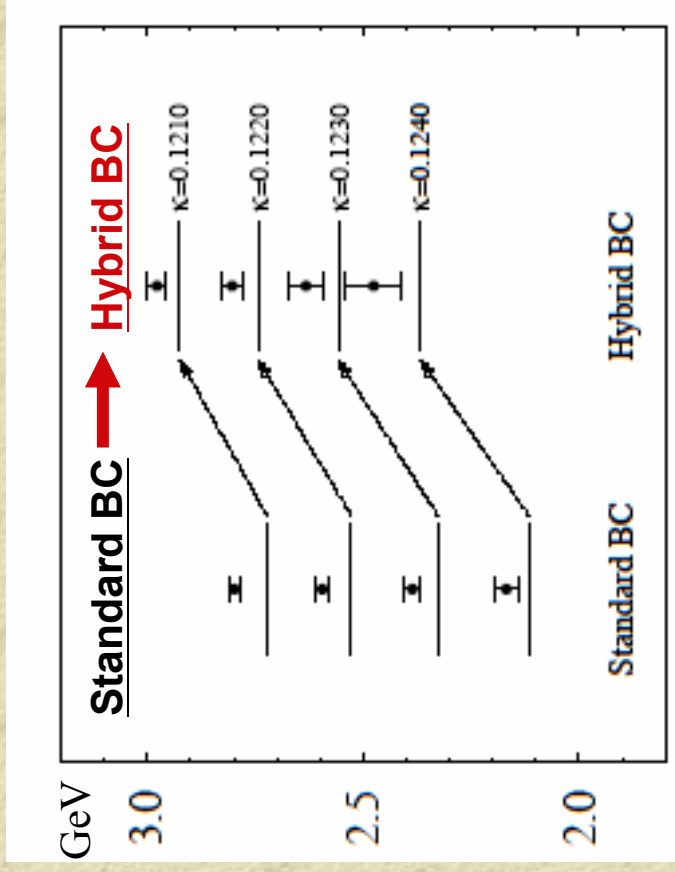


No low-lying positive parity Θ^+



Lattice QCD for Θ^+

N. Ishii *et al.*, hep-lat/0408030



Hybrid BC

antisymmetric for u, d
symmetric for s

NK threshold is artificially raised in HBC, while local pentaquarks are not affected.

The observed plateau is a **NK scattering state**. **No localized resonance state** exist in the interval:

$$m_N + m_K \leq m(\Theta^+) \leq \sqrt{m_N^2 + \vec{p}_{\min}^2} + \sqrt{m_K^2 + \vec{p}_{\min}^2}$$



QCD SR for Θ^+

✧ Deep Euclid region

$$(p_E^2 \equiv -p^2 \rightarrow \infty)$$

$$\alpha_s(p_E^2) \rightarrow 1/\ln(p_E^2) \Rightarrow \text{perturbative}$$

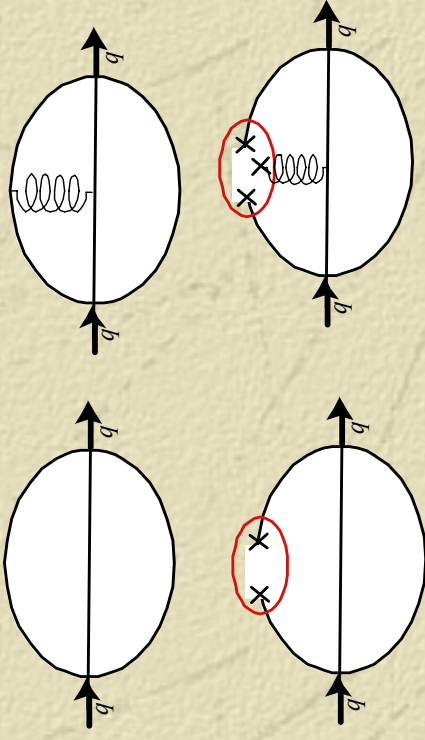
✧ Nonperturbative effects are taken into account as vacuum condensates

✧ OPE

expansion wrt $1/p_E^2$

$$\Pi(p_E^2) = \sum_n C_n(p_E^2) \langle 0 | O_n(0) | 0 \rangle$$

O_n : local operator



QCD SR for Θ^+

 physical region ($p^2 > 0$)

represented by spectral function $\rho(s)$ at $s = p^2 = m^2$

$$\rho(s) = \frac{1}{\pi} \text{Im} \Pi(s) = \sum_i \lambda_i^2 \delta(s - m_i^2)$$

i : hadrons (poles) and scattering states (cut)

