

Collective modes investigated by inelastic scattering and charge-exchange reactions with magnetic spectrometers

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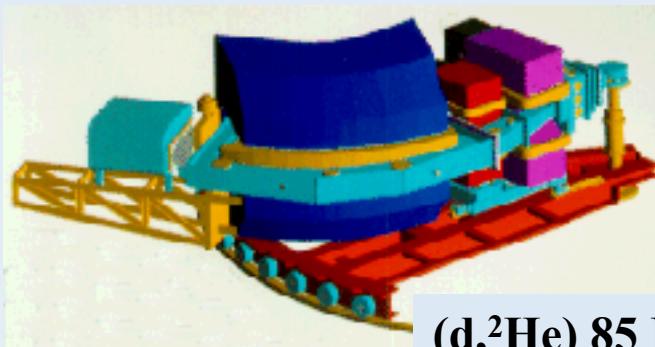
Kernfysisch Versneller Instituut

4th CNS International Summer School (CISS05)

Outline

- **Introduction on Giant Resonances (GRs)**
What are GRs? Why study GRs?
- **Isoscalar giant resonances; compression modes**
Inelastic α -scattering \Rightarrow ISGMR, ISGDR
- **Isovector charge-exchange modes**
($^3\text{He},t$) & ($d,^2\text{He}$) charge-exchange reactions
 \Rightarrow GTR, IVSGMR, IVSGDR; GT⁺ Strength
- **Outlook**



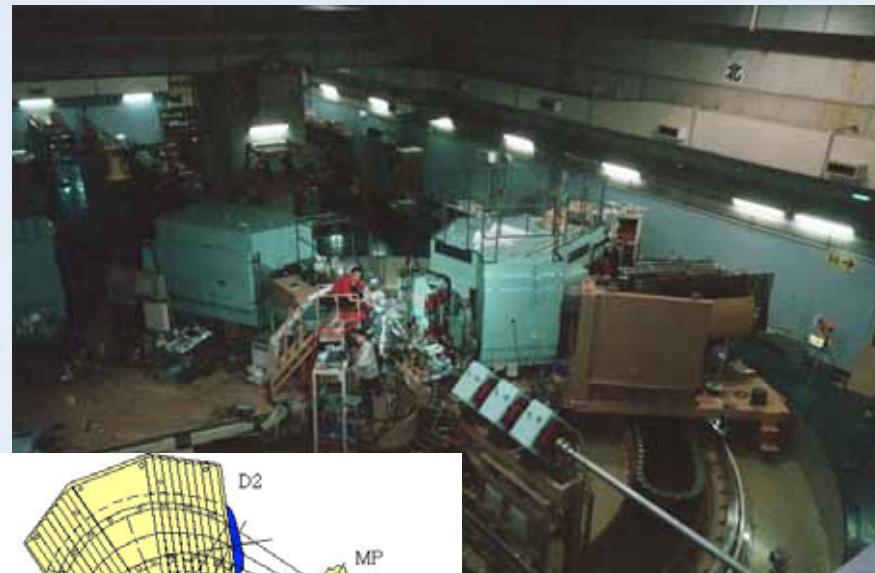


BBS@KVI

(d,²He) 85 MeV/u
(³He,t) 60 MeV/u
(t,³He) 43 MeV/u

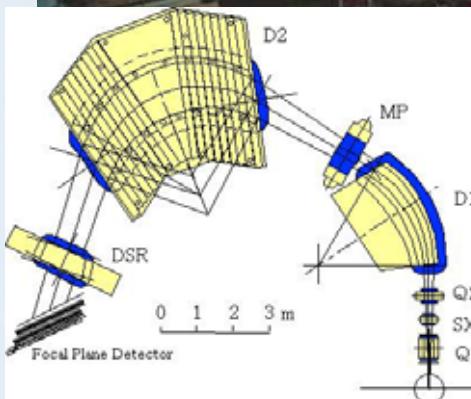


S800@NSCL ($t,{}^3\text{He}$) 115 MeV/u



Grand Raiden@RCNP

(³He,t) 140 MeV/u
Also done
(p,n) (n,p) 300 MeV



TRIUMF (n,p) (π -CEX)
IUCF (p,n) (³He,t) K600
iThemba (p,p) K600
Texas A&M
...etc



SHARAQ *Spectroscopy of HAdron systems with RadioActive Quantum beams*

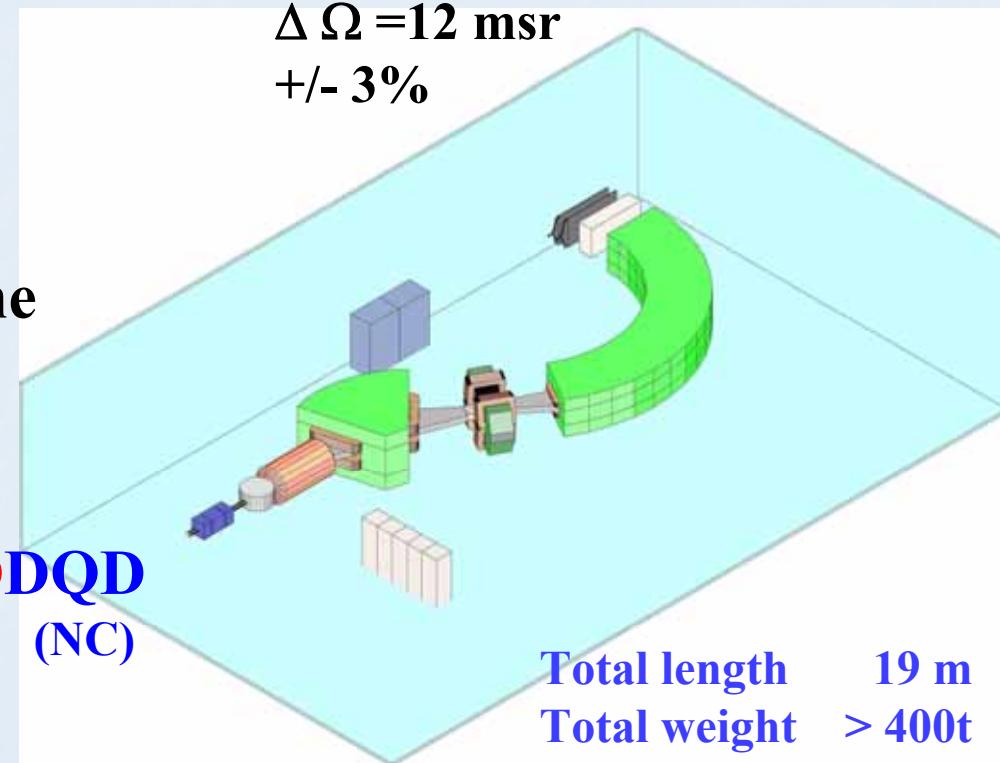


Parameters

- Momentum resolution $\Delta p/p < 1/15000$
 $\delta E_x \sim 300 \text{ keV}$ for $A=8, E=300\text{MeV/u}$
- Maximum rigidity $B\rho = 6.8\text{Tm}$ ($\rho = 4.8\text{m}$)
- Solid angle $\Delta \Omega = 12 \text{ msr}$
- Momentum acceptance $\pm 3\%$

High-quality RI beam-line

- Dispersion matching



Giant resonances: fundamental high-frequency modes of nuclear excitation

- Interesting to study on own merit: E_x , Γ , %EWSR;
Microscopic structure \Rightarrow compare to theory:
HF+RPA; RMFT, RRPA, etc.
- Effective interactions, medium effects (meson-coupling constants); IVSGDR 0⁻ strength (pion quantum numbers)

Moreover,

- Equation of state (EOS)
ISGMR, ISGDR \Rightarrow Incompressibility, symmetry energy

$$K_A = K_{\text{vol}} + K_{\text{surf}} A^{-1/3} + K_{\text{sym}} ((N-Z)/A)^2 + K_{\text{Coul}} Z^2 A^{-4/3}$$



- IVGDR, IVSGDR, IVGMR \Rightarrow n-skin thickness; isospin mixing
- Collective-flow in H.I. collisions
- Astrophysics: Input for supernova explosions, neutron stars

Microscopic picture: GRs are coherent (1p-1h) excitations induced by single-particle operators.

- Excitation energy depends on i) multipole L ($L\hbar\omega$, except for isoscalar monopole and isoscalar dipole GRs, $2\hbar\omega$ & $3\hbar\omega$, respectively), ii) strength of effective interaction and iii) collectivity.
- Exhaust appreciable % of EWSR and/or NEWSR
- Acquire a width due to coupling to continuum and to underlying 2p-2h configurations.



Nucleus → **Many-body system with a finite size**

Vibrations → **Multipole expansion with r, Y_{lm}, τ, σ**

$\Delta S=0, \Delta T=0$ $\Delta S=0, \Delta T=1$ $\Delta S=0, \Delta T=1$ $\Delta S=1, \Delta T=1$ $\Delta S=1, \Delta T=1$

L=0: Monopole	ISGMR $r^2 Y_0$	IAS τY_0	IVGMR $\tau r^2 Y_0$	GTR $\tau \sigma Y_0$	IVSGMR $\tau \sigma r^2 Y_0$
L=1: Dipole	ISGDR $r^3 Y_1$		IVGDR $\tau r Y_1$		IVSGDR $\tau \sigma r Y_1$
L=2: Quadrupole	ISGQR $r^2 Y_2$		IVGQR $\tau r^2 Y_2$		IVSGQR $\tau \sigma r^2 Y_2$
L=3: Octupole	LEOR, HEOR $r^3 Y_3$				



Isoscalar Excitation Modes of Nuclei

Giant Resonance: Coherent vibration of nucleons in a nucleus.

Compression modes : **ISGMR, ISGDR**

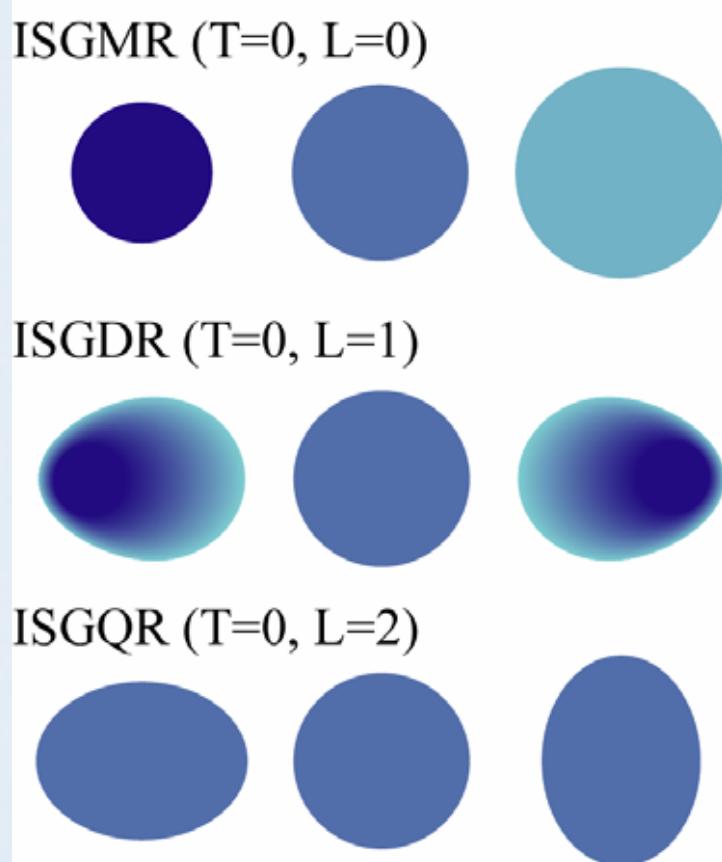
$$E_{ISGMR} = \hbar \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

$$E_{ISGDR} = \hbar \sqrt{\frac{7}{3} \frac{K_A + \frac{27}{25} \epsilon_F}{m \langle r^2 \rangle}}$$

The nucleus incompressibility:



$$K_A = \left[r^2 \left(\frac{d^2(E/A)}{dr^2} \right) \right]_{r=R_0}$$



In HF+RPA calculations,

$$K_{nm} = \left[9\rho^2 \frac{d^2(E/A)}{d\rho^2} \right]_{\rho=\rho_0}$$

E/A: binding energy per nucleon

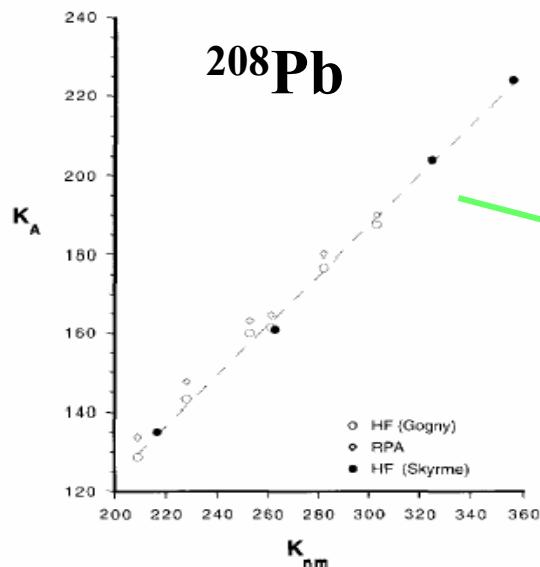
ρ : nuclear density
energy

ρ_0 : nuclear density at saturation

Nuclear matter

K_A : incompressibility

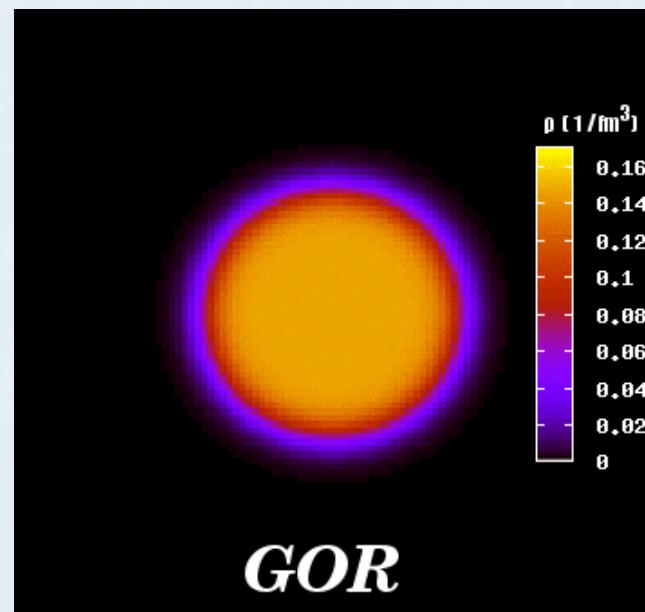
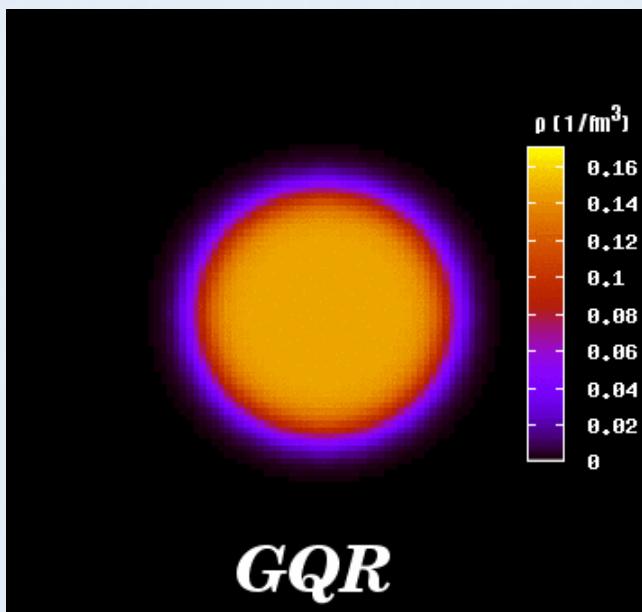
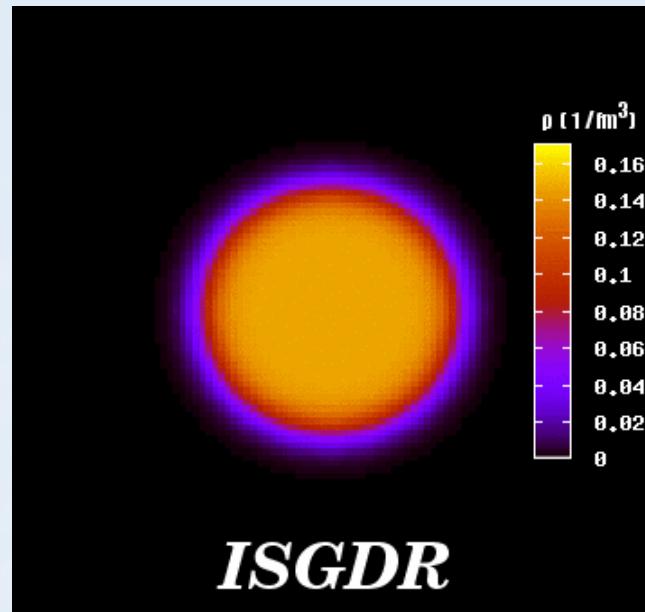
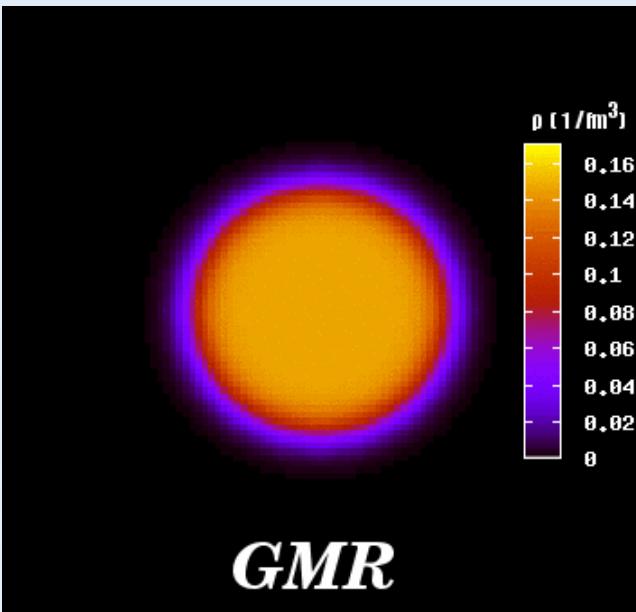
ε_F : Fermi



K_A is obtained from excitation energy of ISGMR & ISGDR

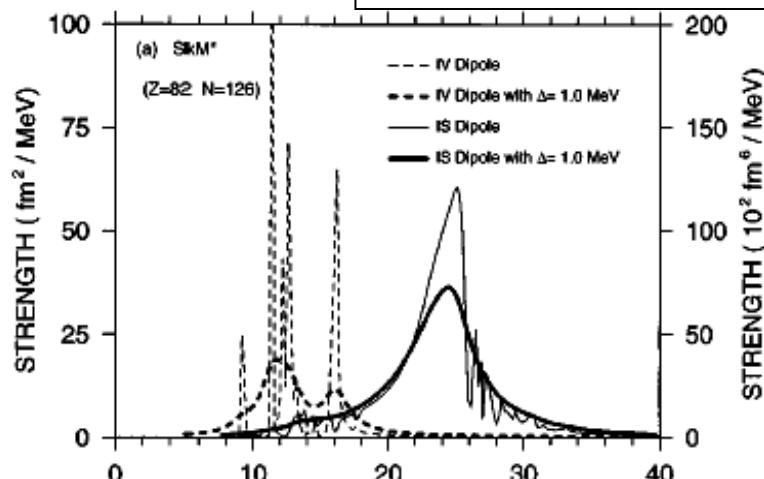
$$K_A = 0.64K_{nm} - 3.5$$

J.P. Blaizot, NPA591 (1995) 435

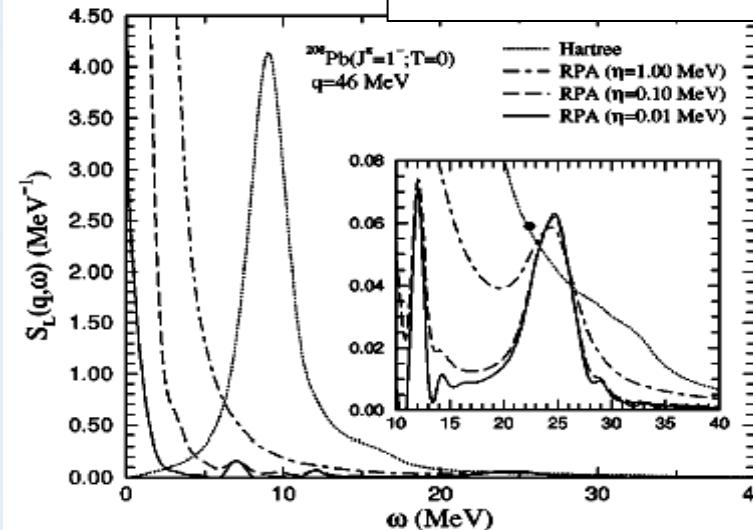


HF+RPA; ISGDR in ^{208}Pb

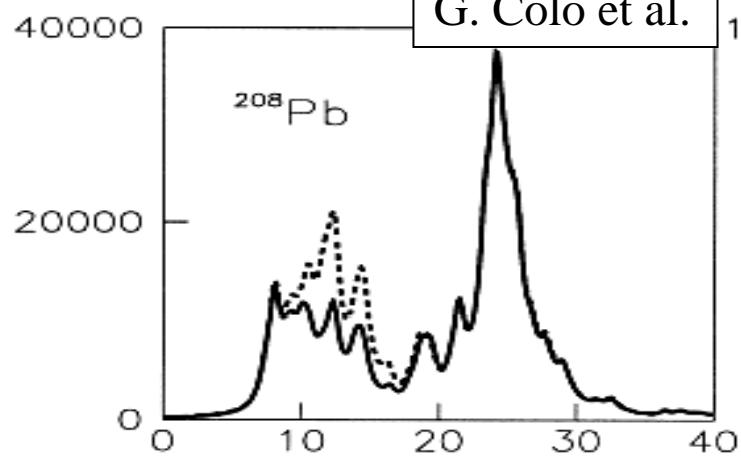
I. Hamamoto et al.



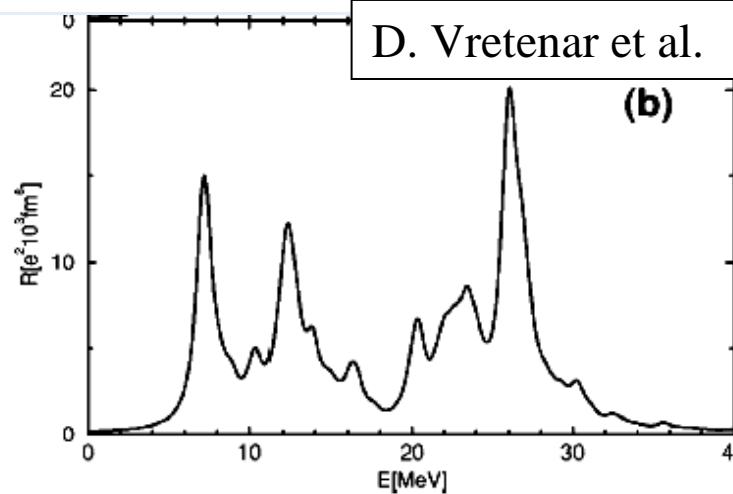
J. Pikarewicz



G. Colo et al.

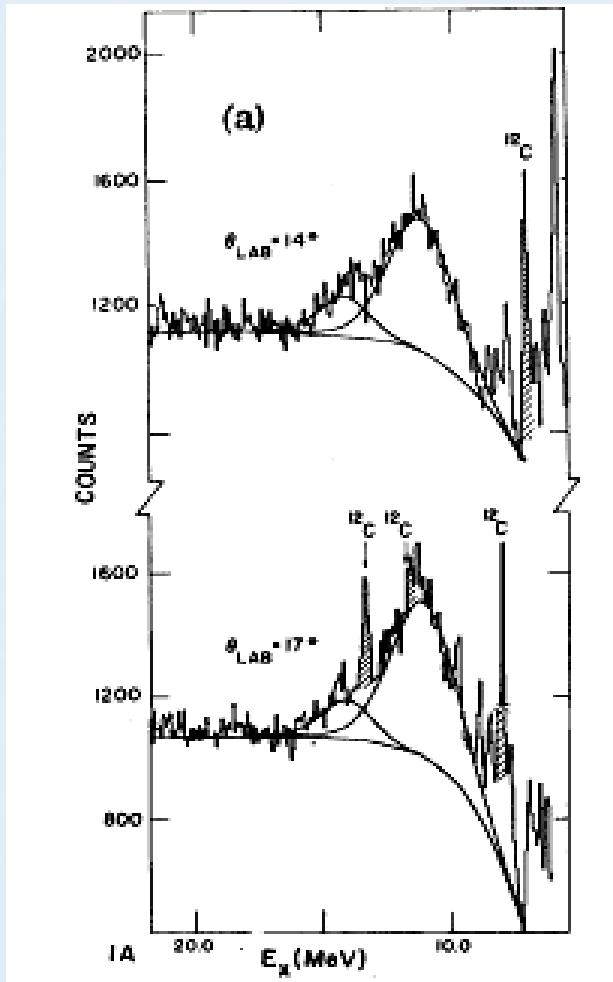


D. Vretenar et al.

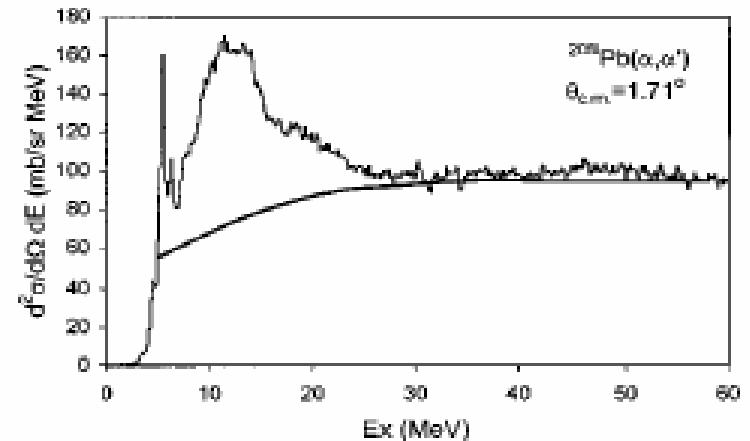


ISGQR, ISGMR, ISGDR

KVI (1977)



TAMU(2000)

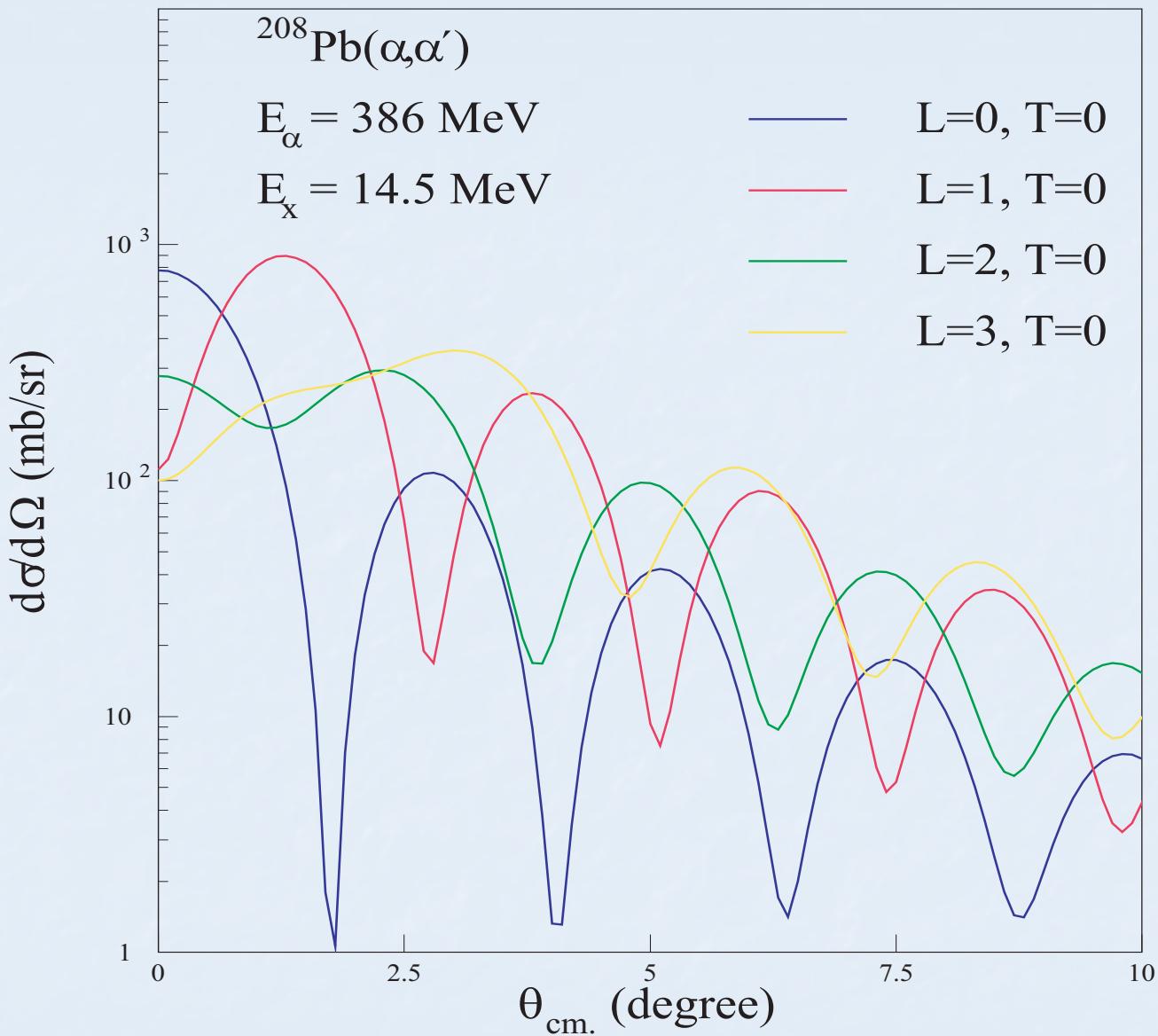


$E_\alpha = 240$ MeV

Large instrumental background!

$\Leftarrow ^{208}\text{Pb}(\alpha, \alpha')$ at $E_\alpha = 120$ MeV





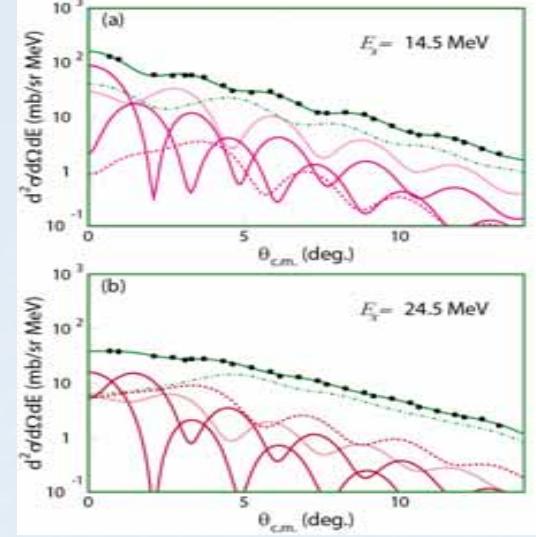
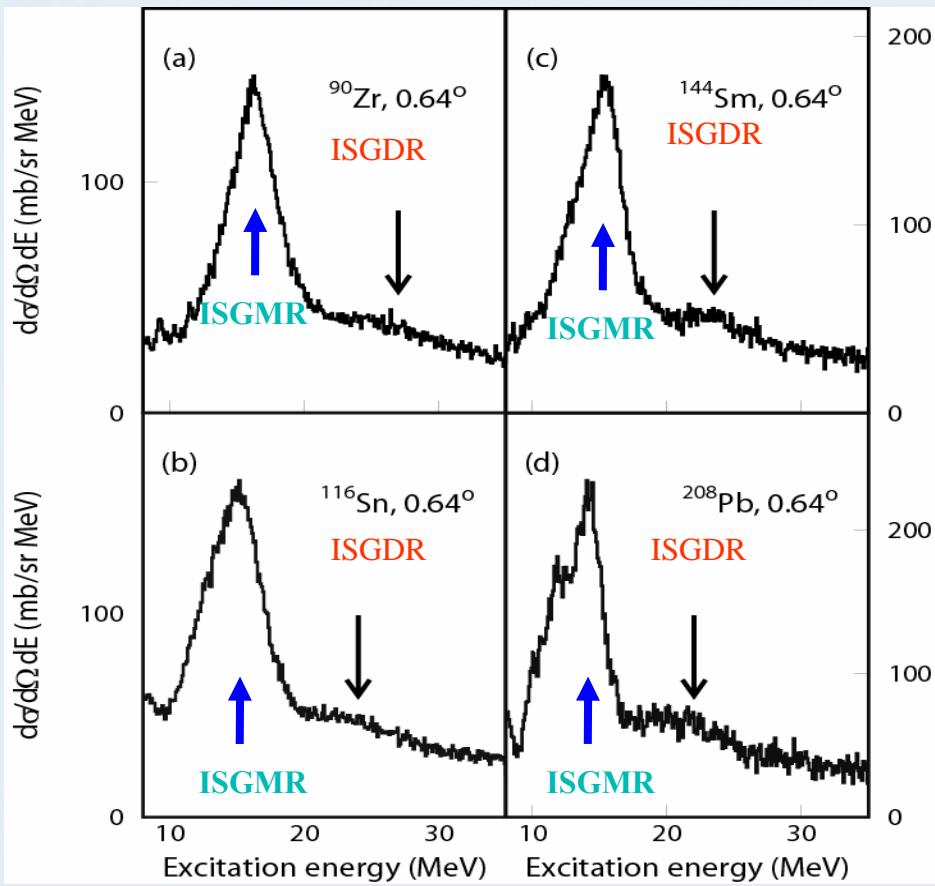
ISGMR, ISGDR

ISGQR, HEOR

100 % EWSR

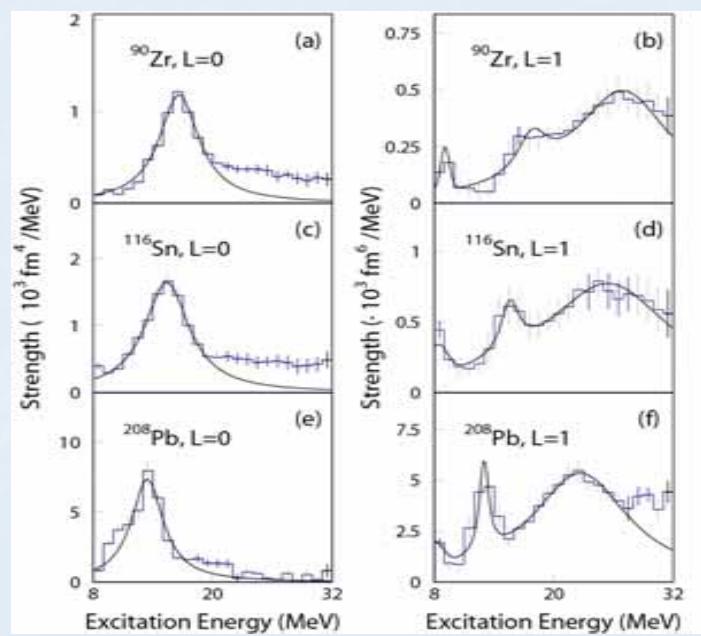
At $E_x = 14.5 \text{ MeV}$

(α, α') spectra at 386 MeV



^{116}Sn

MDA results for L=0 and L=1



Multipole decomposition analysis (MDA)

$$\left(\frac{d^2\sigma}{d\Omega dE}(\mathcal{G}_{c.m.}, E) \right)^{ex.} = \sum_L a_L(E) \left(\frac{d^2\sigma}{d\Omega dE}(\mathcal{G}_{c.m.}, E) \right)_L^{calc}$$

$\left(\frac{d^2\sigma}{d\Omega dE}(\mathcal{G}_{c.m.}, E) \right)^{ex.}$: Experimental cross section

$\left(\frac{d^2\sigma}{d\Omega dE}(\mathcal{G}_{c.m.}, E) \right)_L^{calc}$: DWBA cross section (unit cross section)

$a_L(E)$: EWSR fraction

- ISGR ($L < 15$) + IVGDR (through Coulomb excitation)
- DWBA formalism; single folding \Rightarrow transition potential

$$\delta U(r, E) = \int d\vec{r}' \delta\rho_L(\vec{r}', E) [V(|\vec{r} - \vec{r}'|, \rho_0(r')) + \rho_0(r') \frac{\partial V(|\vec{r} - \vec{r}'|, \rho(r'))}{\partial \rho_0(r')}]$$

$$U(r) = \int d\vec{r}' V(|\vec{r} - \vec{r}'|, \rho_0(r')) \rho_0(r')$$

Transition density

- ISGMR Satchler, Nucl, Phys, A472 (1987) 215

$$\delta\rho_0(r, E) = -\alpha_0 [3 + r \frac{d}{dr}] \rho_0(r)$$

$$\alpha_0^2 = \frac{2\pi}{mA} \frac{\hbar^2}{\langle r^2 \rangle / E}$$

- ISGDR Harakeh & Dieperink, Phys. Rev. C23 (1981) 2329

$$\delta\rho_1(r, E) = -\frac{\beta_1}{R\sqrt{3}} [3r^2 + \frac{d}{dr} + 10r - \frac{5}{3}\langle r^2 \rangle \frac{d}{dr} + \varepsilon(r \frac{d^2}{dr^2} + 4 \frac{d}{dr})] \rho_0(r)$$

$$\beta_1^2 = \frac{6\pi l^2}{mA E} \frac{R^2}{(11\langle r^4 \rangle - (25/3)\langle r^2 \rangle^2 - 10\varepsilon\langle r^2 \rangle)}$$

- Other modes Bohr-Mottelson (BM) model

$$\delta\rho_L(r, E) = -\delta_L \frac{d}{dr} \rho_0(r)$$

$$\delta_L^2 = (\beta_L c)^2 = \frac{l(2l+1)^2}{(l+2)^2} \frac{2\pi}{mA E} \frac{\hbar^2}{\langle r^{l-2} \rangle^2}$$

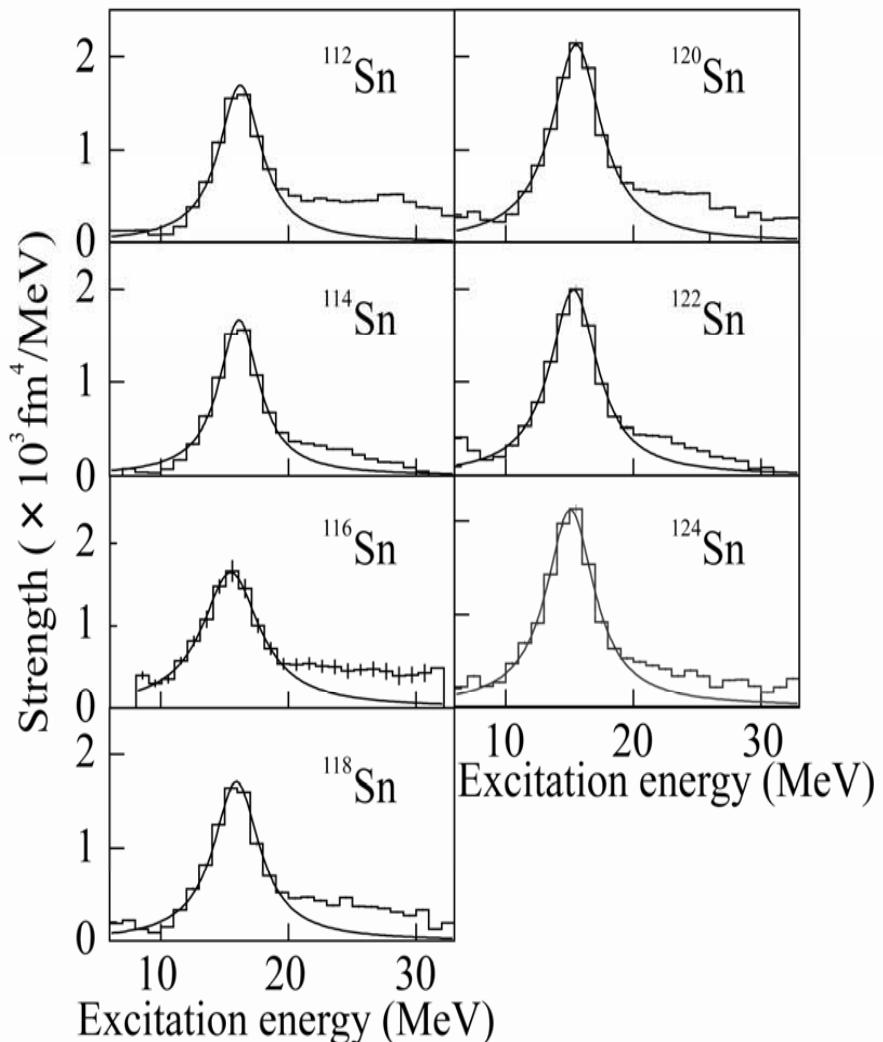
Results of ISGMR/ISGDR

	E _{GMR} (MeV)	Γ (MeV)	EWSR (%)	E _{LEGDR} (MeV)	Γ (MeV)	EWSR (%)	E _{HEGDR} (MeV)	Γ (MeV)	EWSR (%)
⁹⁰ Zr	16.6 (0.1)	4.9 (0.2)	101 (3)	17.8 (0.5)	3.7 (1.2)	7.9 (2.9)	26.9 (0.7)	12.0 (1.5)	67 (8)
¹¹⁶ Sn	15.4 (0.1)	5.5 (0.3)	95 (4)	15.6 (0.5)	2.3 (1.)	4.9 (2.2)	25.4 (0.5)	15.7 (2.3)	68 (9)
¹⁴⁴ Sm	15.3 (0.1)	3.71 (0.12)	84 (4,-25)	14.2 (0.2)	4.8 (0.8)	23 (4,-10)	25.0 (1.7)	19.9 (1.4)	91 (25,-17)
²⁰⁸ Pb	13.5 (0.2)	4.0 (0.4)	104 (9)	13.0 (0.1)	1.1 (0.4)	7.0 (0.4)	22.8 (0.3)	11.9 (0.4)	111 (6)



ΔE is about 0.1 MeV

$L=0$



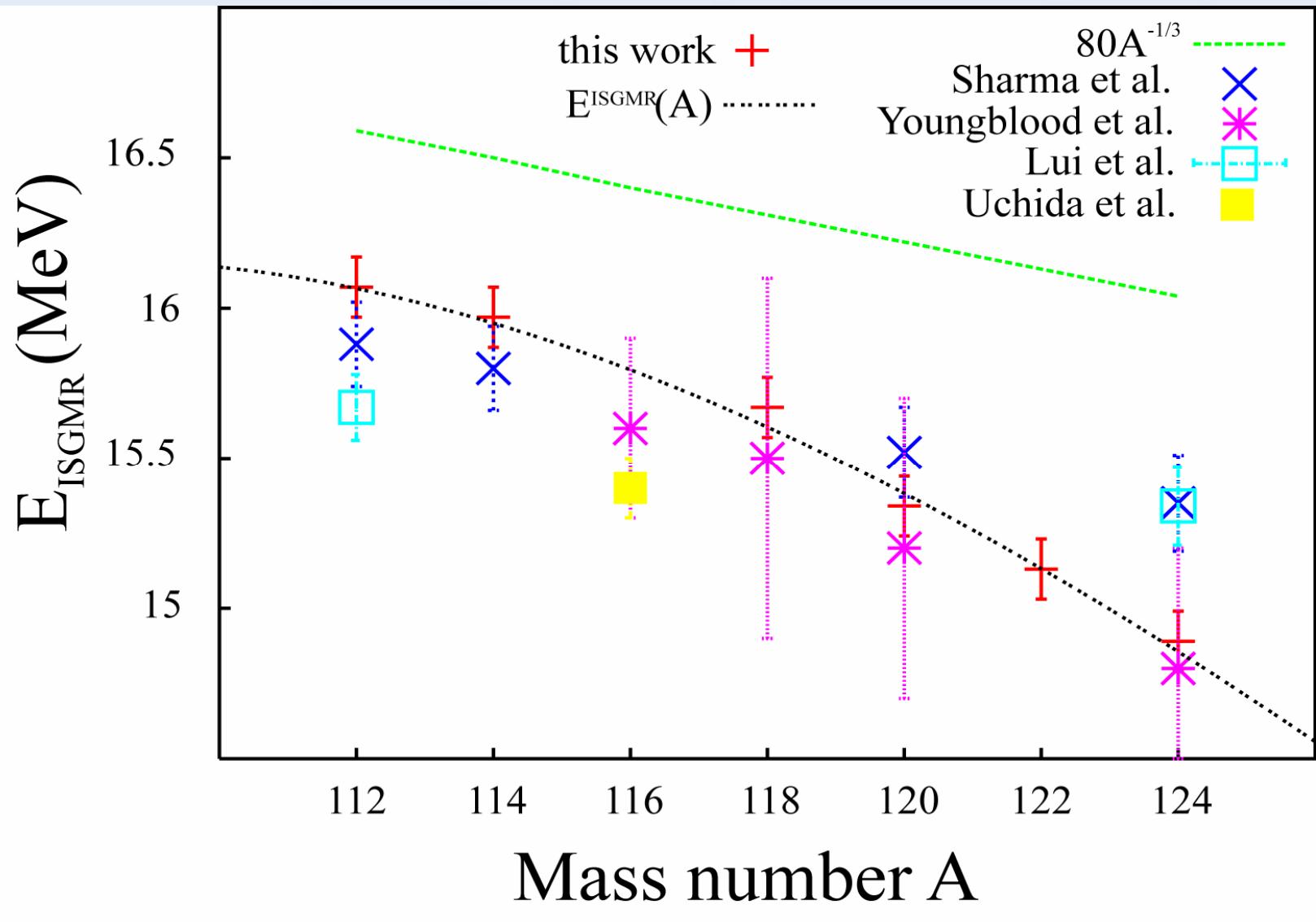
Breit-Wigner shape:

	$E_m(\text{MeV})$	Γ_m
^{112}Sn	16.1 ± 0.1	4.4 ± 0.2
^{114}Sn	16.0 ± 0.1	4.0 ± 0.1
^{116}Sn	15.4 ± 0.1	5.7 ± 0.3
^{118}Sn	15.7 ± 0.1	4.6 ± 0.2
^{120}Sn	15.3 ± 0.1	4.8 ± 0.2
^{122}Sn	15.1 ± 0.1	4.5 ± 0.2
^{124}Sn	14.9 ± 0.1	4.5 ± 0.3

$^{112,114,118,120,122,124}\text{Sn}$: this work

^{116}Sn : Uchida et al.

ISGMR energy E_{ISGMR}



208Pb

	E_{ISGDR} (MeV)	E_{ISGMR} (MeV)	K_{nm} (MeV)
	HE	LE	
Uchida et al. (Breit-Wigner)	22.8 (0.3)	13.0 (0.2)	13.5 (0.2)
Morsch et al.	21.3 (0.8)		13.8
Djalali et al.	21.5 (0.2)		13.9
Adams et al.	22.6 (0.2)		
Davis et al.	22.4(0.5)		
Clark et al.	19.9(0.8)	12.2(0.6)	14.17(0.28)
Hamamoto et al.	23.4	~14	14.1
Colo et al.	23.9(22.9 *)	10.9	14.1
Piekarewicz et al.	24.4	~8	13.1
Vretenar et al.	26.01	10.4	14.1
Shlomo and Sanzhur	~25.0	~15	14.48
Gorelik and Urin	22.7	11.1	14.3

*** Including the effect of 2p-2h coupling**

Decay of giant resonances

- Width of resonance

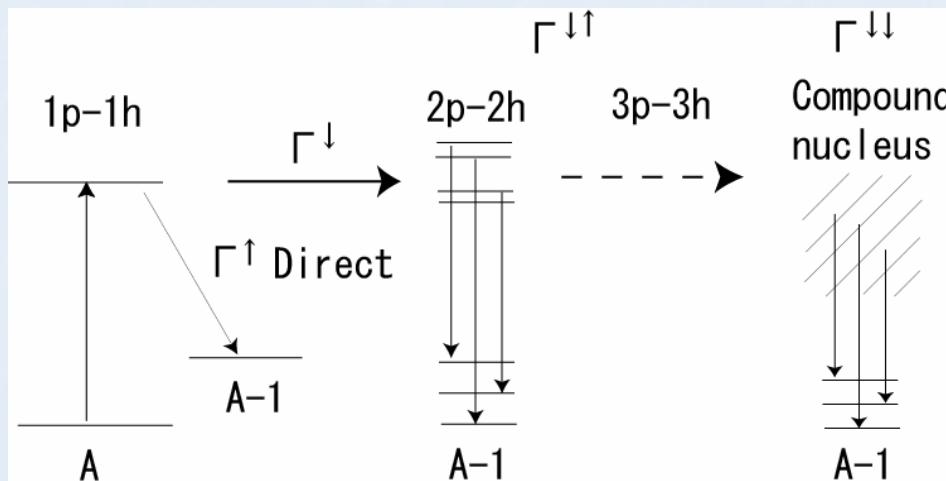
$$\Gamma, \Gamma^{\uparrow}, \Gamma^{\downarrow} (\Gamma^{\downarrow\uparrow}, \Gamma^{\downarrow\downarrow})$$

- Γ^{\uparrow} : direct or escape width
- Γ^{\downarrow} : spreading width
- $\Gamma^{\downarrow\uparrow}$: pre-equilibrium, $\Gamma^{\downarrow\downarrow}$: compound