 QCD SR assumes a sharp resonance + continuum ($s > s_0$)

$$\rho(s) \sim \lambda^2 \delta(s - m^2) + \theta(s - s_0) \rho(s)$$

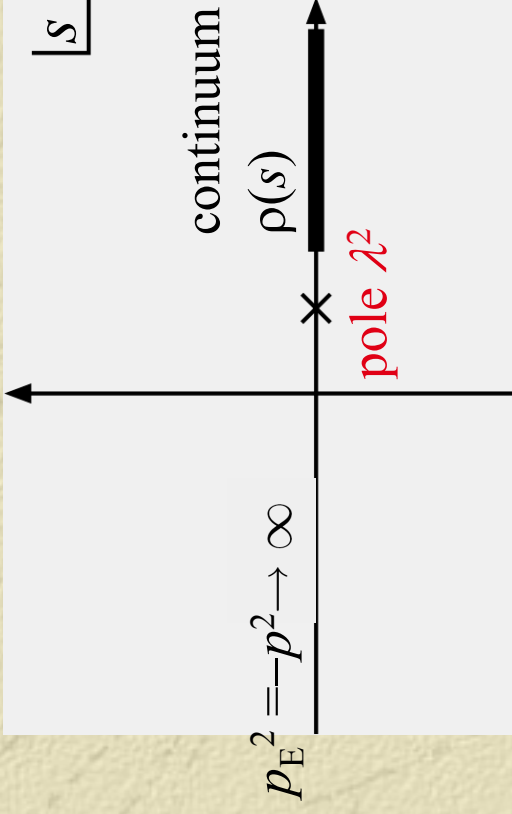
QCD SR for Θ^+

✚ Analyticity of $\Pi(p^2)$
in complex $s = p^2$

✚ analytic continuation by a dispersion
relation

$$\Pi(p_E^2) = \frac{1}{\pi} \int_0^\infty ds \frac{\text{Im } \Pi(s)}{s + p_E^2}$$

(with appropriate subtractions)



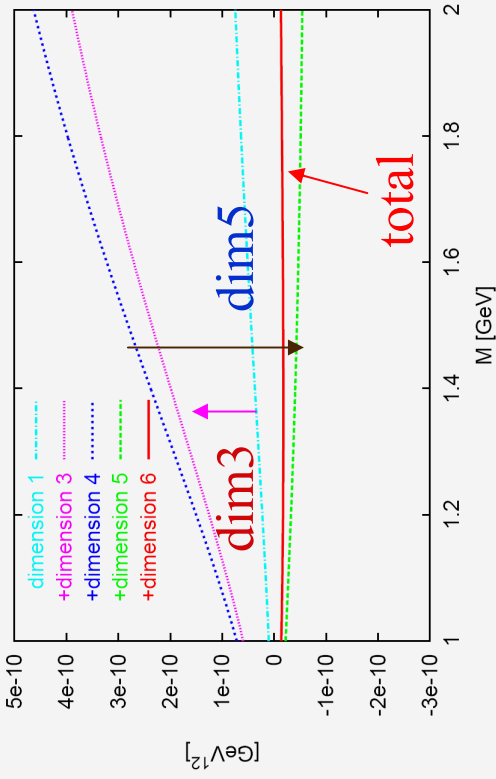
Parity of Θ^+



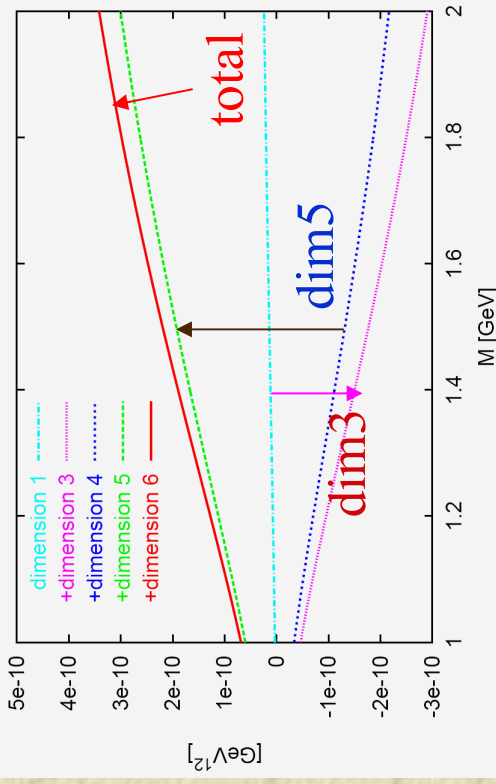
positivity of $|\lambda_{\pm}|^2 e^{-\frac{m_{\pm}^2}{M^2}}$