- Decay measurements

⇒ Direct reflection of damping processes

Allows detailed comparison with theoretical calculations



Decay and microscopic structure of ISGDR

Transition operator

$$O^{L=1} = \sum_i r_i Y_0^1 + \sum_i r_i^3 Y_0^1 + \dots$$

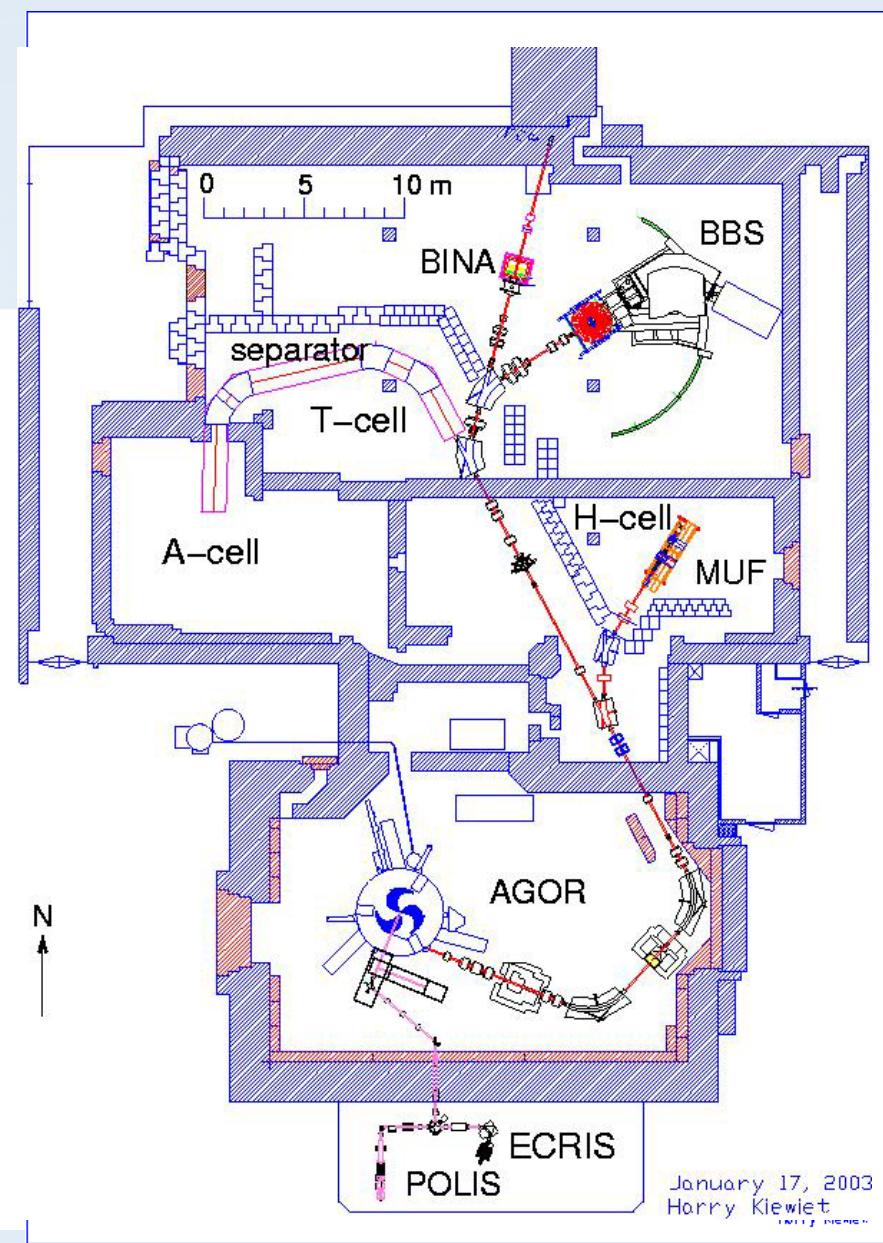
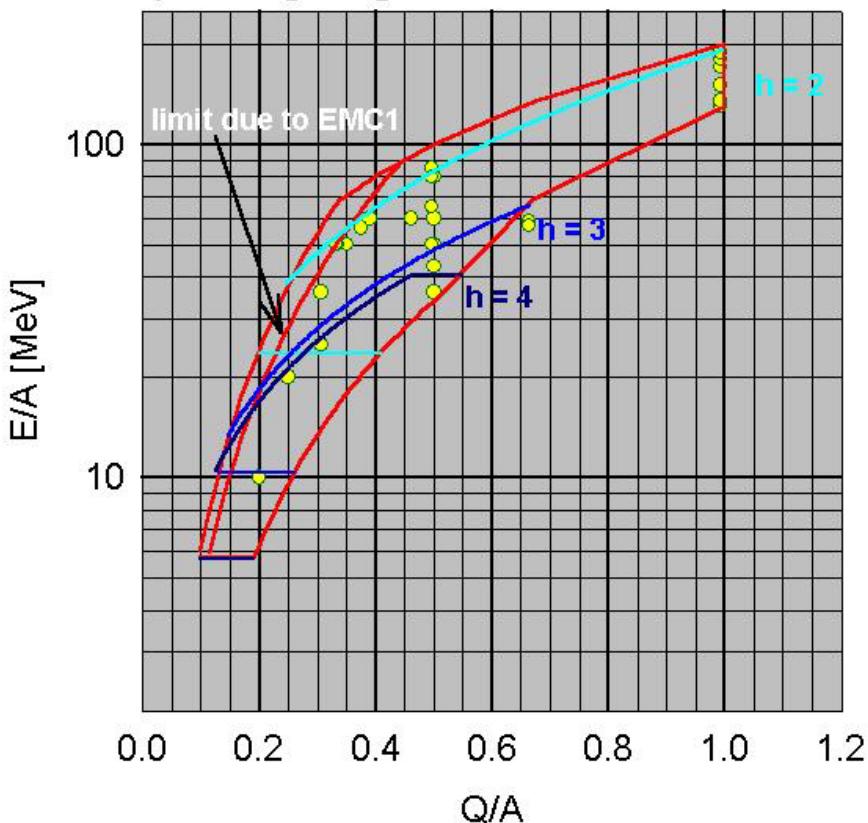
**spurious center
of mass motion**

overtone

$3\hbar\omega$ excitation (overtone of c.o.m. motion)

AGOR cyclotron K = 600
 light and heavy ions
 $p < 190$ MeV
 $q/A = 1/2$; $E/A < 90$ MeV

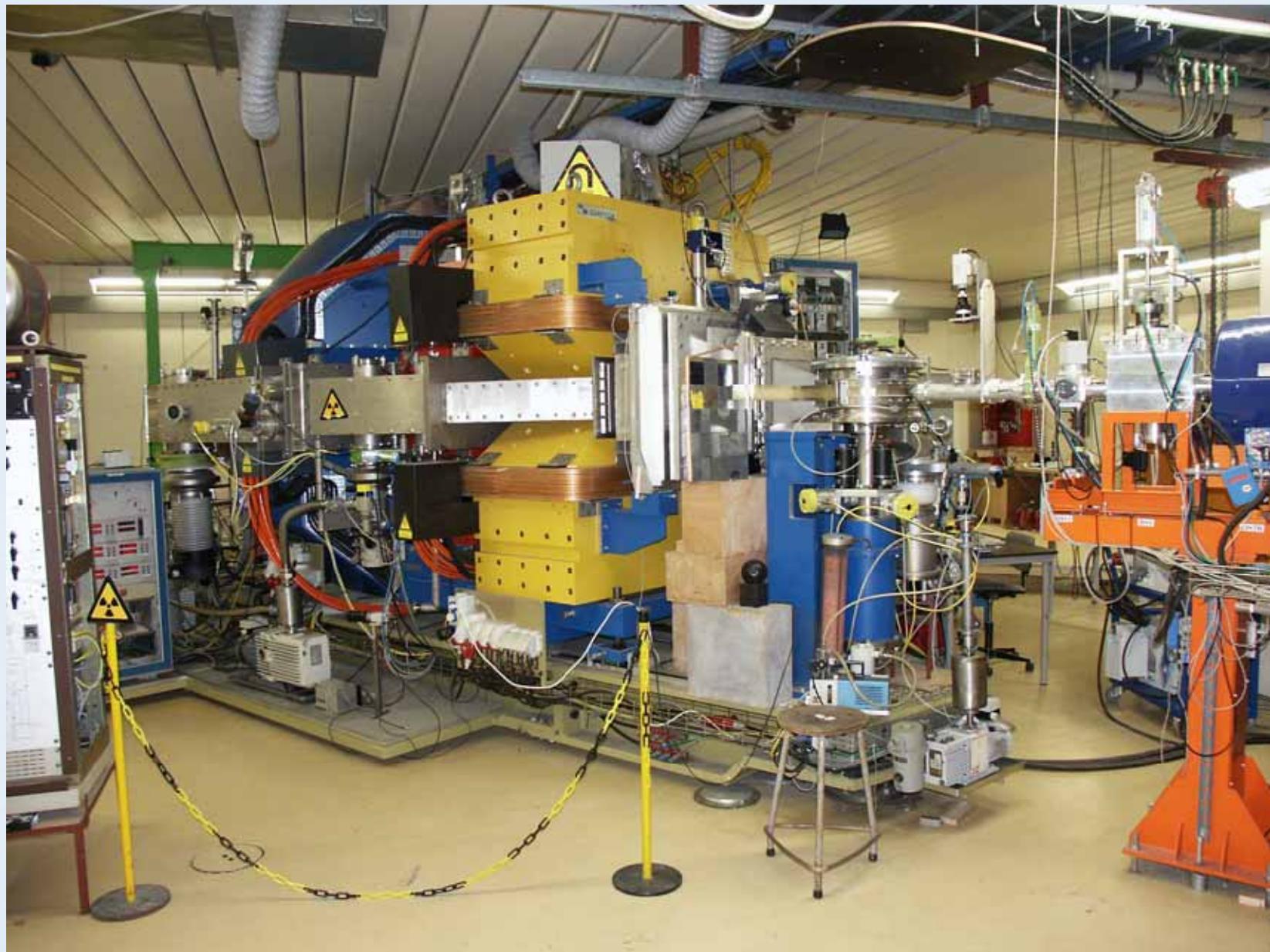
operating diagram and available beams



Equipment

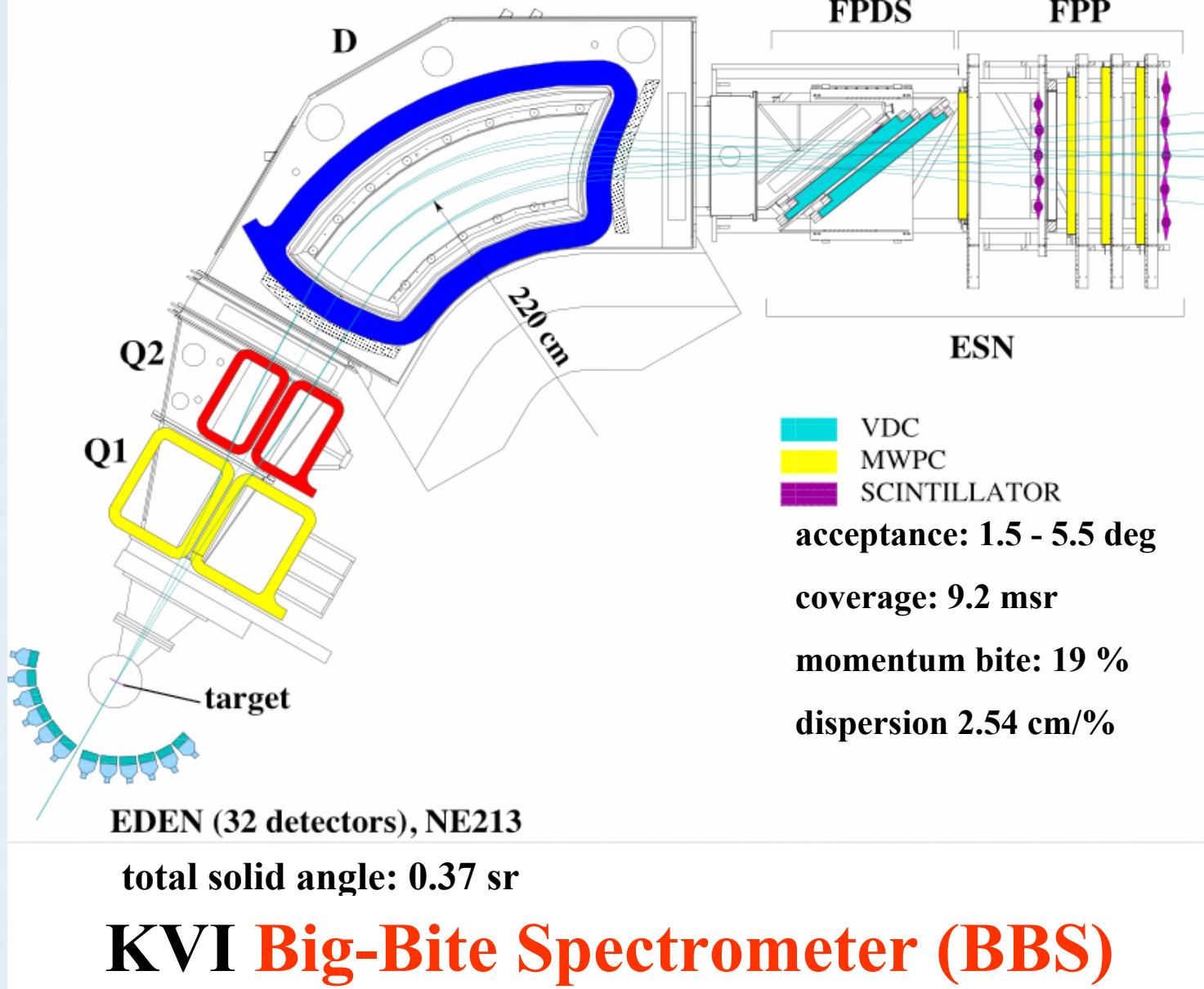
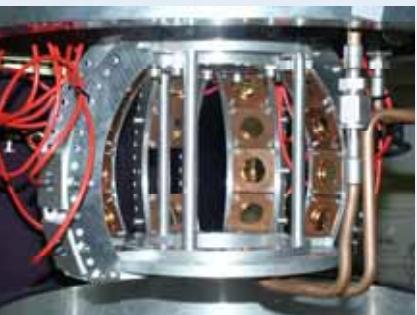
- Big momentum Bite magnetic Spectrometer (BBS)
- its associated Focal-Plane Detectors (ESN)
- various coincidence detectors
 - neutrons (EDEN), protons (SiLi-ball), γ 's (GeLi, Clover), phoswich detectors
- readout electronics
 - DSP, VME, FERA



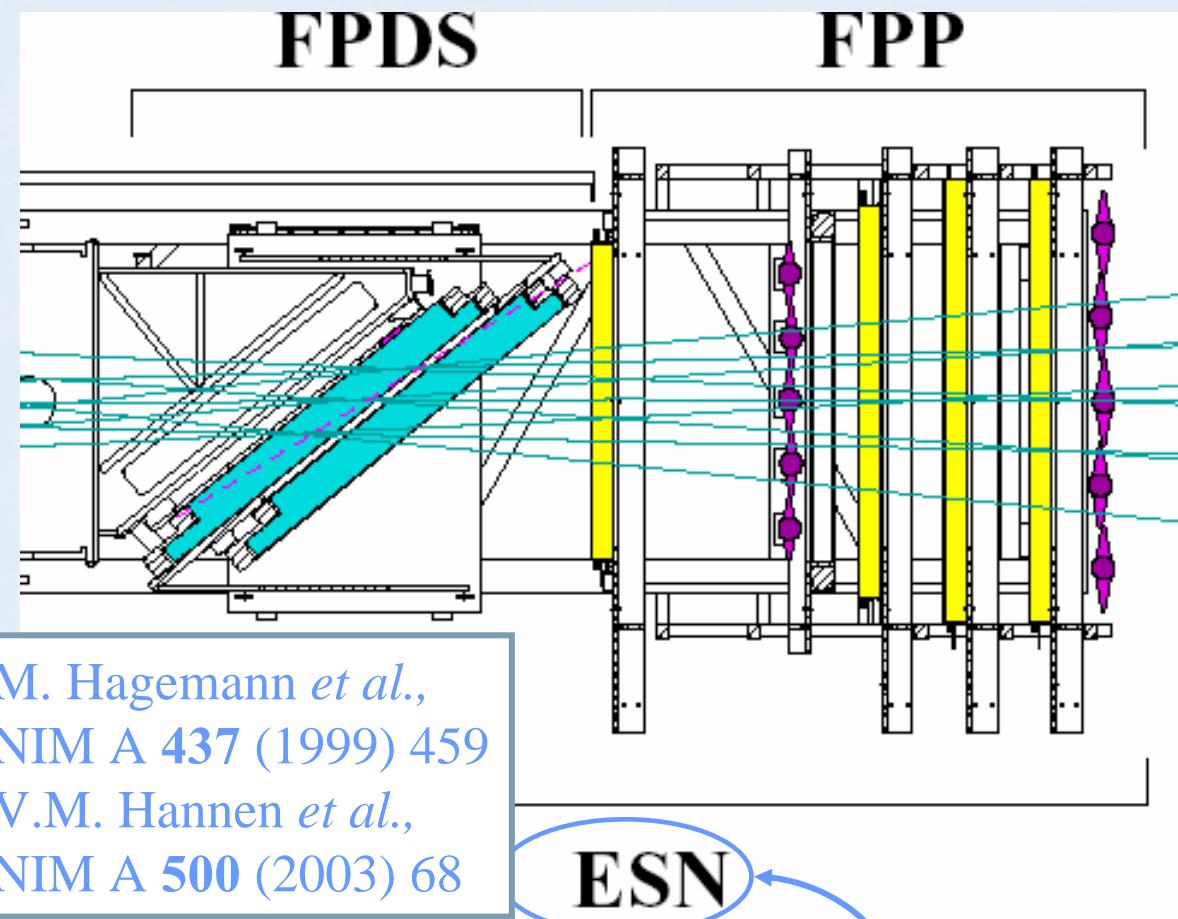


CISS05, 18-23 August 2005; Wako, Tokyo, Japan

Si-ball
16 Si-detectors at
10 cm from the target
total solid angle: 1 sr



Setup: ESN detector



Focal-Plane Detector:
(FPDS): 2 VDCs

Focal-Plane Polarimeter:
(FPP): 4 MWPCs &
graphite analyzer

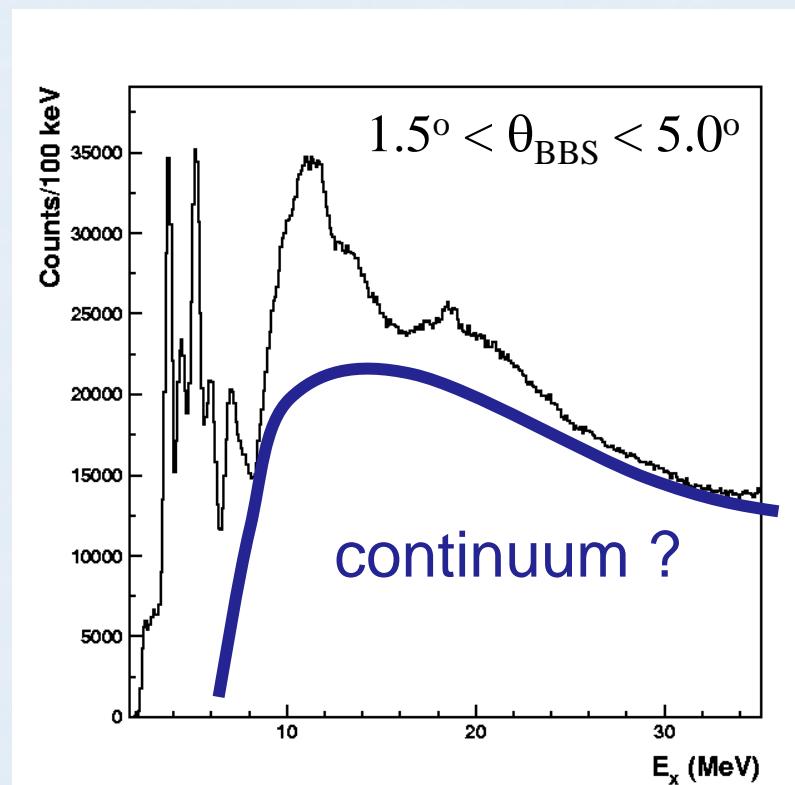
features a.o.:
fast readout
VDC readout **pipeline TDCs**
VDC decoding using
imaging techniques
DSP based online analysis

Bari, Darmstadt, Gent, Iserlohn, KVI, Milano, Münster, TRIUMF

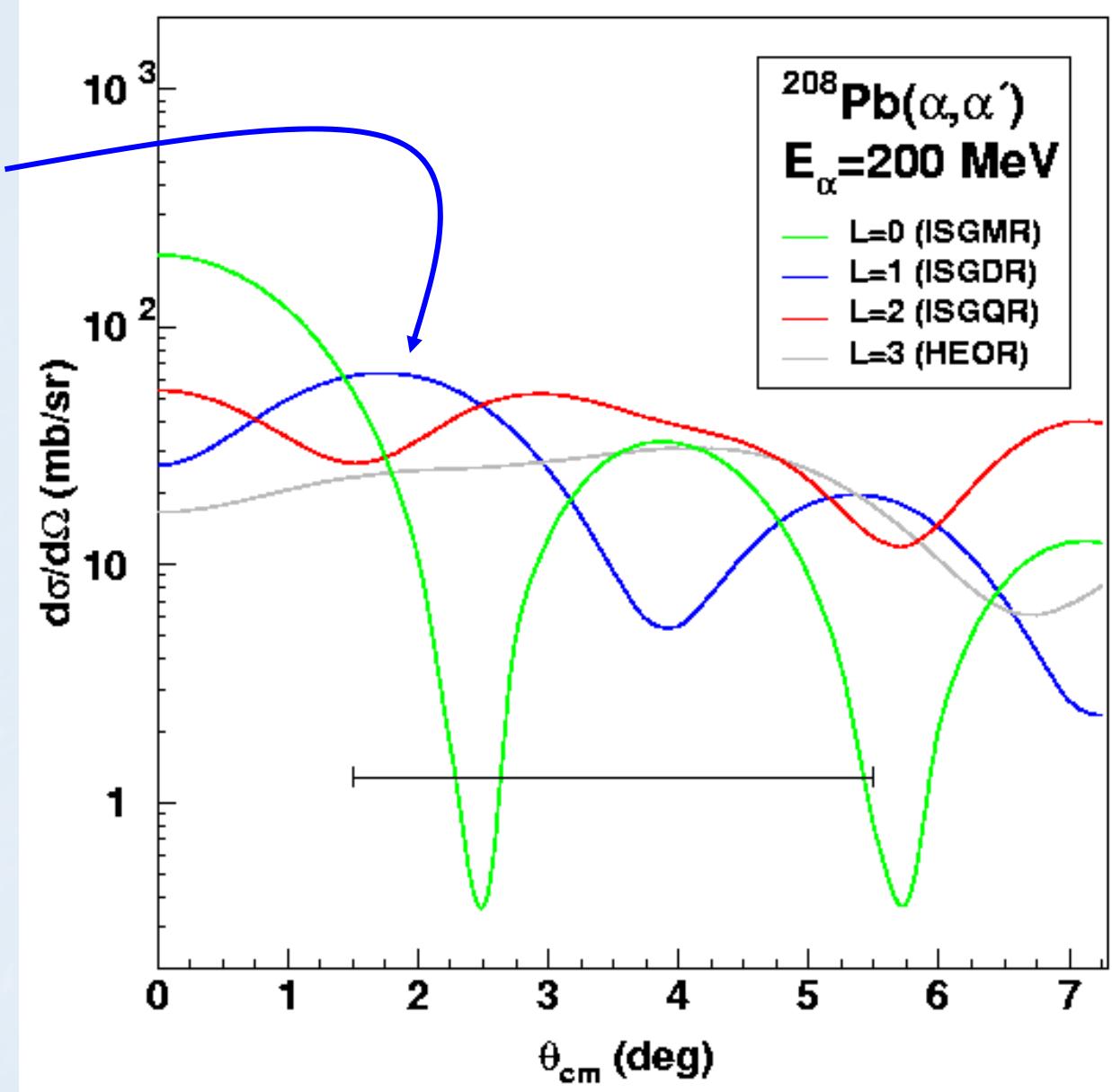
Excitation of ISGDR in ^{208}Pb

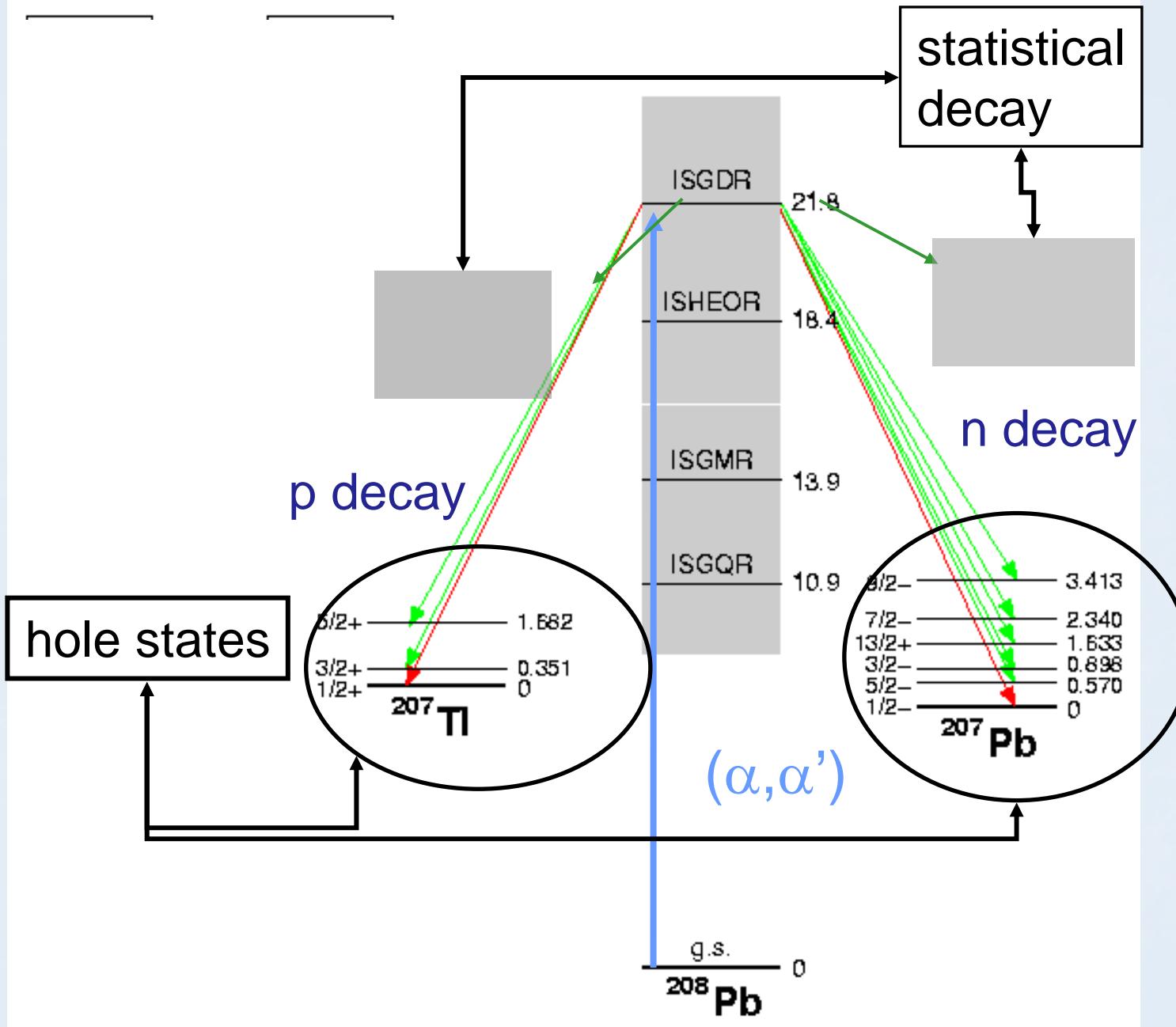
- In ^{208}Pb located around 22 MeV and width of 4 MeV
- $L=1$ angular distribution peaks close to a scattering angle of 0°
- Difficult to identify in nuclear continuum and rides on instrumental background

Singles $^{208}\text{Pb}(\alpha, \alpha') 200 \text{ MeV} \Rightarrow$

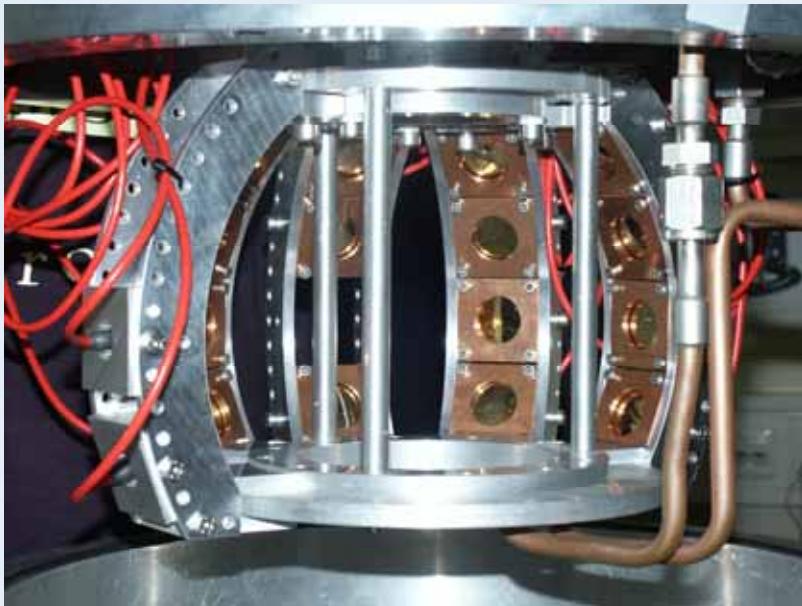


ISGDR L = 1

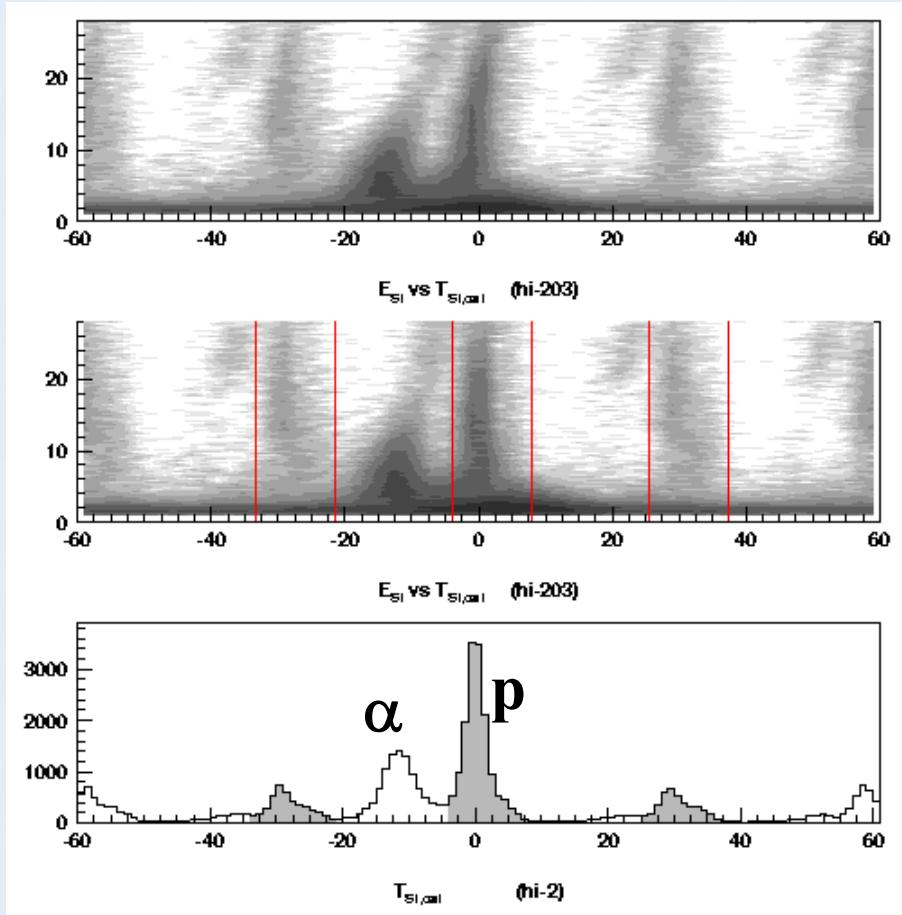




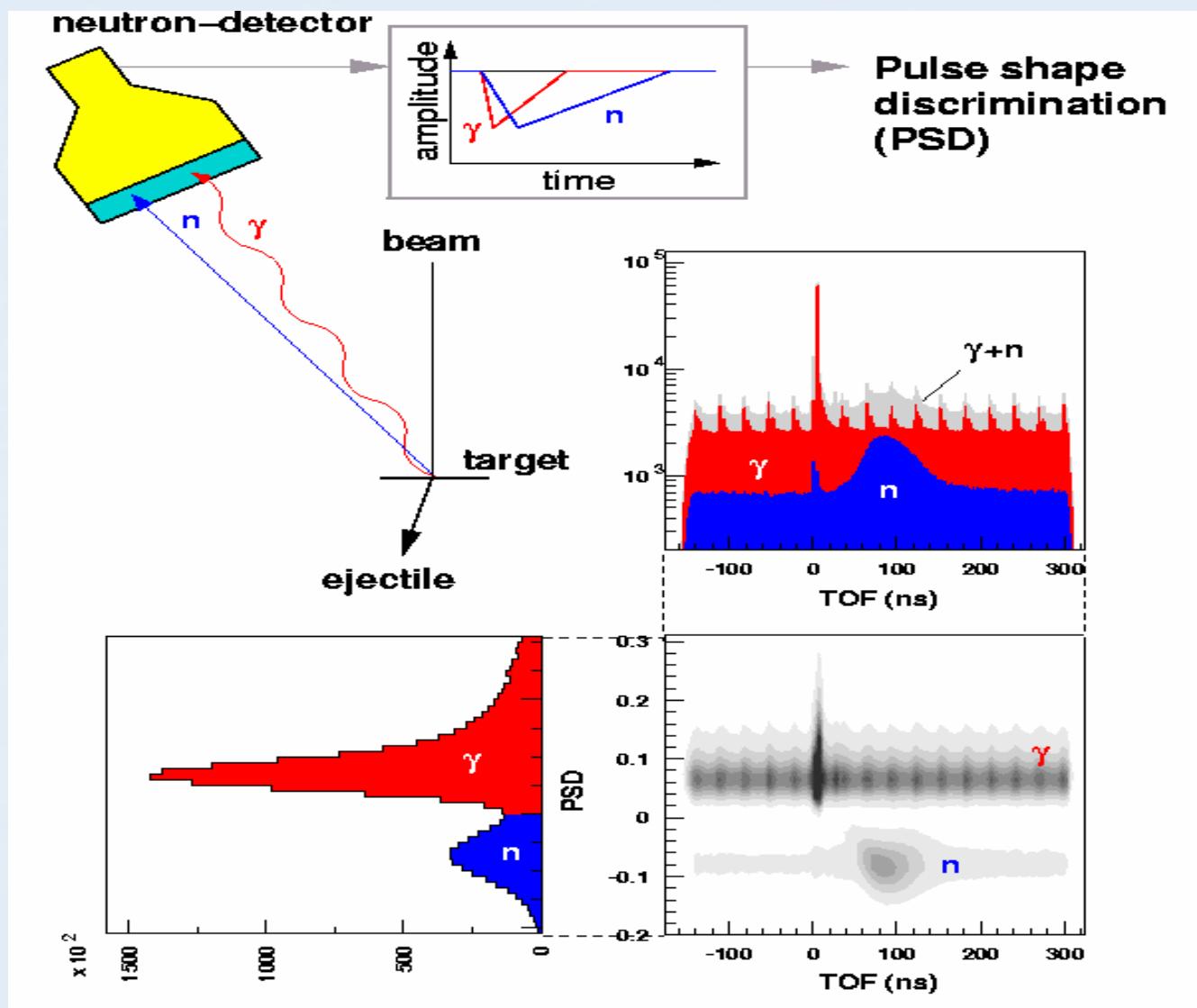
Proton-decay detection



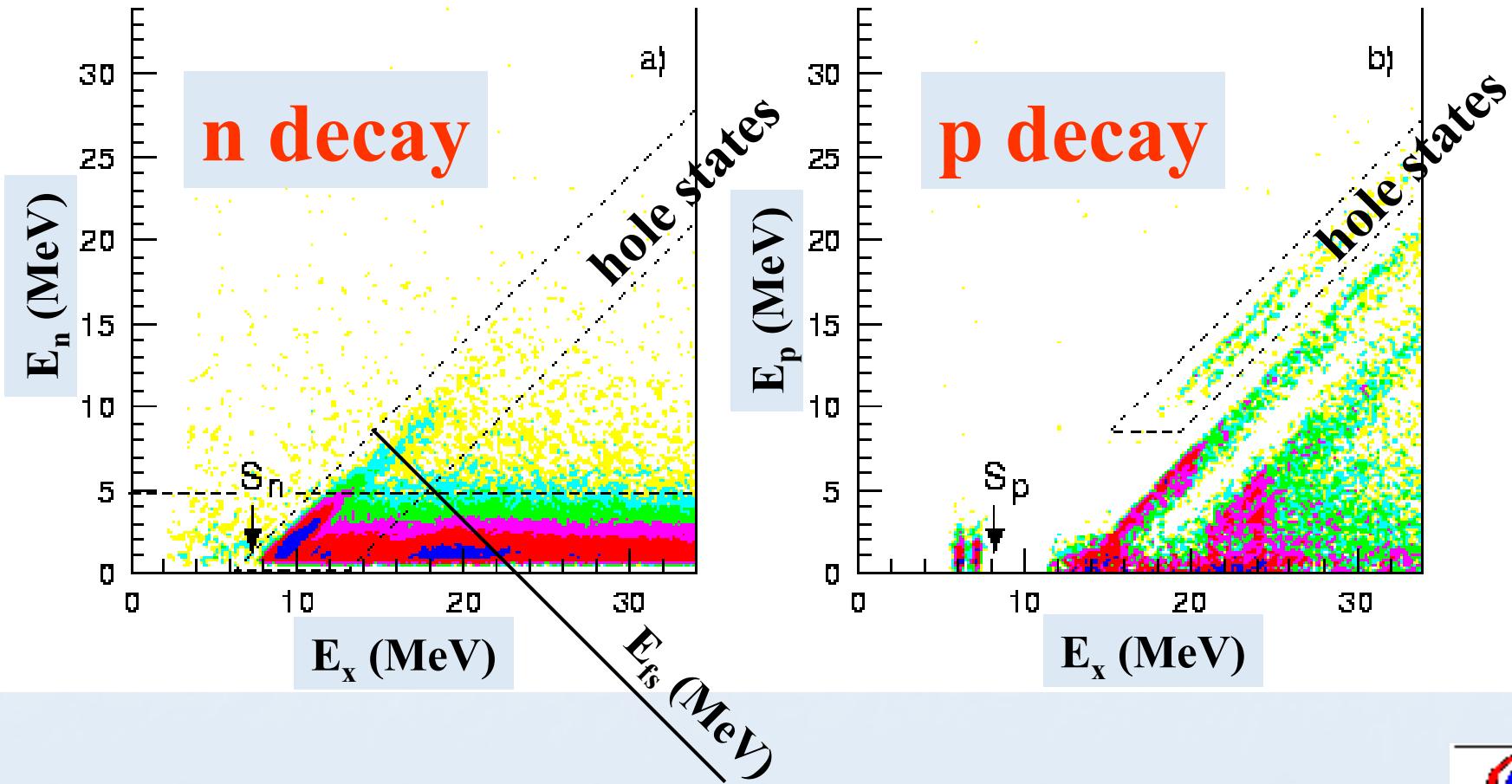
**α -p separation using
rise time of signal SiLi**



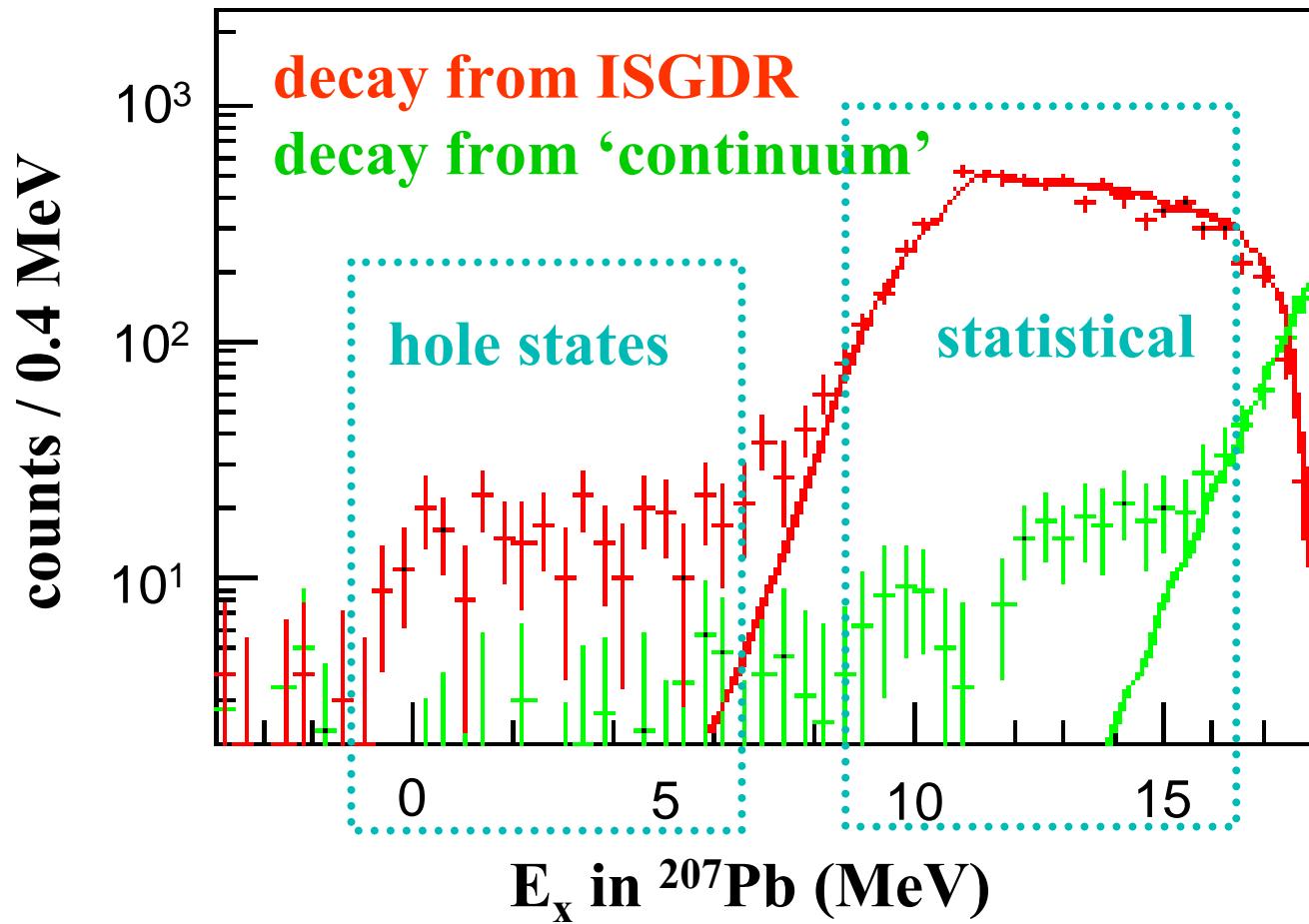
Neutron-decay detection



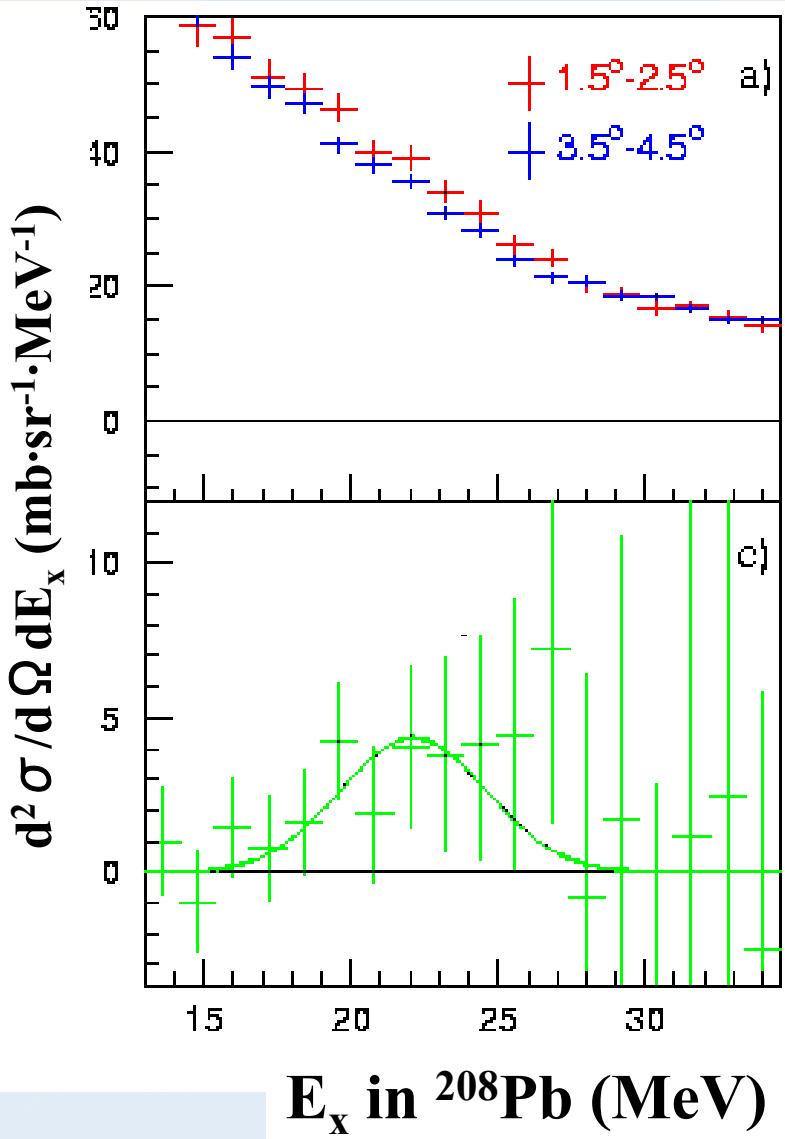
$^{208}\text{Pb}(\alpha, \alpha' \text{ p or n})$



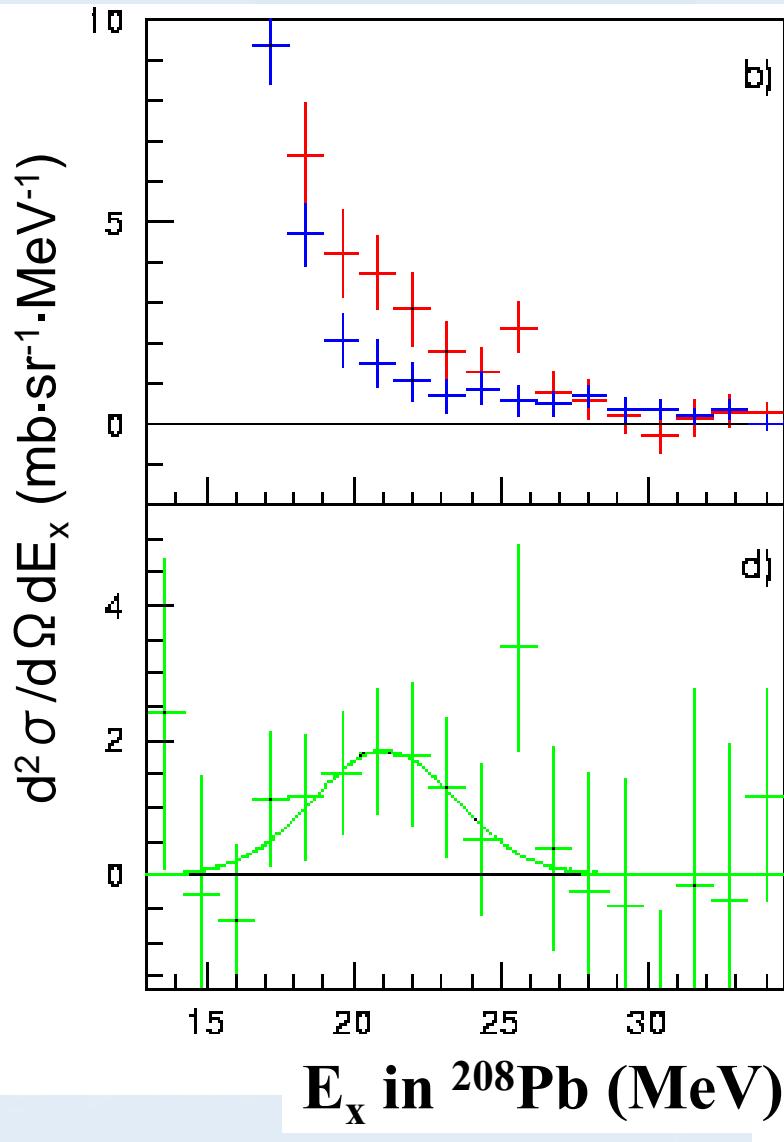
$^{208}\text{Pb}(\alpha, \alpha')$ followed by n decay