M : Borel mass

positive parity



negative parity



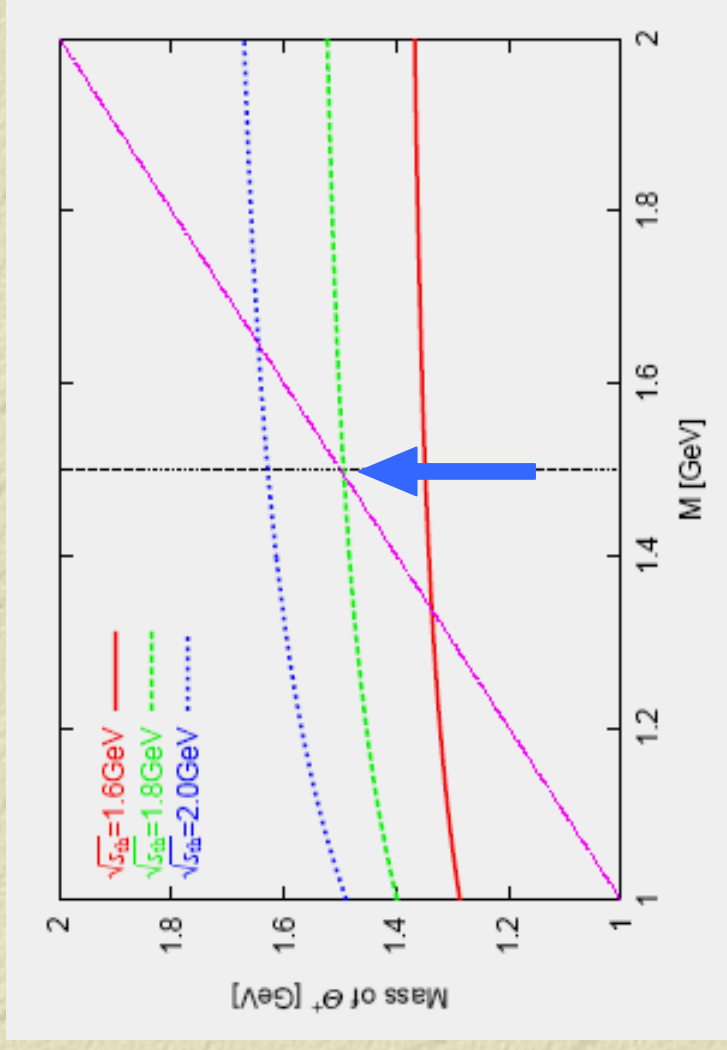
no positive parity pentaquark



Mass of Θ^+

Θ^+ mass (negative parity)

$M \sim 1.3\text{--}1.7\text{ GeV}$



Resonance or continuum?

✱ Is there a sharp state? or just NK continuum?

LQCD or SR cannot predict the width directly.

✱ It is important to choose J_{Θ} so that the coupling to NK is **weak**.

The NK continuum contribution to the *non* NK operator is less than 5% of the OPE. (S.H. Lee)

$$\langle 0 | J_{\Theta} | NK \rangle = -\frac{i}{f_{\pi}} \langle 0 | [Q_5^K, J_{\Theta}] | N \rangle = -\frac{i}{f_{\pi}} \langle 0 | J_{N;5q} | N \rangle$$



PP vs NP

✚ Why do people believe (or like) positive-parity?

- ✚ So predicted by Diakonov et al.

- ✚ To explain the tiny width

$$\Gamma(\text{NP}) \sim 300 \sim 1000 \text{ MeV}$$

Centrifugal barrier suppresses $\Gamma(\text{PP})$.