Statistical decay

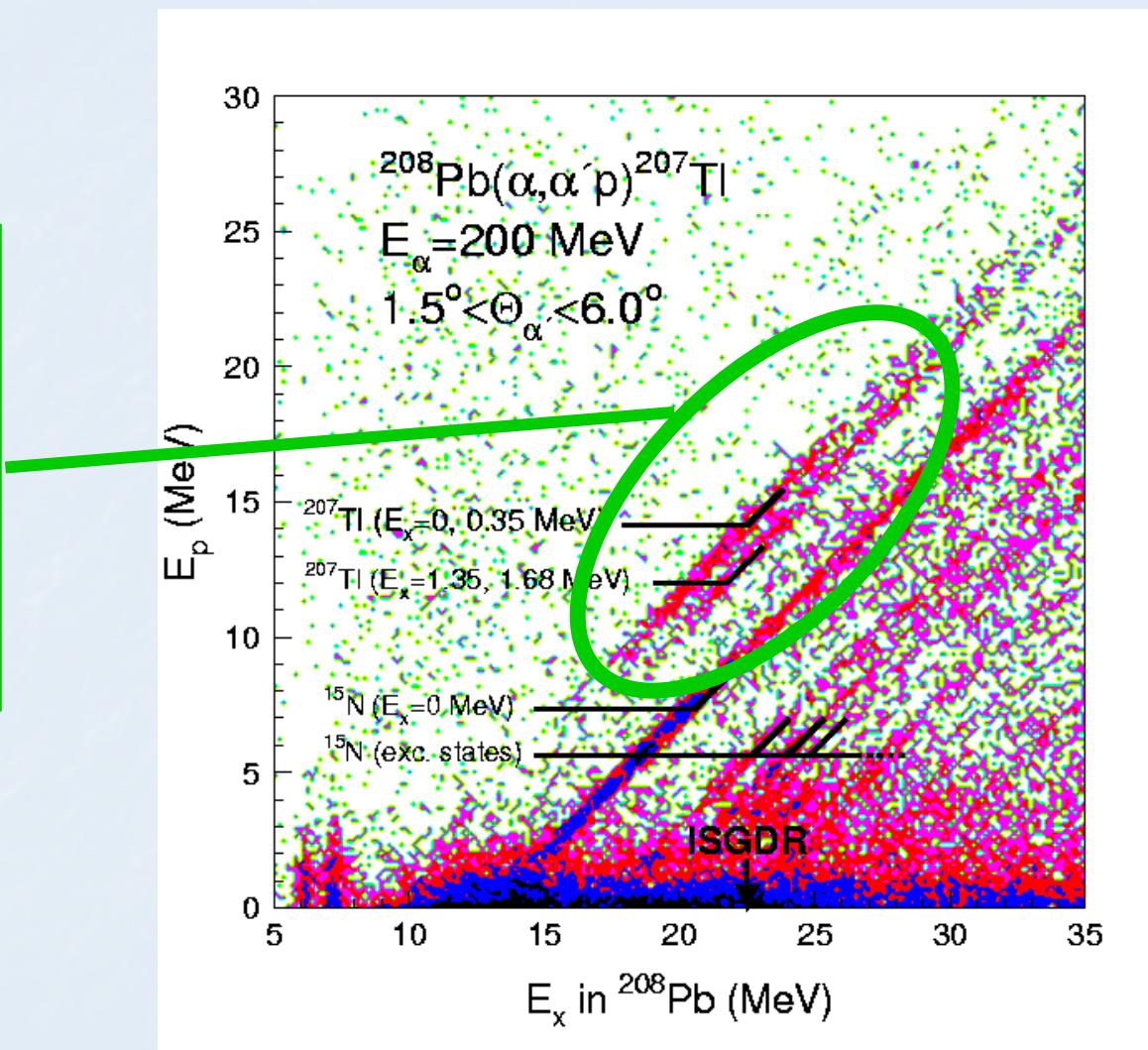


Direct decay

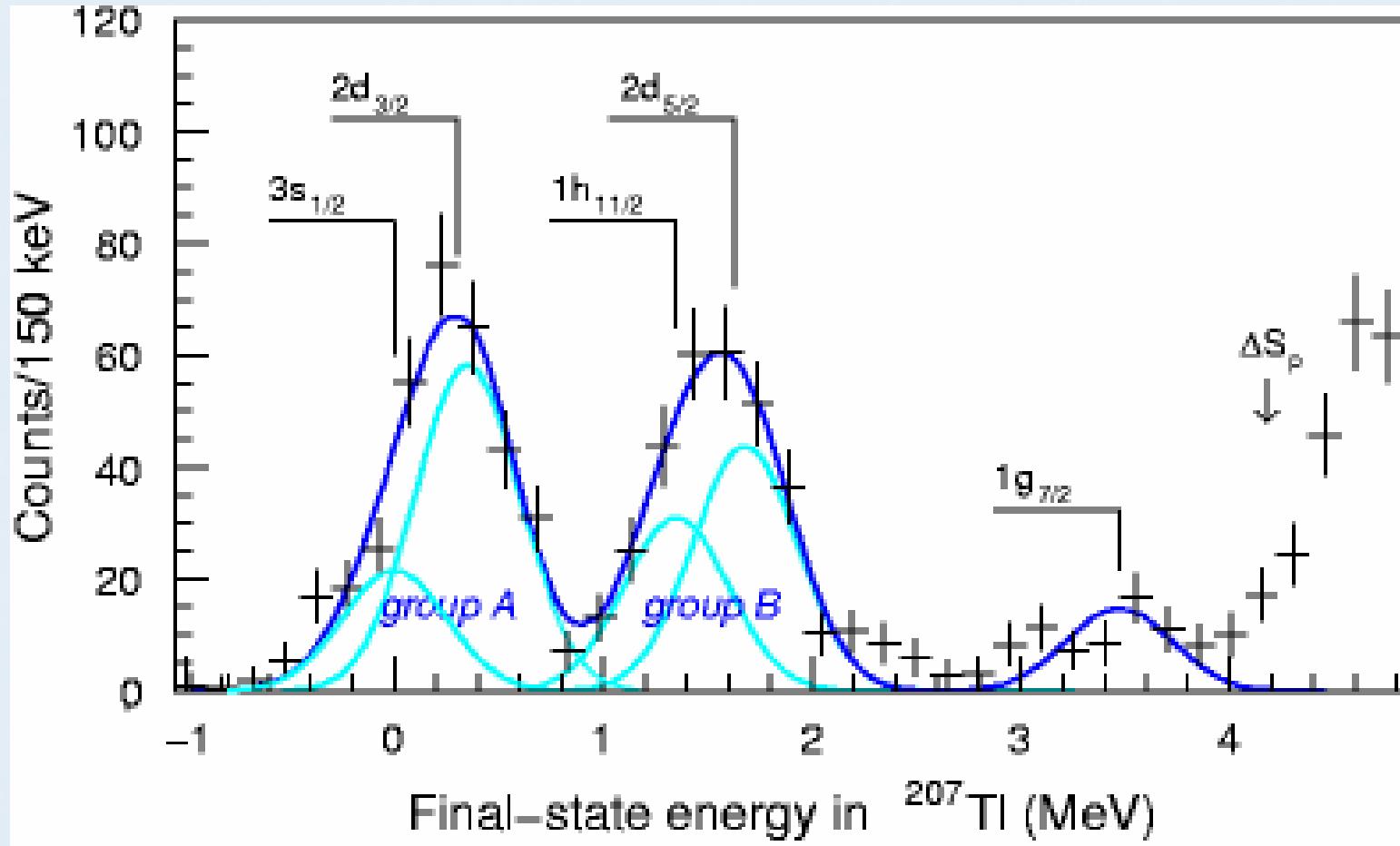


$^{208}\text{Pb}(\alpha, \alpha')$ followed by p decay

Decay to hole states in ^{207}Tl ;
branching ratios
predicted by
Gorelik et al.

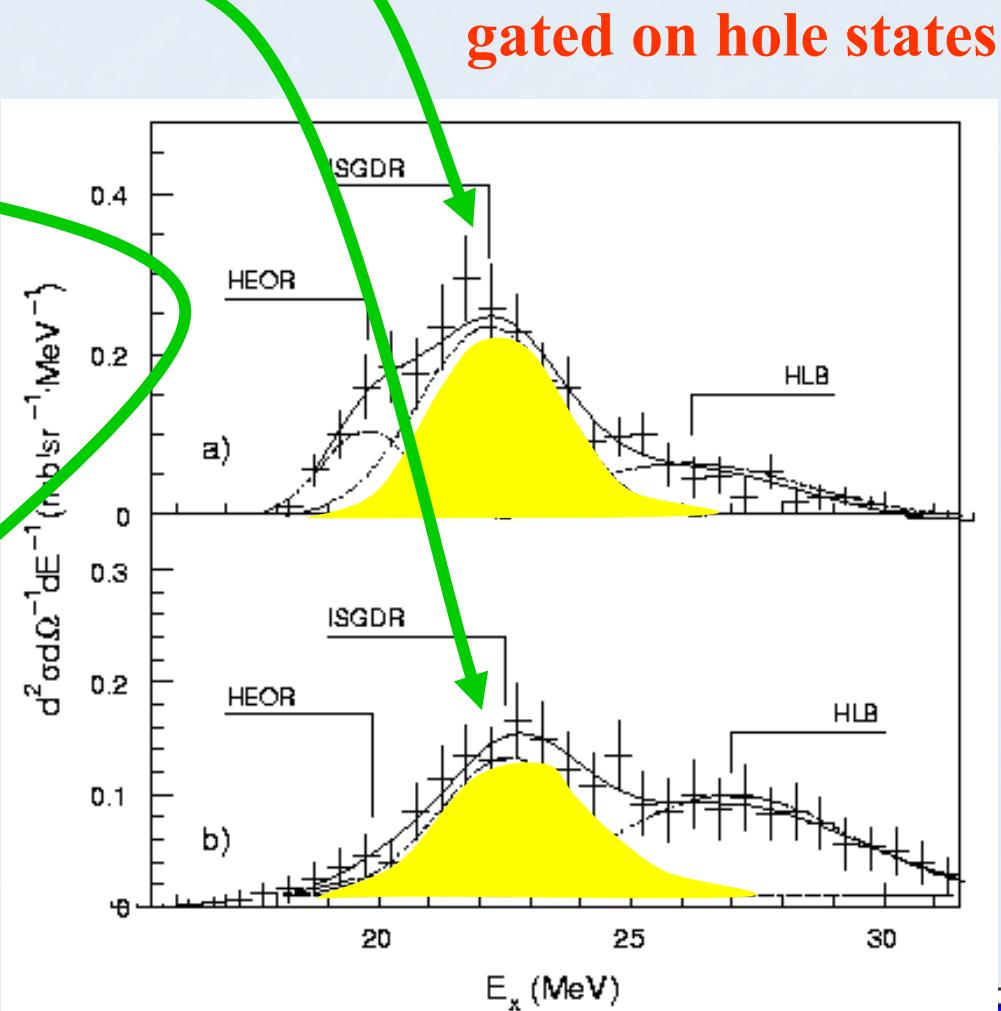
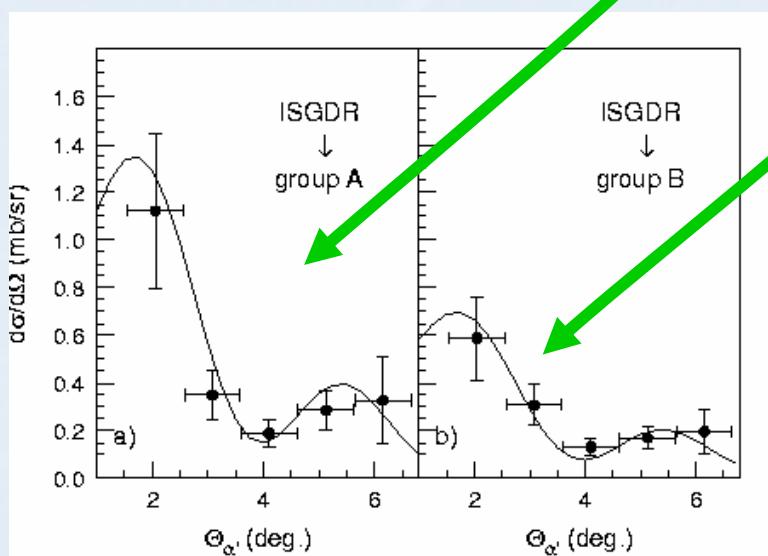


Branching ratios for decay

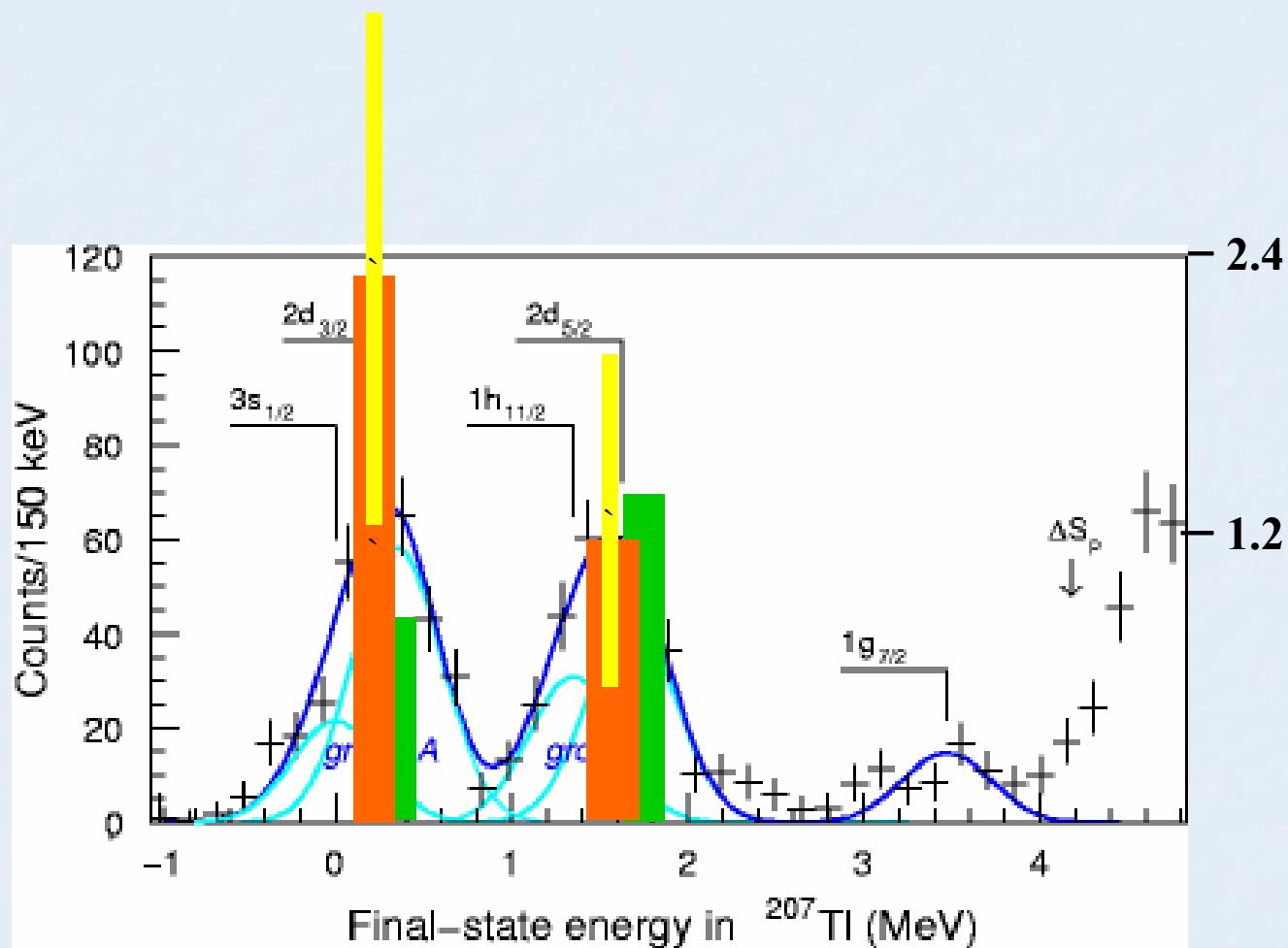


ISGDR in ^{208}Pb in p decay

$E_x = 22.1 \pm 0.3 \text{ MeV}$
 $L = 1$ transition



Branching ratios for decay



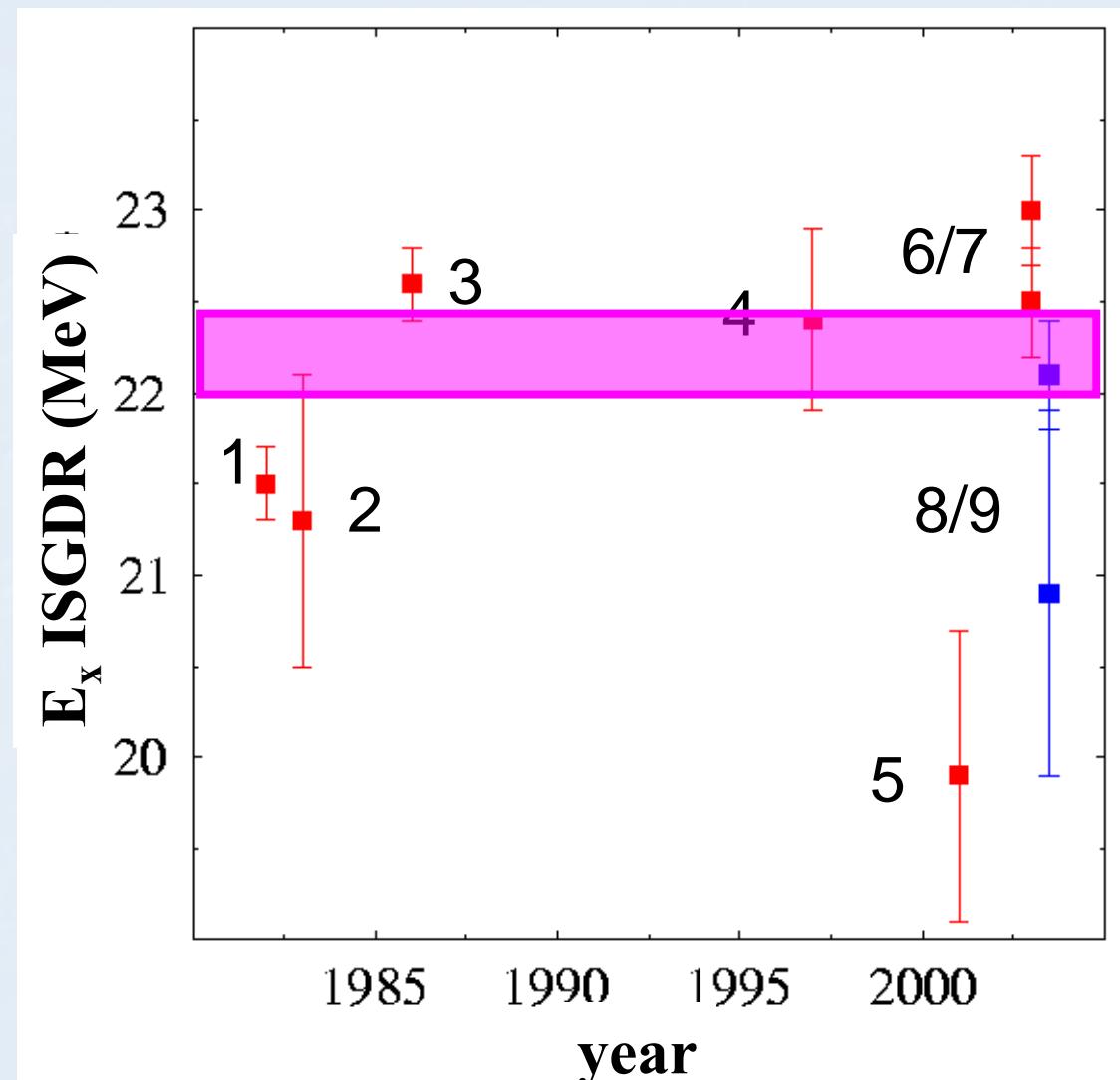
This work

Gorelik et al
PRC 62 (2000)
047301

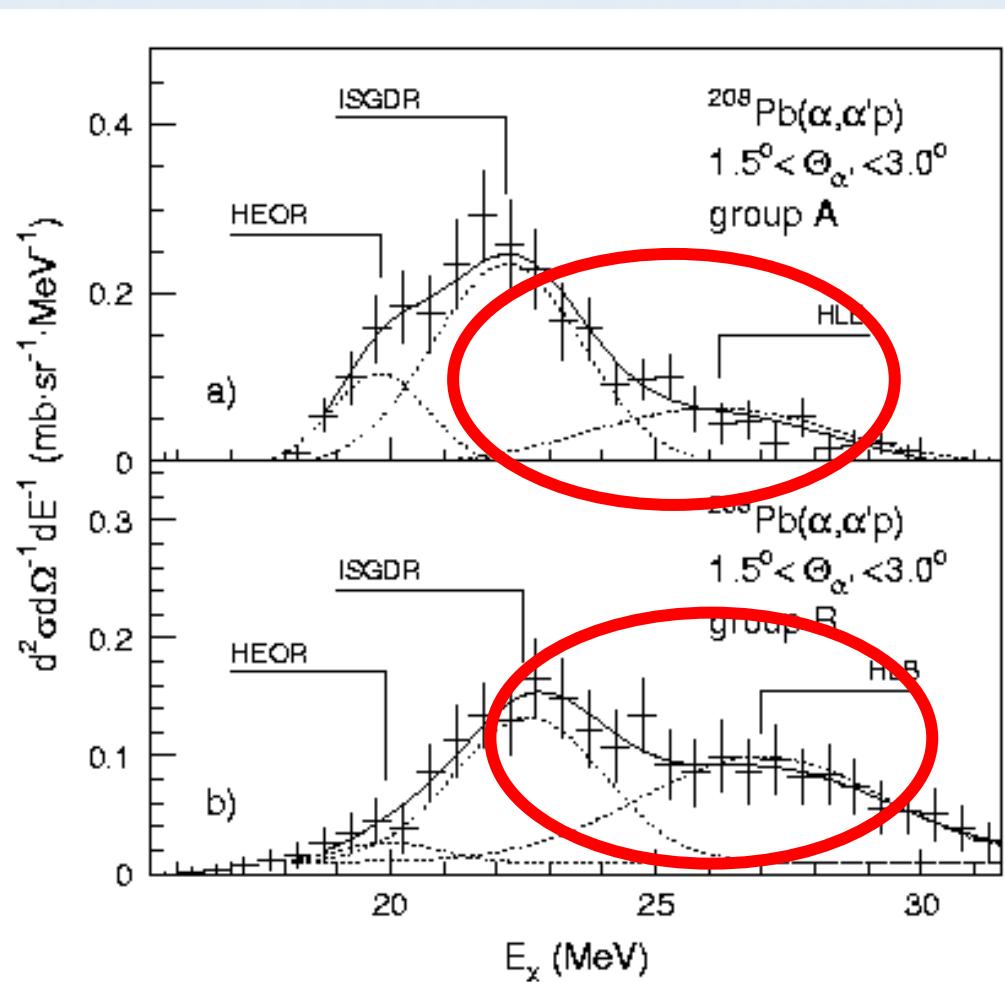
Present Status

- 1 Djalali et al., 1982
- 2 Morsch et al., 1983
- 3 Adams et al., 1986
- 4 Davis et al., 1997
- 5 Clark et al., 2001
- 6 Uchida et al., 2003
- 7 Uchida et al., 2003
- 8 this work *p*
- 9 this work *n*

$$E_x = 22.2 \pm 0.2 \text{ MeV}$$

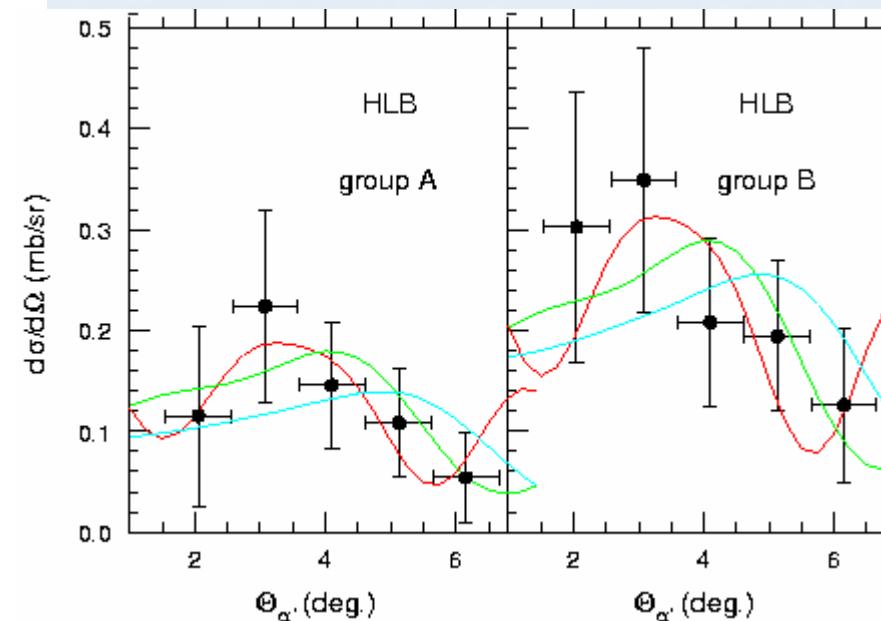


Overtone of the ISGQR? [r^4Y_2]



$$E_x = 26.9 \pm 0.7 \text{ MeV}$$

Muraviev and Urin
Bull. Acad. Sci. USSR
Phys. Ser. 52 (1988) 123
 $E_x = 28.3 \text{ MeV}$



Conclusions!

- There has been much progress in understanding ISGMR & ISGDR

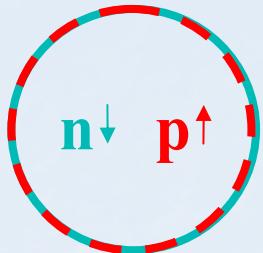
Systematics: $E_x, \Gamma, \%EWSR$

$$\Rightarrow K_{nm} \approx 220 \text{ MeV}$$

Microscopic Structure for a few nuclei

- Observation of a new excitation that could be the quadrupole compressional mode, i.e. overtone of ISGQR

Spin-isospin excitations

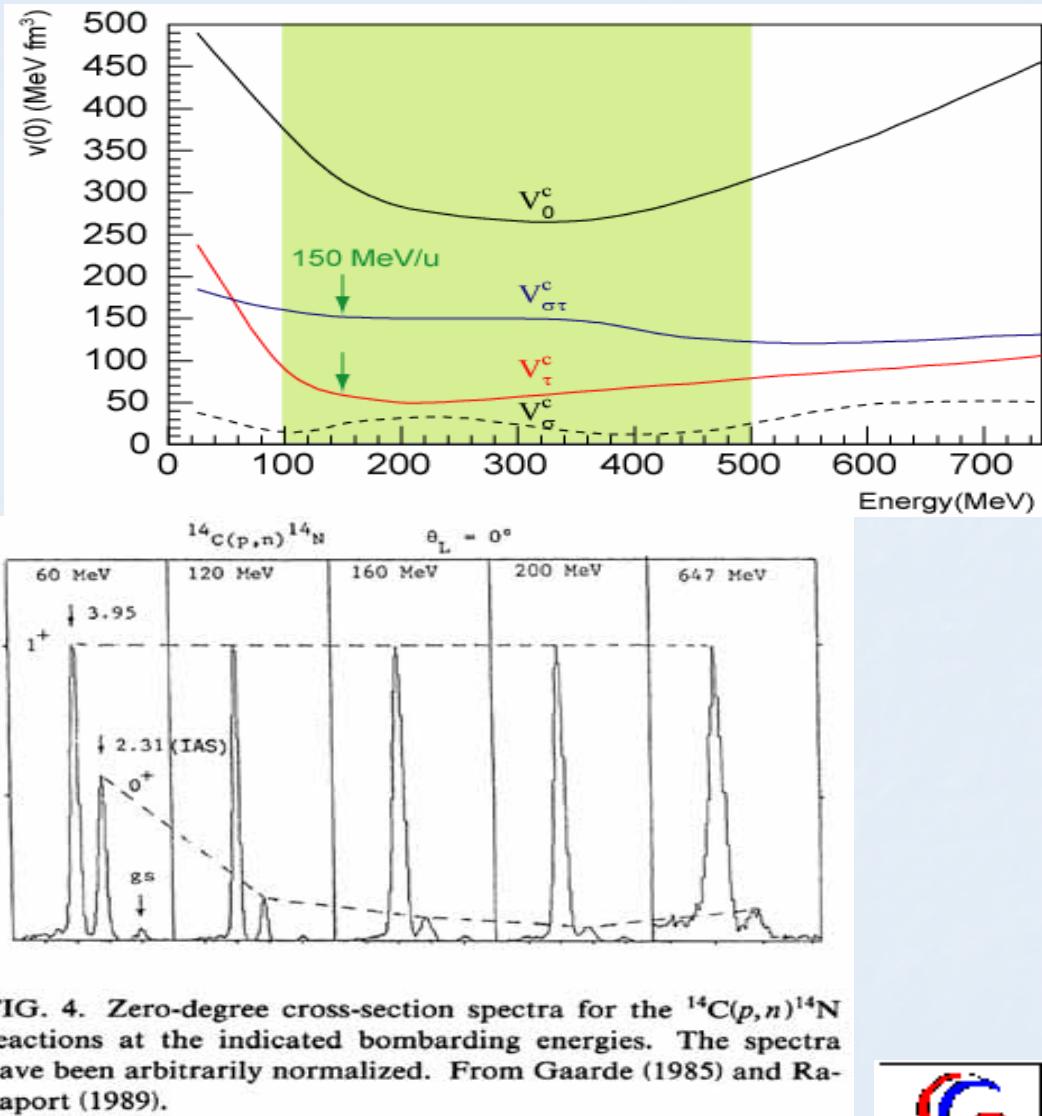


$\Delta L=0$ $\Delta S=1$ $\Delta T=1$
GTR

- Gamow-Teller transitions;
Isospin ($\Delta T=1$)
Spin ($\Delta S=1$)

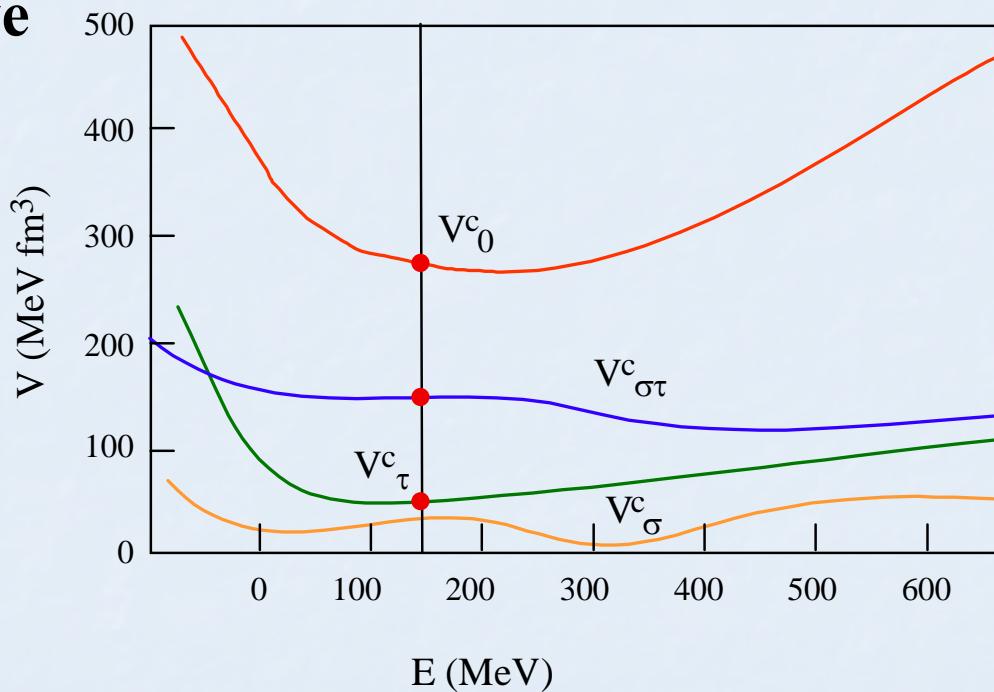
Advantages

- Cross section peak at $\theta=0$ deg. ($\Delta L=0$)
- Strong excitation of GT states at $E/A=100-500$ MeV/u



$(^3\text{He},t)$ Reaction above 150 MeV/u

- Energy dependence of effective interactions.
- 150 MeV/A
 - V_0 part: Minimum.
 - $V_{\sigma\tau}$ part: Relatively large.
 - V_t part: Minimum.



Charge-exchange probes

(p,n)-type ($\Delta T_z = -1$)

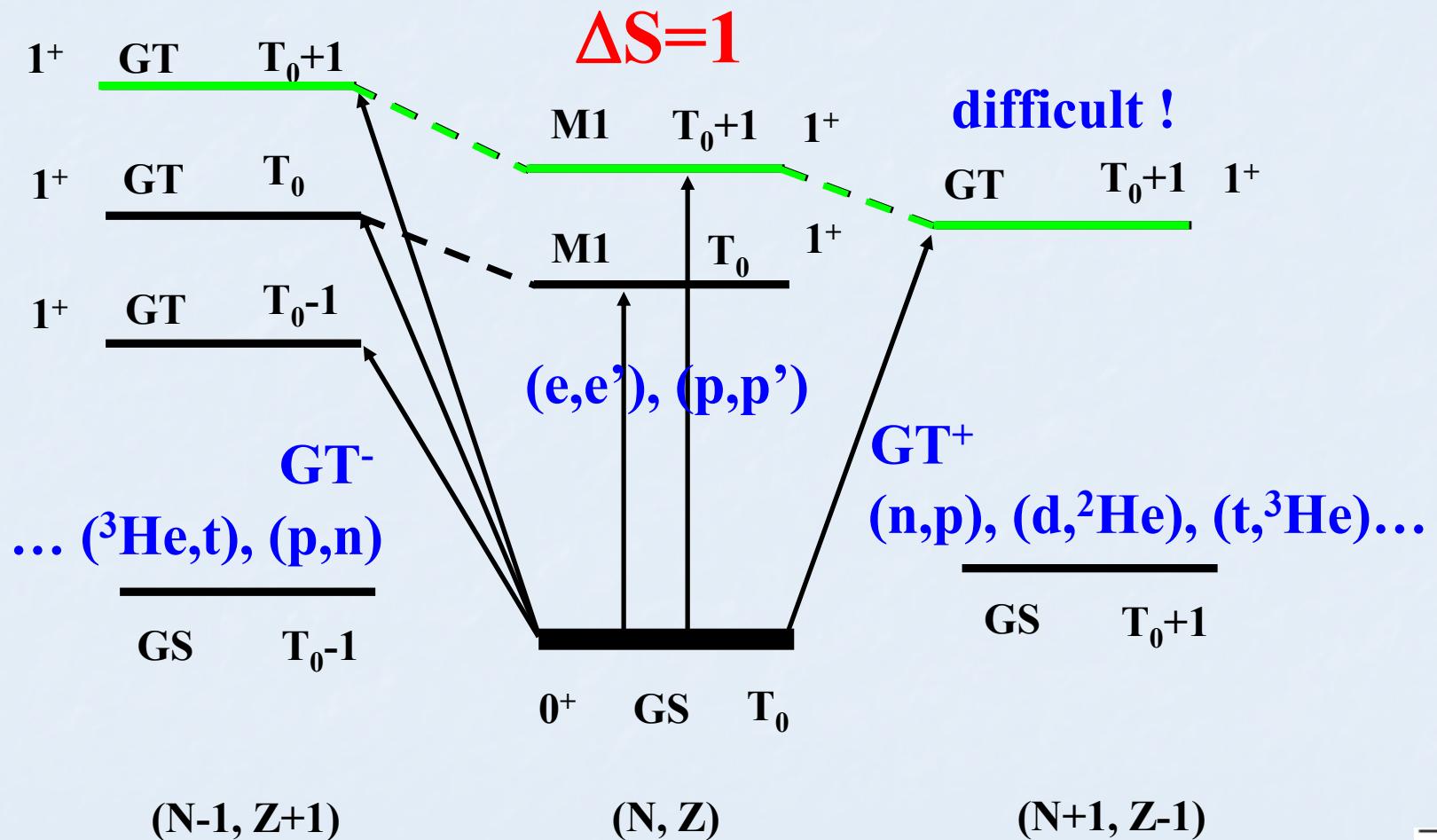
- β^- -decay
- (p,n)
- ($^3\text{He},\text{t}$)
- heavy ion
- (π^+, π^0)

(n,p)-type ($\Delta T_z = +1$)

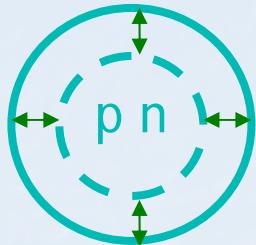
- β^+ -decay
- (n,p)
- (d, ^2He)
- (t, ^3He)
- heavy ion ($^7\text{Li}, ^7\text{Be}$)
- (π^-, π^0)

- Energy per nucleon (>100 MeV/n)
- Spin-flip vs non-spin-flip
- Complexity of reaction mechanism
- Experimental considerations

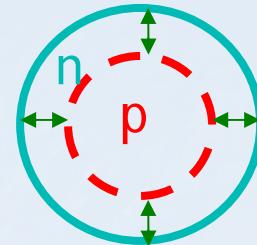
Nuclear Classics: spin-flip & GT transitions



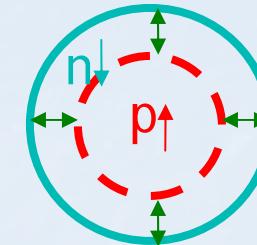
Isovector giant monopole resonances



$\Delta L=0 \Delta S=0 \Delta T=0$
ISGMR



$\Delta L=0 \Delta S=0 \Delta T=1$
IVGMR



$\Delta L=0 \Delta S=1 \Delta T=1$
IVSGMR

$$O = r^\lambda [\sigma \otimes Y_L]_J \tau_-$$

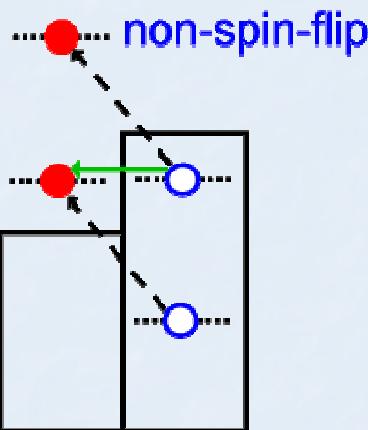
IAS: $\lambda=0 S=0 L=0 J=0$

GTR: $\lambda=0 S=1 L=0 J=1$

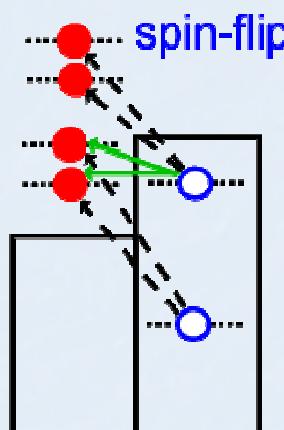
IVGMR: $\lambda=2 S=0 L=0 J=0$

IVSGMR: $\lambda=2 S=1 L=0 J=1$

IVSGDR: $\lambda=1 S=1 L=1 J=0,1,2$



IAS IVGMR
 $\Phi_{njl} \rightarrow \Phi_{njl}$ $\Phi_{njl} \rightarrow \Phi_{n+1jl}$



GTR IVSGMR
 $\Phi_{njl} \rightarrow \Phi_{njl}$ $\Phi_{njl} \rightarrow \Phi_{n+1jl}$
 $\Phi_{njl} \rightarrow \Phi_{nj+1l}$ $\Phi_{njl} \rightarrow \Phi_{n+1j+1l}$



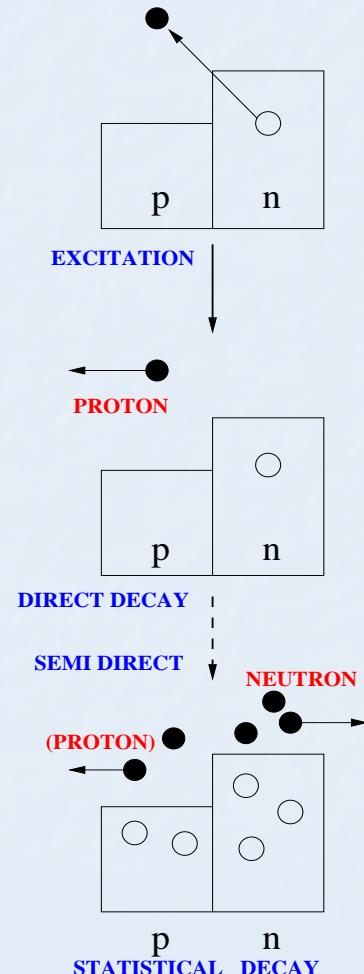
Decay studies

■ Successful:

- GTR, IVSGDR in $^{208}\text{Pb}(\text{He}^3, \text{t}+\text{p})$ at 450 MeV (Akimune et al.)
- IVGMR/IVSGMR in $\text{Pb}(\text{He}^3, \text{t}+\text{p})$ at 177 MeV at KVI & 410 MeV at RCNP (Zegers et al.)

■ Unsuccessful:

- IVGMR/IVSGMR $^{124}\text{Sn}(\text{He}^3, \text{t}+\text{n})$ at 200 MeV at IUCF



Microscopic Structure of GTR and IVSGDR in ^{208}Bi

- Proton decay of ^{208}Bi :
 - Direct decay dominant
 - $E_x > E_{\text{th}}(n) > E_{\text{th}}(p)$
 - High Coulomb Barrier ($Z=83$)
 - Statistical proton decay negligible.
- Angular correlations
 - For IAS and GTR decay isotropic $\Delta L=0$
 - For IVSGDR anisotropic but not strongly
- Direct decay is influenced by:
 - Low n-decay threshold
 - High Coulomb barrier.

- $\Gamma_{\text{GTR}}^{\uparrow}/\Gamma \ll \Gamma_{\text{IAS}}^{\uparrow}/\Gamma \approx 0.5$
 - IAS n-decay: isospin forbidden.
 - Centroid energy shift: cut off by Coulomb barrier
- $\Gamma_{\text{IVSGDR}}^{\uparrow}/\Gamma > \Gamma_{\text{GTR}}^{\uparrow}/\Gamma$
 - Higher proton energy
- Width Γ :

$$\Gamma = \Gamma^{\uparrow} + \Gamma^{\downarrow}$$

Escape: Direct decay
Spreading: Statistical Decay

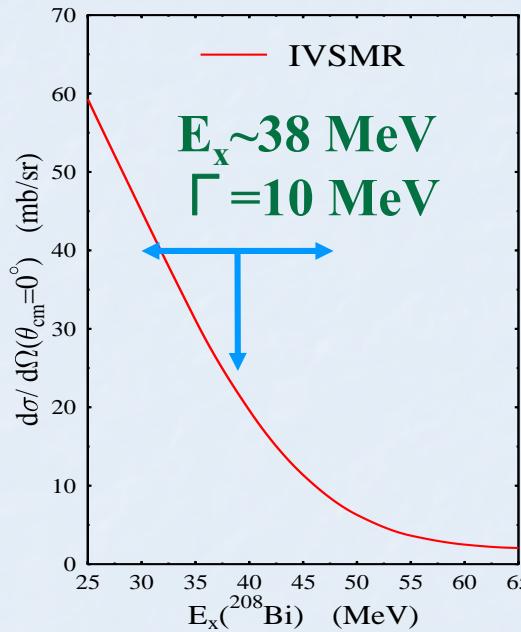
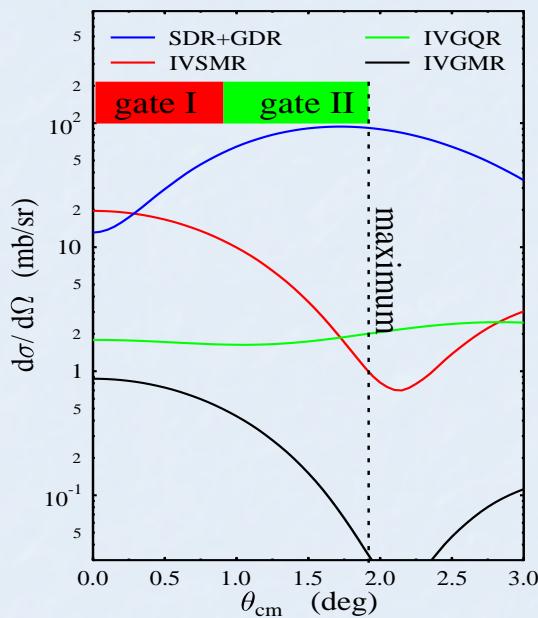
$$\Gamma^{\uparrow} = \Gamma_p^{\uparrow} = \sum_i \Gamma_{pi}^{\uparrow}$$

Partial Escape Width

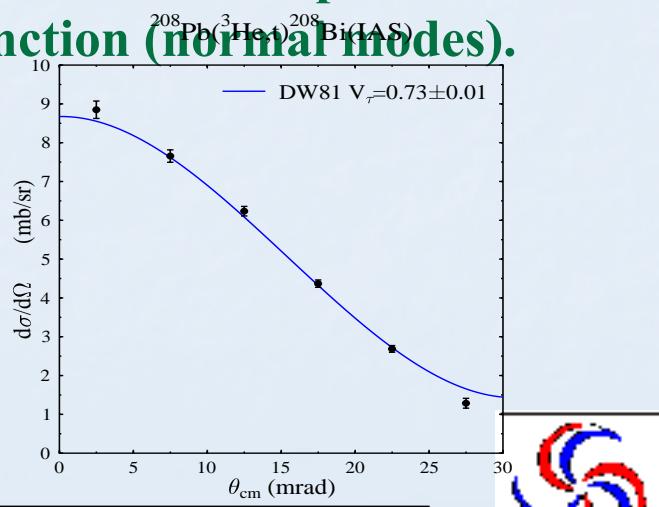
$$\frac{\Gamma_{pi}^{\uparrow}}{\Gamma} = \frac{\int d^2\sigma_{pi}/(d\Omega_t d\Omega_p) d\Omega_p}{d\sigma/d\Omega_t}$$

Branching ratio

Measurement of IVSGMR via $^{208}\text{Pb}(\text{He}^3, t+p)$



- DW81 (Raynal)
- Effective $^3\text{He-N}$ potential
 - $V_\tau = 0.73 \pm 0.01 \text{ MeV}$ (IAS)
 - $V_{\sigma\tau} = -2.1 \pm 0.2 \text{ MeV}$ (known ratio to V_τ)
 - $V_{T\tau} = -2.0 \text{ MeV/fm}^2$
- most coherent 1p-1h wavefunction (normal modes).



Use *difference-of-angle* to identify the monopole excitations

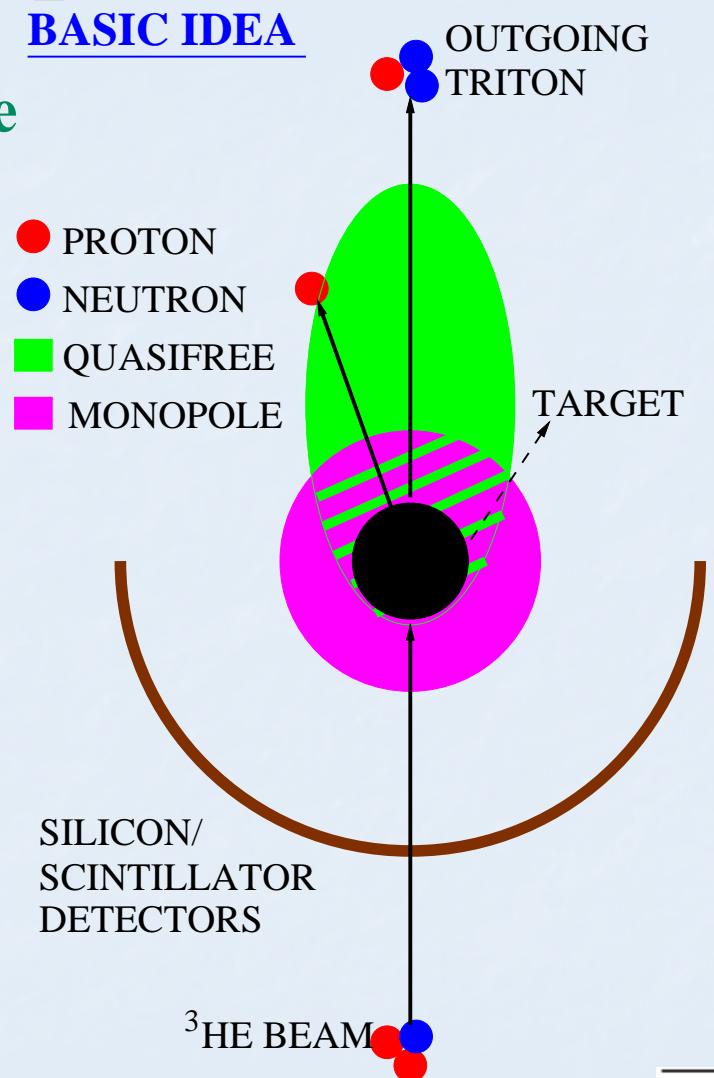
Continuum suppression

Physical background (continuum) due to:

- breakup-pickup reactions
- quasifree knock-on reactions

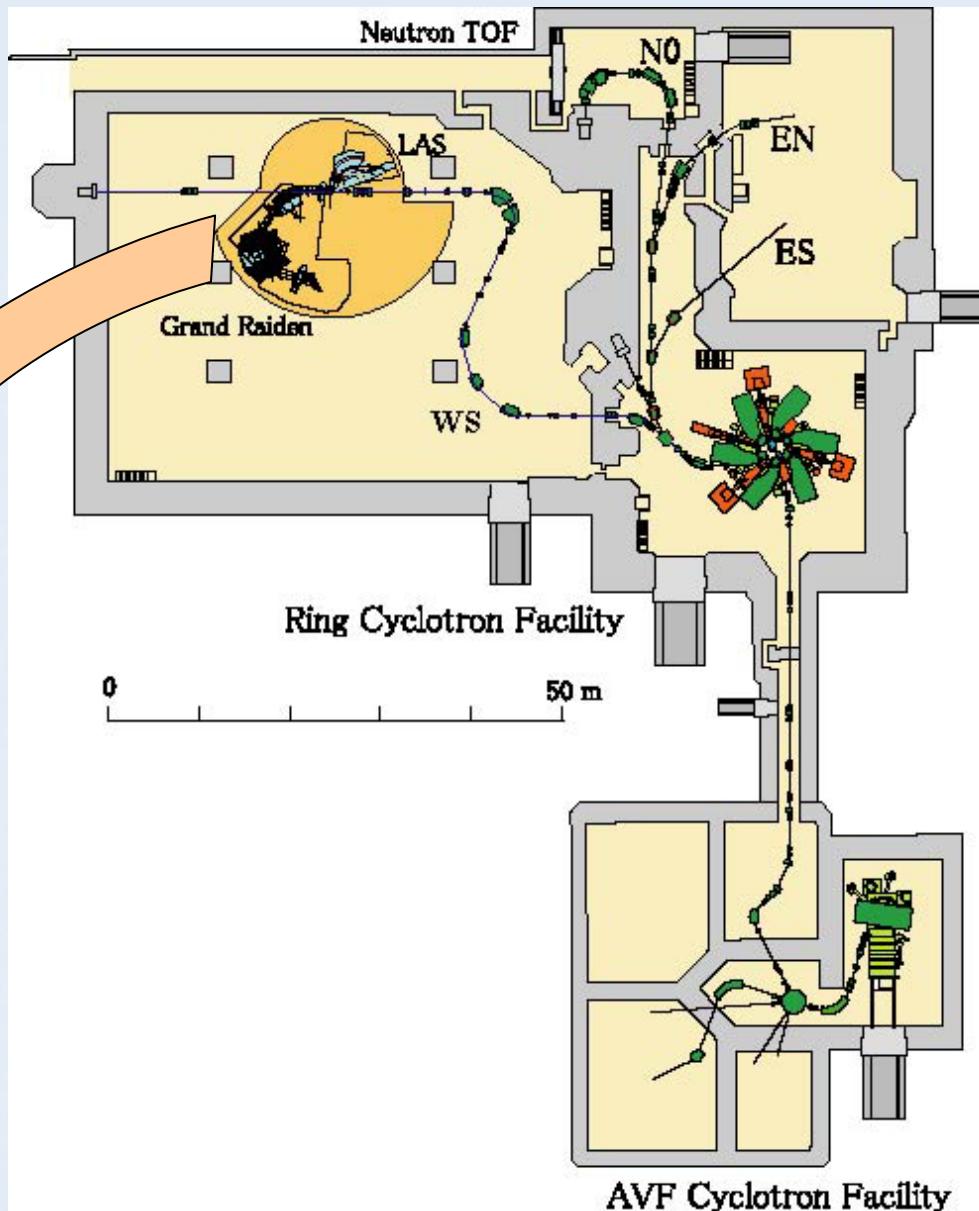
IVSGMR is wide (~10 MeV) and lying on top of the continuum.