- ✚ Production cross sections

$$\sigma(\text{PP}) \sim 10 \sigma(\text{NP}) \quad \text{Hosaka et al.}$$



Chiral Soliton Model for Θ^+

✱ Excited states of the SU(3) rotated Skyrme soliton

Diakonov, Petrov, Polyakov, ZP A359 (1997) 305

$$F = 10^*$$

$$J^\pi = 1/2^+$$

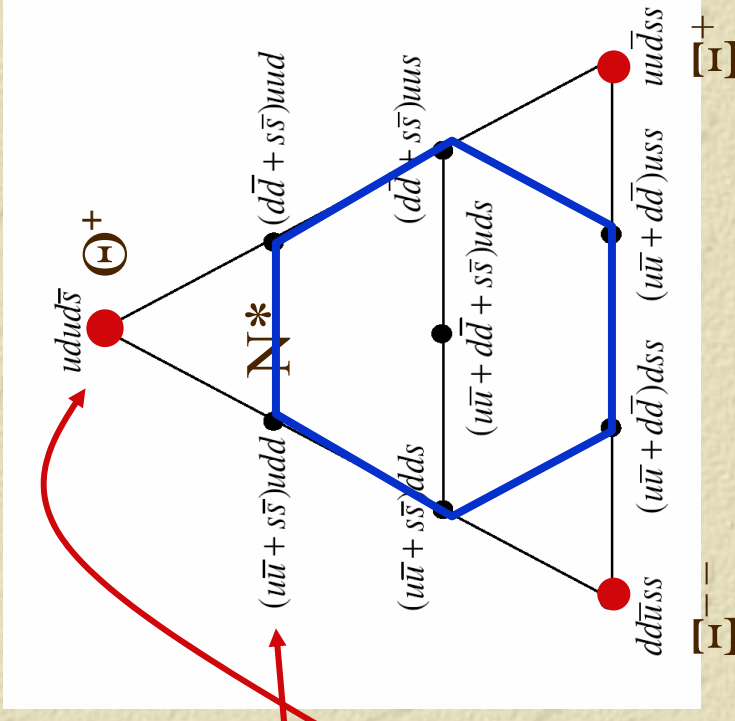
$$M = (1890 - 180 Y) \text{ MeV}$$

$$N^*(1710)$$

$$\rightarrow \Theta^+ = 1530 \text{ MeV}$$

$$\Gamma < 15 \text{ MeV}$$

10- 8 mixing expected*



Mixing of 10^* and 8

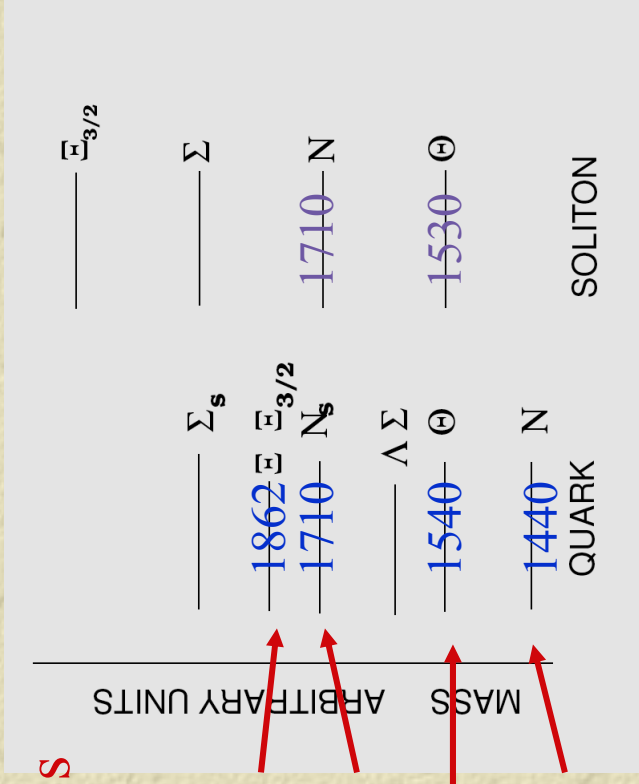


Ideal Mixing of 10^* and 8

Jaffe-Wilczek

of s

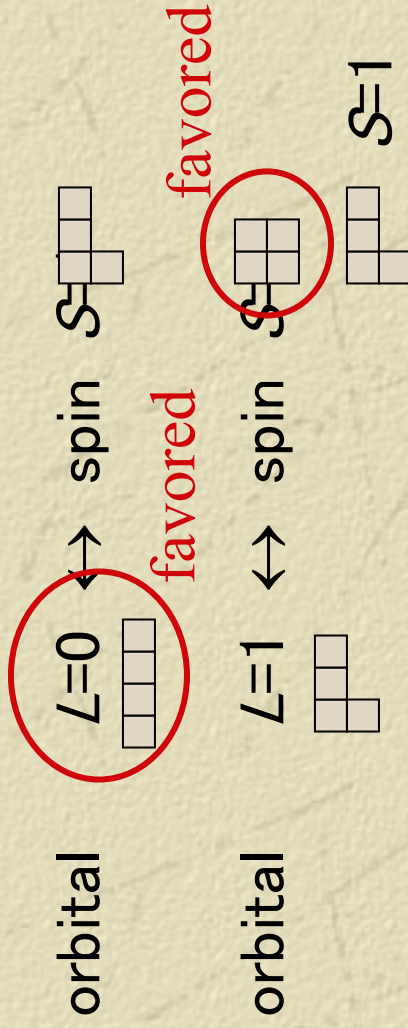
Σ, Σ_s	$udsu\bar{u}, uds\bar{s}\bar{s}$	1, 3
$\Xi, \Xi_{3/2}$	$\bar{u}ssu\bar{u}, u\bar{u}ssd$	2
$N_s (1710)$	$uud\bar{s}\bar{s} (8+10^*)$	2
Θ^+	$uudd\bar{s} 10^*$	1
$N (1440)$	$uudd\bar{d} (8+10^*)$	0



Quark Model Symmetry

 tetraquark $(q^4)_{F=10^*, C=3} \sim (u^2 d^2)_{I=0, C=3}$

$$F=10^* \begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array} \times C=3 \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \end{array} = \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \end{array} + \begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array}$$



Carlson et al. PL B573, 101 (2003)



Quark Model Symmetry

possible low energy (u^2d^2) _{$I=0$, $C=3$} states

	<i>spin-color</i>	<i>spin-flavor</i>
$L=0, S=1, J=1$ 1^+		
$L=1, S=0, J=1$ 1^-		

possible quantum numbers for Θ^+
 favored by CMI favored by FD

$$u^2d^2(1^+) + \bar{s}(1/2^-) = 1/2^-, 3/2^- \leftrightarrow \text{KN } S\text{-wave}$$

$$u^2d^2(1^-) + \bar{s}(1/2^-) = 1/2^+, 3/2^+ \leftrightarrow \text{KN } P\text{-wave}$$



Questions



Orbital excitation energy

How large is $\Delta E (\Delta L=1) = E(L=1) - E(L=0)$?



Spin-dependent interaction(s)
hyperfine (HF) interaction

How large is $\Delta E (\Delta S=1) = E(S=1) - E(S=0)$?



How large is the splitting of $J = 3/2$ and $1/2$?