Continuum has a very flat angular distribution at forward angles



Experiment

- RCNP facility
K=400 MeV ring cyclotron
Grand Raiden spectrometer
- Beam: ${}^3\text{He}^{++}$, 450 MeV
- Target: ${}^{208}\text{Pb}$ foil

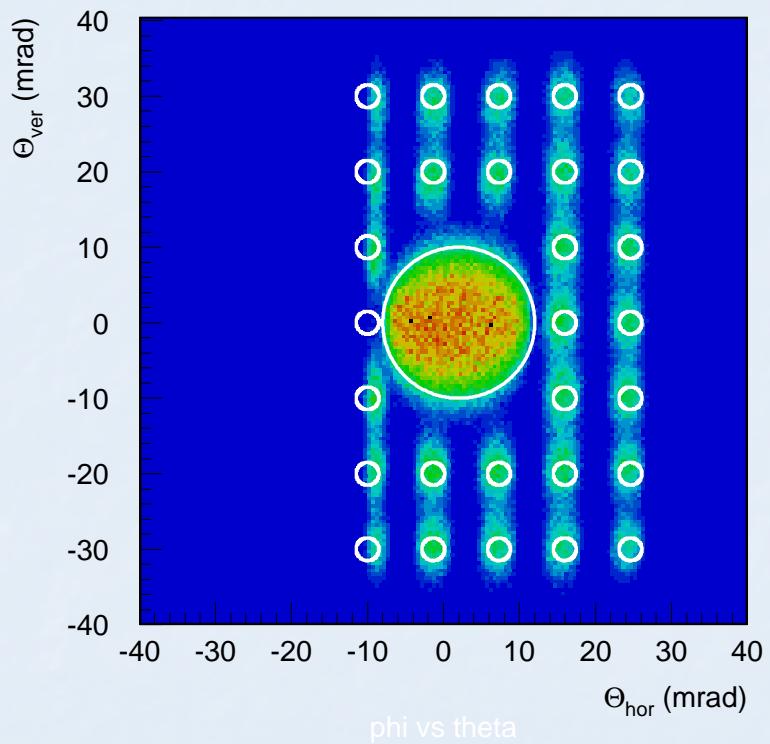


M. Fujiwara et al., NIM A 422 (1999) 484

Experimental Considerations

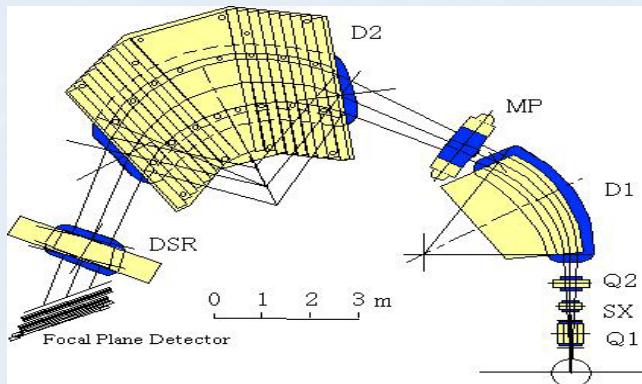
- Over-focus mode for Grand-Raiden (H. Fujita et al. NIM A469, 55) refined to get:

- Vertical angle resolution:
0.4° FWHM
- Horizontal angle resolution:
0.2° FWHM
- Negligible systematic errors
- Sieve-slit is used for calibrations

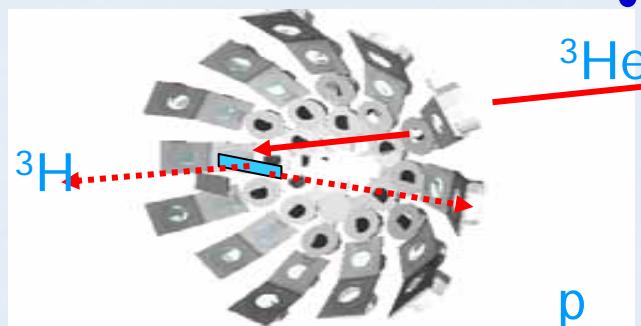


Experiment: 410 MeV ^3He -beam @ RCNP

Grand Raiden @ RCNP

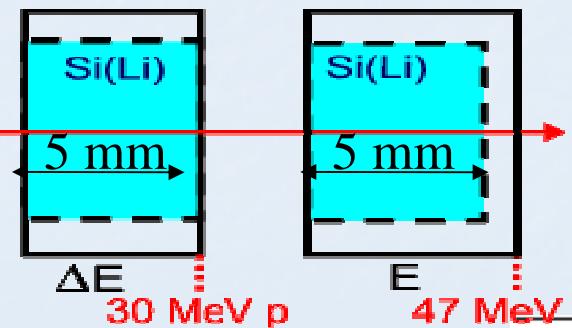


- Measure both t-singles and t-p coincidences.



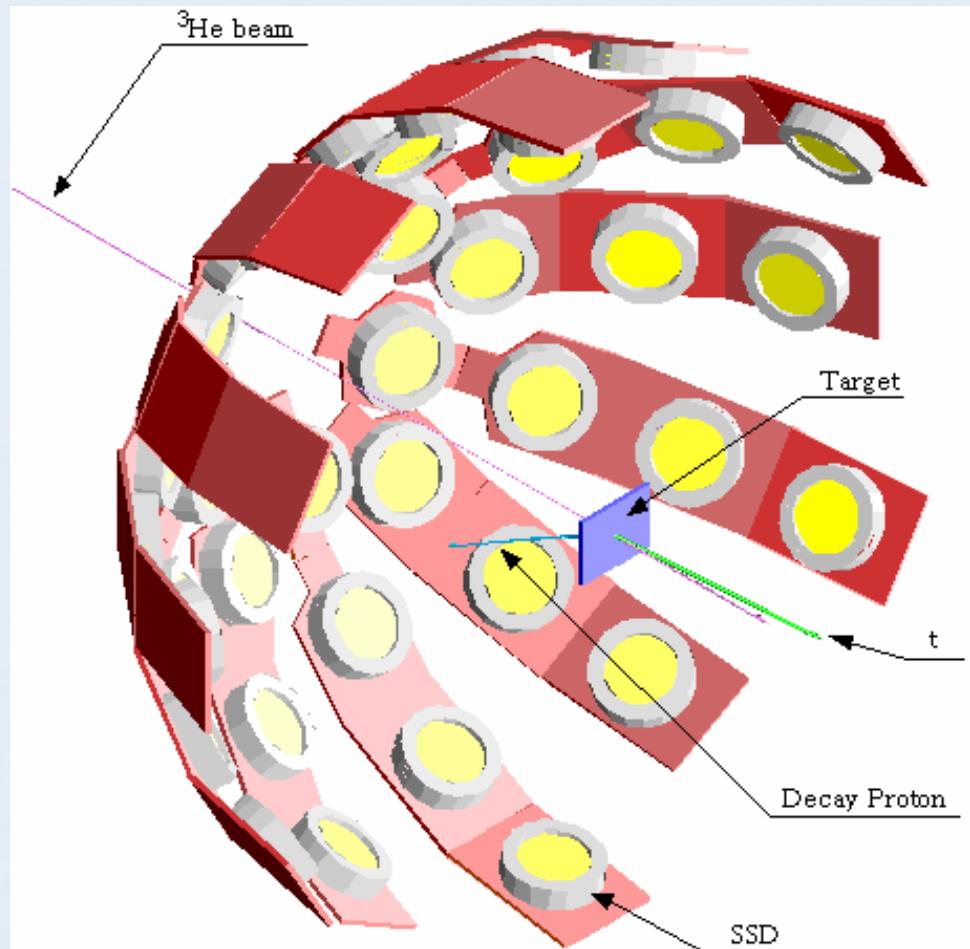
- 8 ΔE -E telescopes

- 4@ 113°
- 4@ 136°



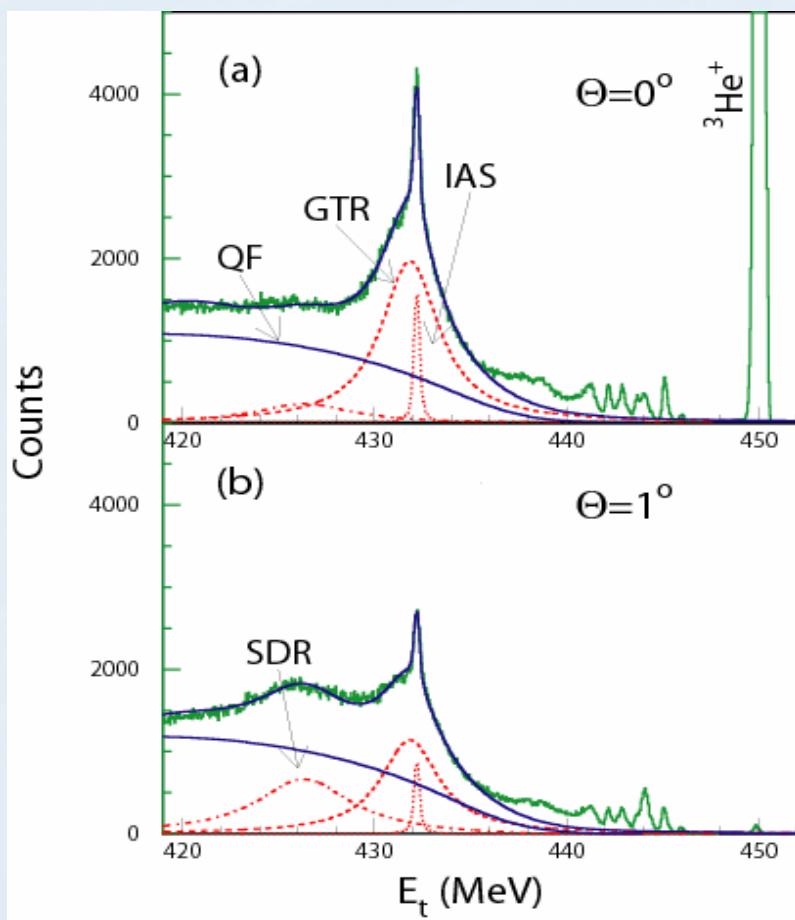
Set-up of the Proton Counter

- Si(Li) detectors with a thickness of 5 mm, covering a solid angle of 5.7% in total.
- 35 keV (^{241}Am test)



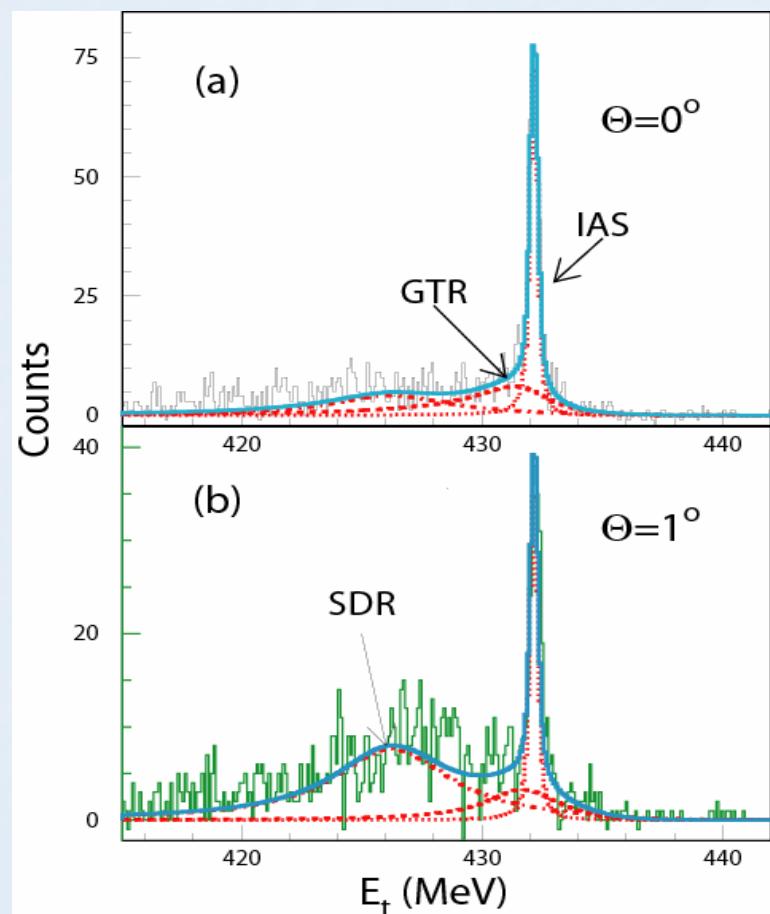
Spin-isospin-flip transitions in charge-exchange reactions and proton decay

$(^3\text{He},t)$ reactions $E(^3\text{He})=450 \text{ MeV}$



- A. Krasznahorkay et al., PRC 64 (2001) 067302.
A. Krasznahorkay et al., PRL 82 (1999) 3216.
H. Akimune et al., PRC 52 (1995) 604.
H. Akimune et al., Phys. Lett. B 233 (1994) 107

$(^3\text{He},tp)$ Coincidence data



Experimental Results and Theoretical Calculations

■ Partial escape width for GTR

channel	E_x (keV)	Theory		This work	
		Γ_i^\uparrow (keV)	branch (%)	Γ_i^\uparrow (keV)	branch (%)
$3p_{1/2}$	0	48.7	1.23	58.4 ± 9.8	1.8 ± 0.5
$2f_{5/2}$	570	46.2	2.12	inc. in $p_{3/2}$	
$3p_{3/2}$	898	44.7	2.5	101.5 ± 31.3	2.7 ± 0.6
$1i_{13/2}$	1633	0.87	3.57	8.3 ± 0.4	0.2 ± 0.2
$2f_{7/2}$	2340	5.89	2.97	15.6 ± 7.6	0.4 ± 0.2
$1h_{9/2}$	3413	0.24	0.63	—	—
Total		146.6	13.02	184 ± 49	4.9 ± 1.3

Theory:
E. Moukhai,
V.A. Rodin,
M.H. Urin
Continuum RPA

● Partial escape width for IVSGDR

channel	E_x (keV)	Theory		This work	
		Γ_i^\uparrow (keV)	branch (%)	Γ_i^\uparrow (keV)	branch (%)
$3p_{1/2}$	0	103.4	1.23	83.4 ± 24.3	0.99 ± 0.29
$2f_{5/2}$	570	178.1	2.12	170.8 ± 49.3	2.12 ± 0.61
$3p_{3/2}$	898	210.1	2.5	240 ± 69.6	2.86 ± 0.83
$1i_{13/2}$	1633	299.8	3.57	330.4 ± 95.7	3.74 ± 1.08
$2f_{7/2}$	2340	249.3	2.97	282.2 ± 86.8	3.36 ± 0.97
$1h_{9/2}$	3413	52.6	0.63	86.7 ± 25.1	1.03 ± 0.29
Total		1209.6	14.4	1180 ± 340	14.1 ± 4.2

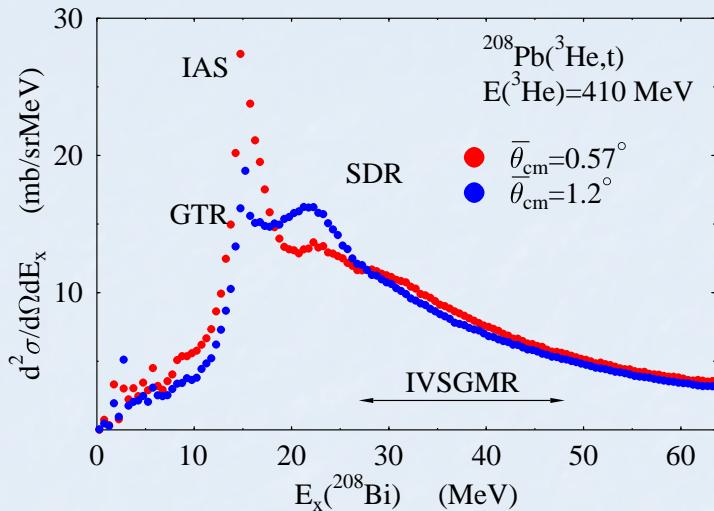
Summary: $^{208}\text{Pb}({}^3\text{He},\text{tp})$

- GTR: $\Gamma^\uparrow/\Gamma \sim 4.9\%$, $\Gamma^\uparrow = 184 \pm 49$ keV
 - Small branching ratio:
 - Spreading effect is very important.
 - Coupling to underlying 2p-2h states.
 - Centroid energy shift caused by
High Coulomb barrier.
- IVSGDR: $\Gamma^\uparrow/\Gamma \sim 14.1\%$, $\Gamma^\uparrow = 1180 \pm 340$ keV
 - Larger p-decay Γ^\uparrow/Γ compared to GTR.
 - E_p : enough higher than
Coulomb barrier, centrifugal barrier.
 - Enhancement of decay to high-spin 1n-hole states

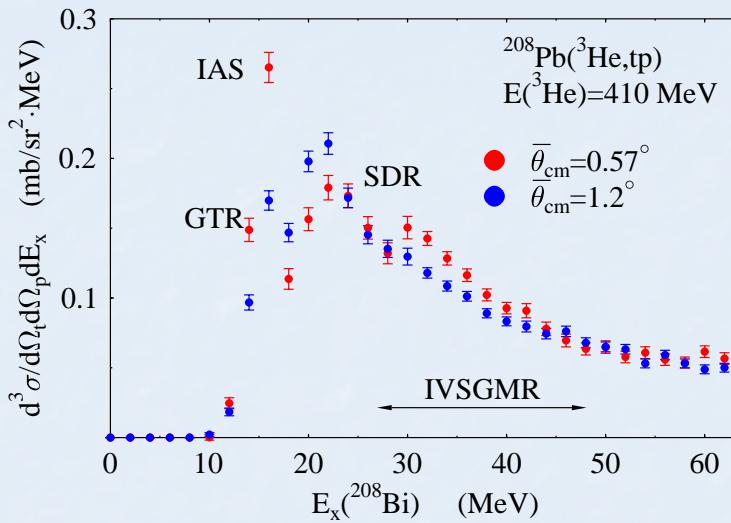
Results

R.G.T. Zegers et al., PRL 90 (2003) 202501

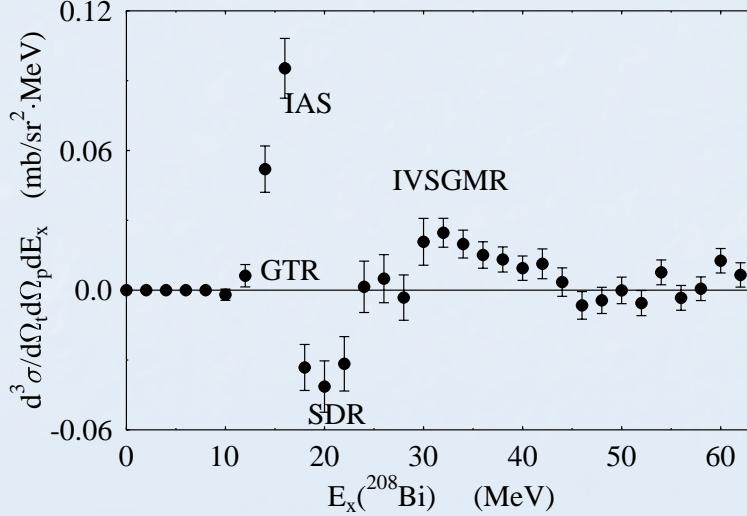
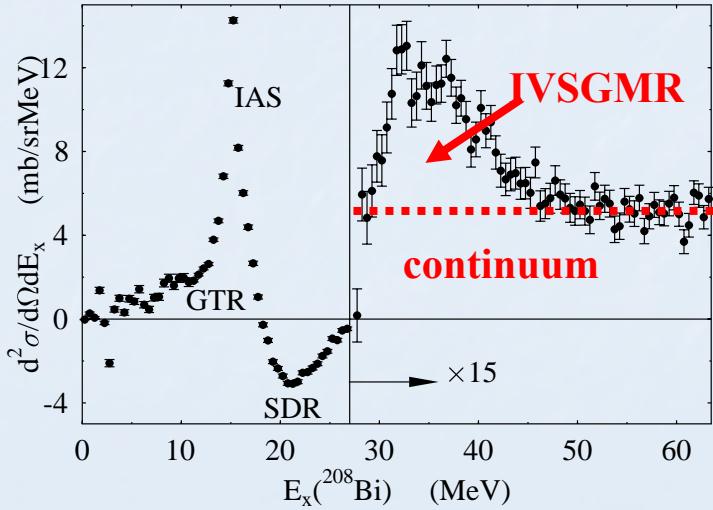
t singles



t-p coincidences

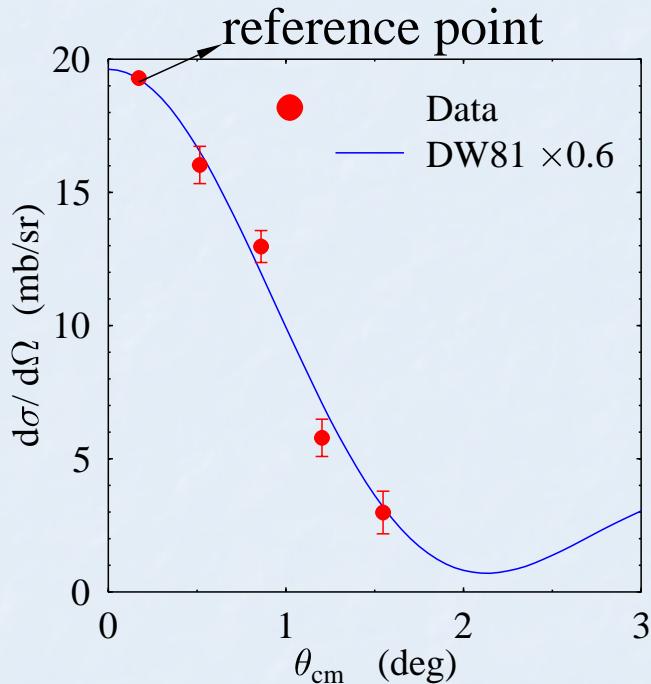


Difference of angles



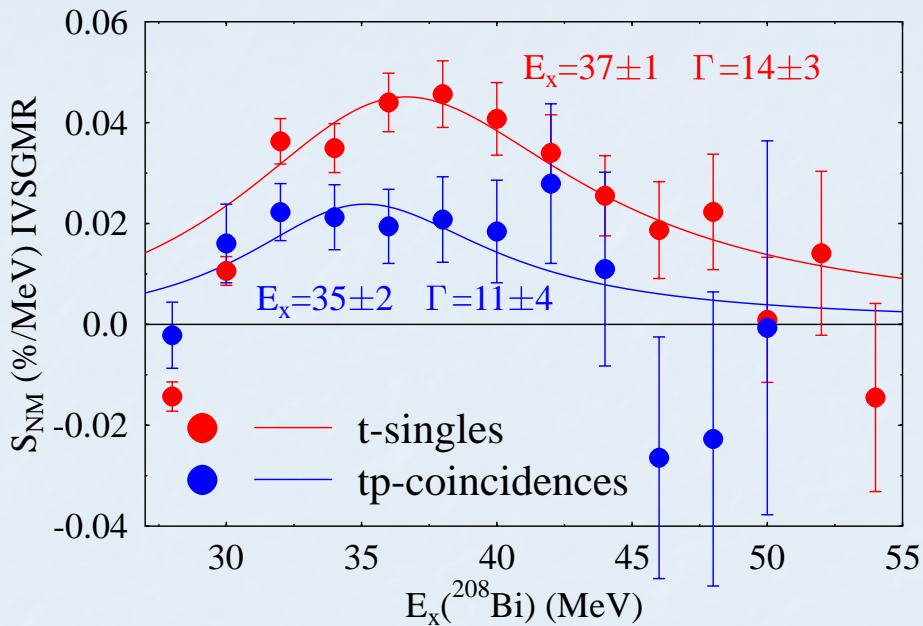
Angular distribution

Use difference-of-angle method between narrow angular bins
to extract angular distribution of the resonance



IVSGMR angular distribution confirmed

Strength exhaustion



Systematic errors:

- extrapolation of continuum: 5%
- high-lying GT strength: small
- tail of the IVSGDR: 10%
- DWBA: 10% of measured value