PP in Quark Models

Fl. Stancu, PLB 595 (2004) 269

PS meson exchanges between quarks

The absolute mass is 510 MeV too large

Jennings, Maltman, Phys.Rev. D69 (2004) 094020

Strong CM or FD interaction favors PP

assume $\Delta E (\Delta L=1) = 210 - 250 \text{ MeV}$

Kanada-En'yo, Morimatsu, Nishikawa, hep-ph/0404144

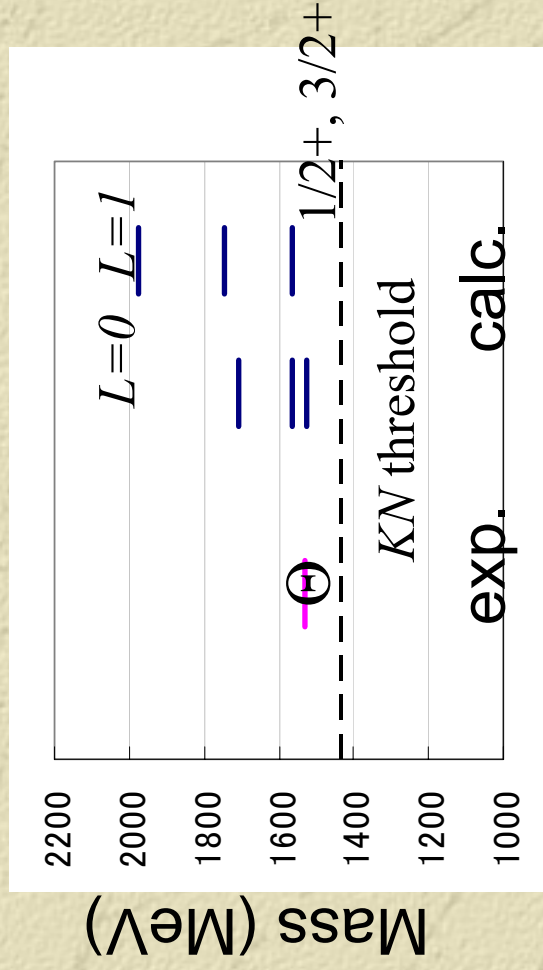
5-body molecular dynamics calculation (AMD)

Narrow positive parity states are predicted.



Diquark correlation?

 Kanada-En'yo et al.



quark and anti-quark

degeneracy of
 $1/2-, 3/2-, 1/2+,$
 $3/2+$ states

Diquarks



Diquark models Jaffe, Wilczek, PRL 91 (2003) 232003
 Karliner, Lipkin, PL B571 (2003) 55

diquarks		color	spin
$I = 1$	$uu, \{ud\}, dd$	3^*	1^+
		6	0^+
$I = 0$	$[ud]$	3^*	0^+
		6	1^+

OgE and/or III are strongly attractive in $[ud]_0 (3^*, 0^+)$

$m_{[ud]0} \sim 420 \text{ MeV}$



Diquark Models

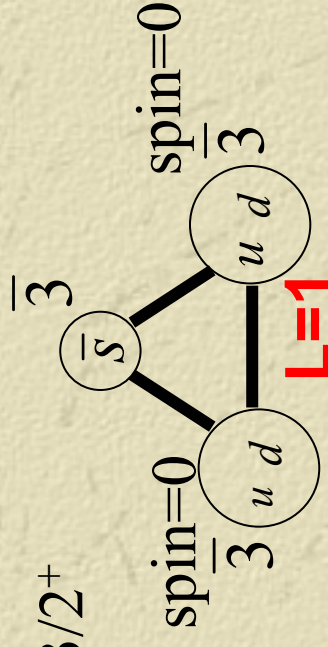
Jaffe, Wilczek

Phys. Rev. Lett. 91 (2003) 232003

2 diquarks + $\bar{s}$ model

$[ud]_0 \times [ud]_0$: color antisymmetric

$$L = 1 \Rightarrow J^\pi = 1/2^+, 3/2^+$$



Karliner, Lipkin

PL B571 (2003) 55

diquark + triquark $L = 1$

Diquark


 Diquark spectrum in lattice QCD (Landau gauge)

Wetzorke, Karsch, hep-lat/0008008

Diquark state	(FSC) [GeV]		
	($\bar{3}0\bar{3}$)	($6\bar{1}\bar{3}$)	($\bar{3}16$) (606)
$ma_{(FSC)}: \boxed{\text{MEM}}$ result	0.60(2)	0.70(3)	0.74(9) —
$ma_{(FSC)}: 2\text{-exp. fit}$	0.62(2)	0.73(4)	0.77(17) 0.50(15)

maxial entropy method

Are they hadronic?

✠ Hadronic bound/resonances

KN (K^+n , K^0p) $I = 0$, $L = 0$ interaction is weak and repulsive

K^+n **weak repulsive**

K^-p strong attractive

NK π bound state 30 MeV below $E_{\text{th}} = 1.57$ GeV

Width must be strongly suppressed for $L=0$ bound state.

Bicudo, Marques, PRD69(2004) 011503

did not find a bound state.



Conclusion

- ✠ QCD is consistent with a $J^\pi = 1/2^-$ $F=10^*$ pentaquarks.
But the width may not be narrow.
- ✠ Many models suggest positive parity baryons.
But the $1/2^+$ state has a larger mass in the quark model.
- ✠ Width
Need to explain the unusually narrow width.
- ✠ We need more experimental data
 - ◆ spin, parity and width
 - ◆ other 10^* members
 - ◆ excited states ex. $J=3/2$ states, $I=1$ states?



More theoretical ideas