Summed strength: $(46 \pm 4 \pm 10) \cdot 10^3 \text{ fm}^4$
(contribution from IVGMR subtracted)

method

Normal modes

Exhaustion(%)
 $(\pm \sigma_{\text{stat}} \pm \sigma_{\text{sys}})$

$60 \pm 5 \pm 14$

Tamm-Dancoff

Hamamoto & Sagawa
PRC 62, 024319

$68 \pm 6 \pm 17$

Continuum RPA

Rodin & Urin
NPA 687, 276c

$103 \pm 9 \pm 25$

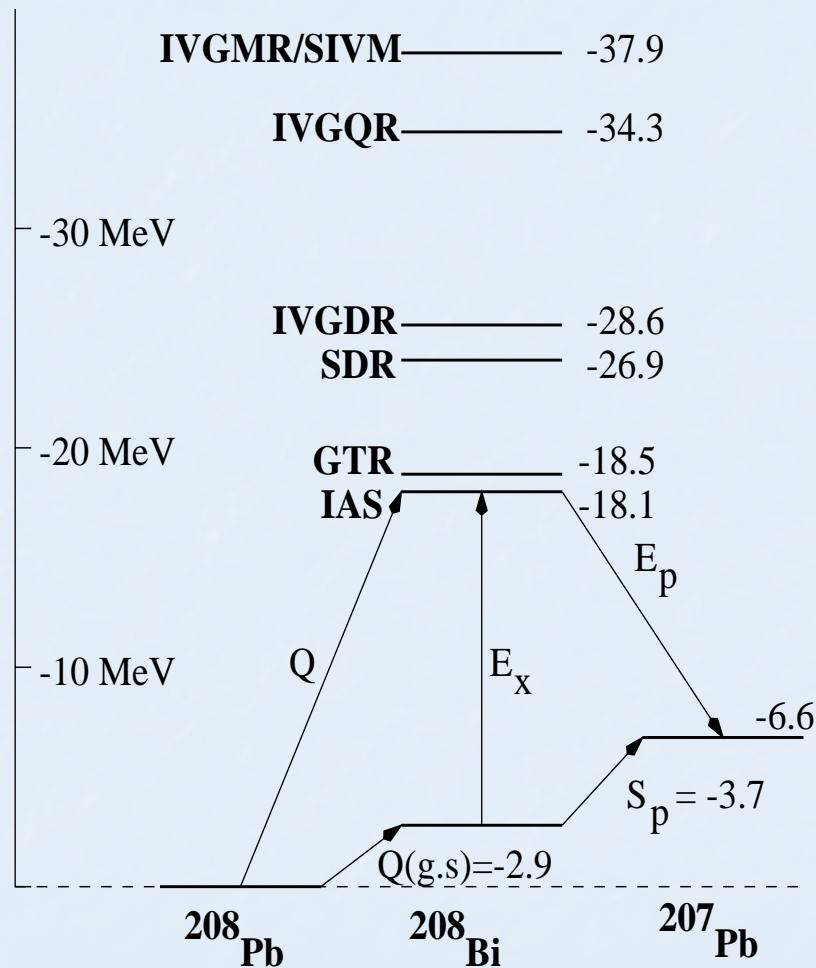
HF-RPA*

Auerbach & Klein
PRC 30, 1032

$210 \pm 16 \pm 45$

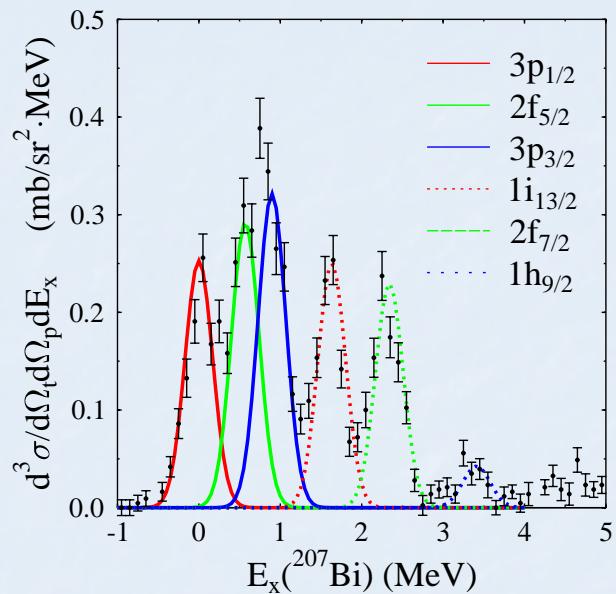
* Different operator, includes GT

Proton decay from the IVSGMR

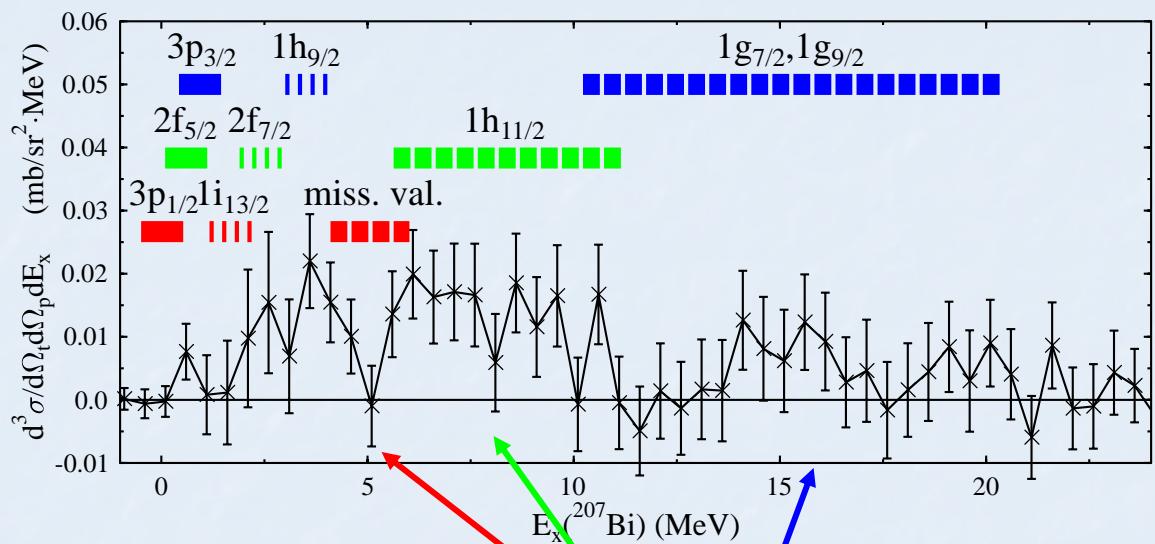


Final state spectra

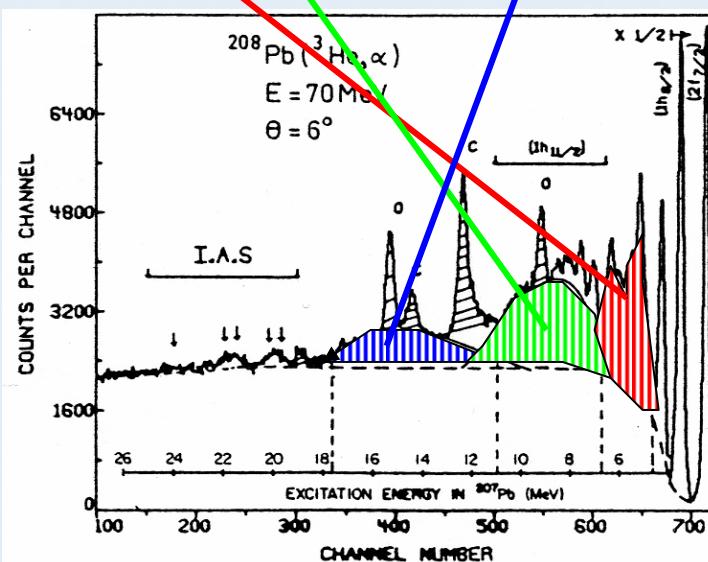
$E_x(^{208}\text{Bi}) < 30 \text{ MeV}$ (IAS+GTR+SDR+....)



IVSGMR ($30 < E_x(^{208}\text{Bi}) < 45 \text{ MeV}$ difference of angles)



Comparison with
 $^{208}\text{Pb}(^3\text{He}, \alpha)$
Galès *et al.*
Phys. Rep. 166, 255



Final state population in ^{207}Pb

Final state	Data(%)	Theory(%)*
$3\text{p}_{1/2} \ 2\text{f}_{5/2} \ 3\text{p}_{3/2}$	< 3	11.3
$1\text{i}_{13/2}$		21.4
$2\text{f}_{7/2} \ 1\text{h}_{9/2}$	13±5	9.5
$1\text{h}_{11/2}$	22±8	22.8
$1\text{g}_{7/2} \ 1\text{g}_{9/2}$	17±8	
All	52±12	66

*Rodin & Urin NPA 687, 276c (continuum RPA)

Large discrepancies for partial branchings!!

The (${}^3\text{He},\text{t}$) reaction at 0 degree.

- Cross sections at $E({}^3\text{He})=450 \text{ MeV}$, $\mathbf{q}=0$ for (${}^3\text{He},\text{t}$) reactions

$$\frac{d\sigma}{d\Omega} = \frac{\mu_i \mu_f}{(\pi \hbar^2)^2} \left(\frac{k_f}{k_i} \right) (N^D |J_\tau|^2 B(F) + N_{\sigma\tau}^D |J_{\sigma\tau}|^2 B(GT))$$

T. N. Taddeucci et al., Nucl. Phys. A469, 125 (1987)

I. Bergqvist et al., Nucl. Phys. A469, 648 (1987)

- Neutrino absorption cross sections

$$\sigma = \frac{1}{\pi \hbar^4 c^3} [G_V^2 B(F) + G_A^2 B(GT)] \times F(Z, E_e) p_e E_e$$

Importance of charge-exchange reactions at intermediate energies



Measuring GT strengths

$$\frac{d\sigma}{d\Omega}(q=0) = KN_D |J_{\sigma\tau}|^2 B(GT)$$

kinematic factor

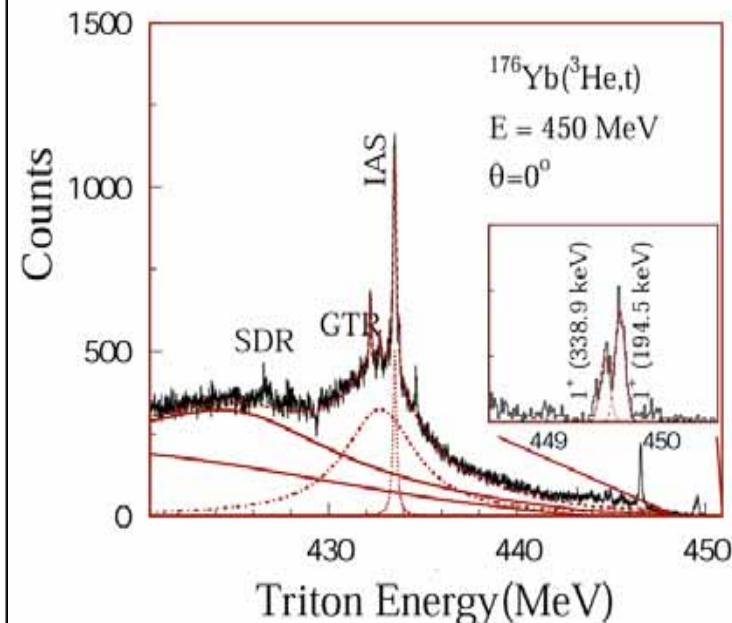
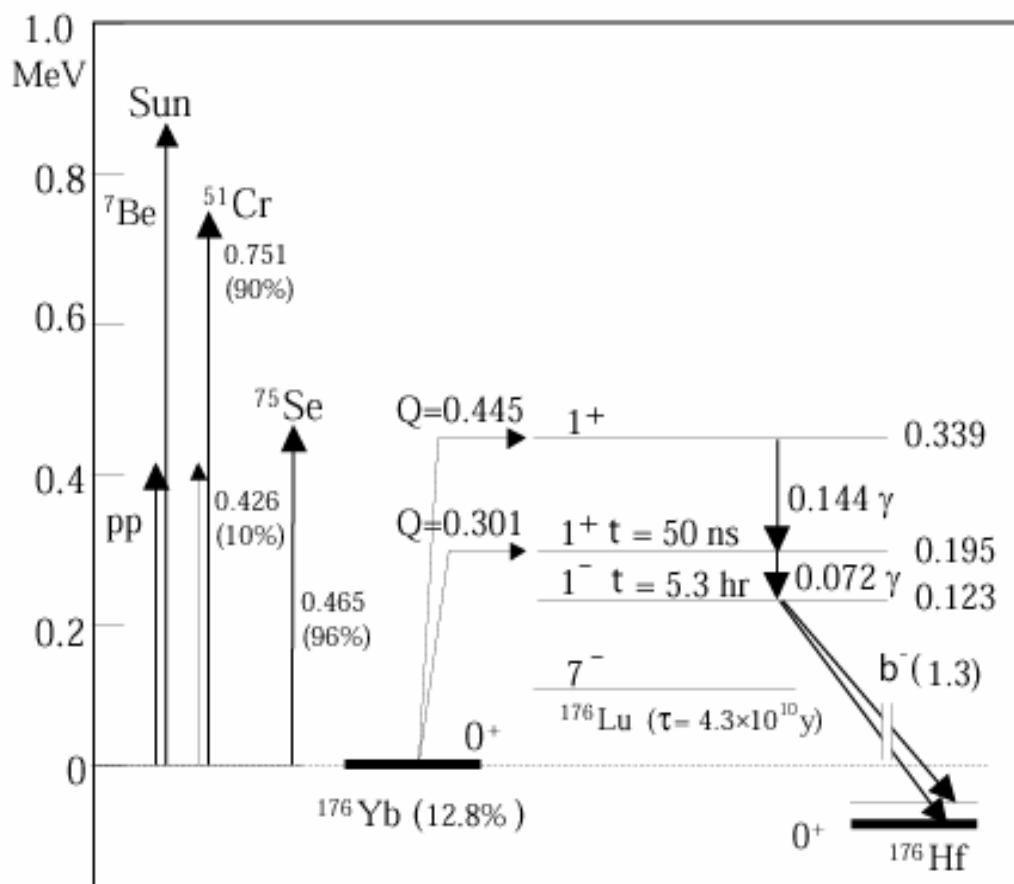
distortion factor

Gamow-Teller strength

nucleon-nucleus interaction

Calibration of $B(GT)$ to cross section for known transitions
(e.g. from β -decay)





$E_x \text{ (MeV)}$	$0.195 + 0.339 \text{ (} p, n \text{)}$	0.195	$0.339 \text{ (} ^3\text{He}, t \text{)}$
$B(\text{GT})$	0.32 ± 0.04	0.20 ± 0.04	0.11 ± 0.02

Why are Gamow-Teller transitions in *pf*-shell nuclei important ?

- **Role of *pf*-shell nuclei in supernova explosions:**
Core of supernova star is composed of *pf*-shell nuclei.

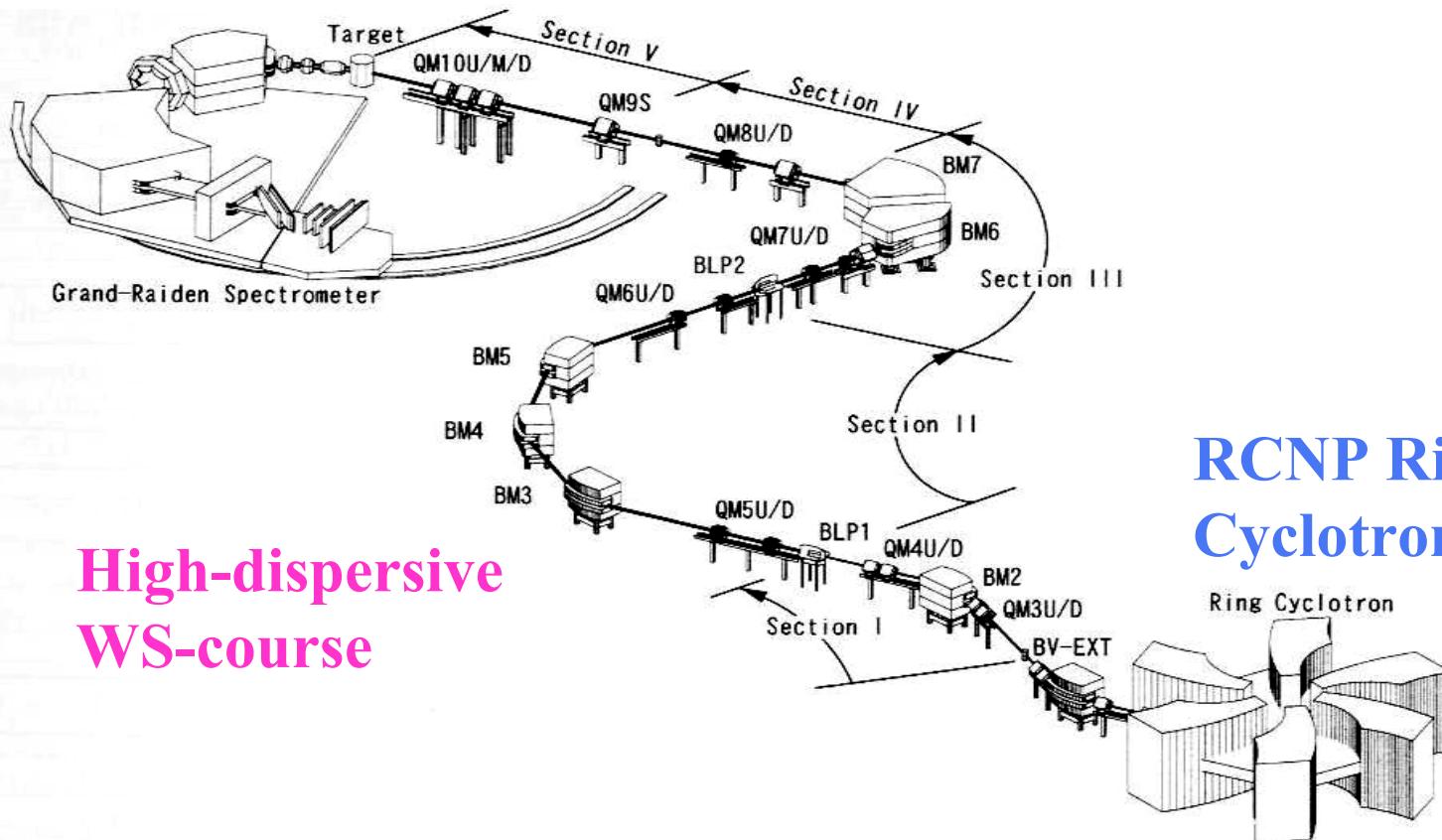
Neutrino absorption cross sections by *pf*-shell nuclei are essential in understanding of nuclear synthesis by Supernova explosions in cosmos.

- **Difficulties in shell model calculations for *pf*-shell nuclei.**
- **Importance of spin-isospin responses of *pf*-shell nuclei**

Beam line WS-course

Grand-Raiden
Spectrometer

T. Wakasa et al., NIM A482 ('02) 79.



RCNP Ring
Cyclotron

Ring Cyclotron

High-dispersive
WS-course

IUCF

Evolution of Resolution
in Charge-Exchange Reactions
at Intermediate Energies

$^{58}\text{Ni}(\text{p},\text{n})$
 $E_{\text{p}}=160\text{MeV}$, 0-deg., IUCF

J.Rapaport et al.,
Nucl.Phys. A410 (1983) 371.

$\Delta E \sim 400\text{keV}$

RCNP

before
WS

WS

$^{58}\text{Ni}(^3\text{He},t)^{58}\text{Cu}$
 $E(^3\text{He})=450\text{ MeV}, t=0$

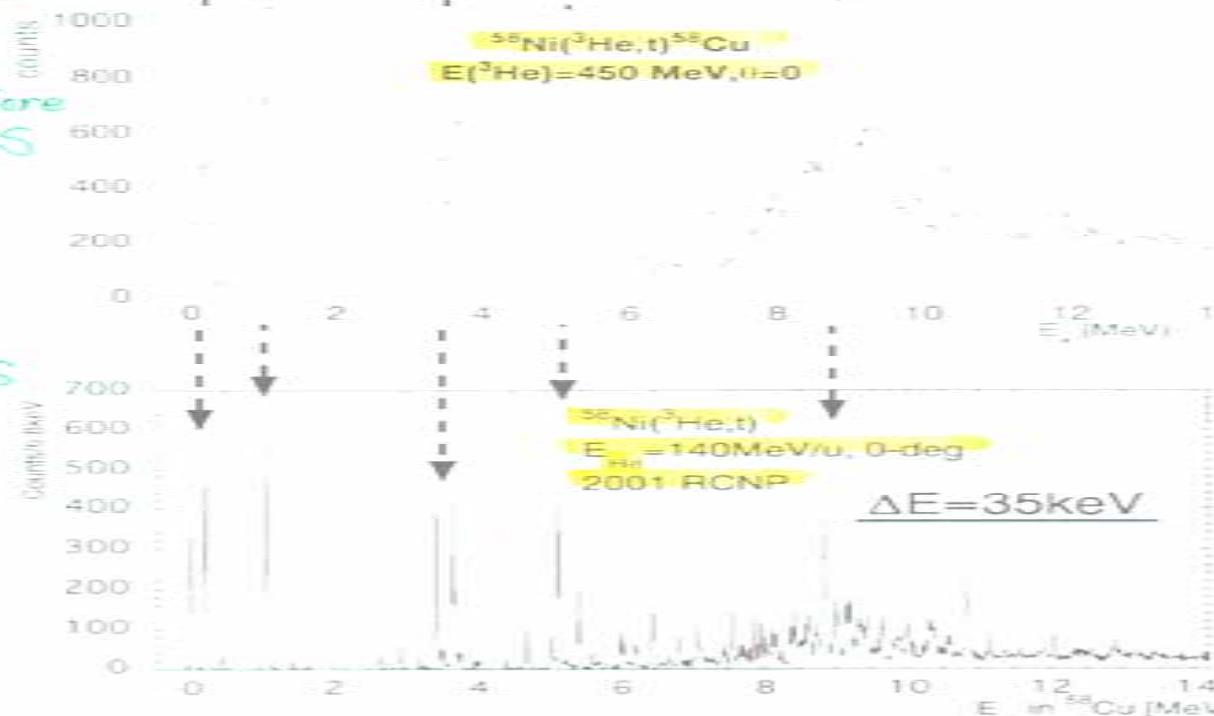
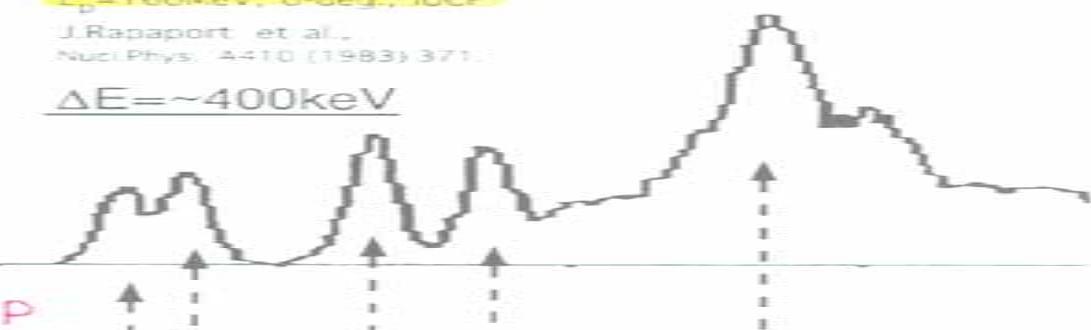
Y.Fujita et al.
Phys.Lett.B365
(1996) 29

$^{58}\text{Ni}(^3\text{He},t)$
 $E(^3\text{He})=140\text{MeV/u, 0-deg}$
2001 RCNP

$\Delta E = 35\text{keV}$

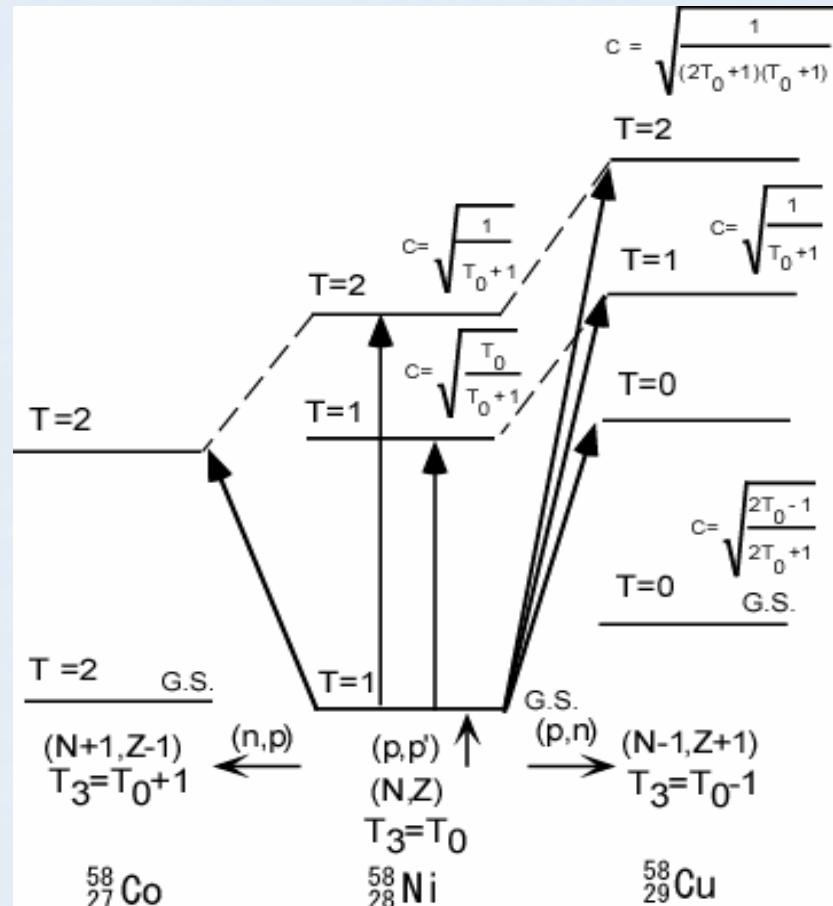
H.Fujita et al.
PhD thesis

Y.Fujita et al.
Euro. Phys. J. A
13 (2002) 411
($E_x \leq 8\text{ MeV}$)



Decomposition of the isospin component of the excited state in ^{58}Cu .

- Isospin of ^{58}Ni g.s. : $T_0=1$
- In principle, comparison among (n,p), (p,p), (p,n) spectra
→ separates isospin components
But, very difficult in practice because of high level density for T=1 and T=2 states.
- CG
- $\sigma_{T=0} : \sigma_{T=1} : \sigma_{T=2} = 2 : 3 : 1$
($T_0=1$)

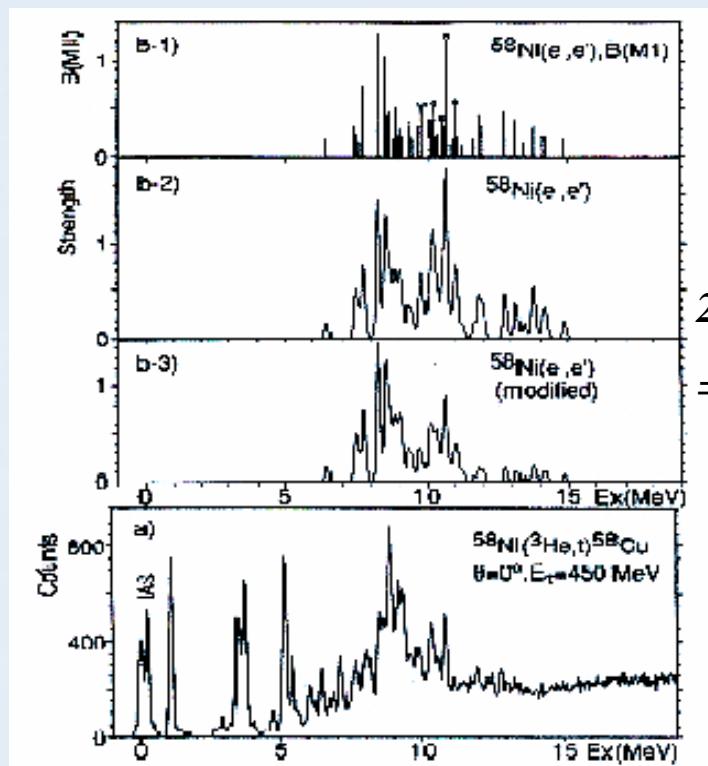


Comparison of ($^3\text{He},\text{t}$) and (e,e') spectra

- Comparison of ($^3\text{He},\text{t}$) with (e,e') spectra → Try to separate isospin components
- At $E_x = 6\text{-}10 \text{ MeV}$ ($T=1$ region)
 - Rather good correspondence
- At $E_x = 10\text{-}15 \text{ MeV}$ ($T=2$ region)
 - No good correspondence

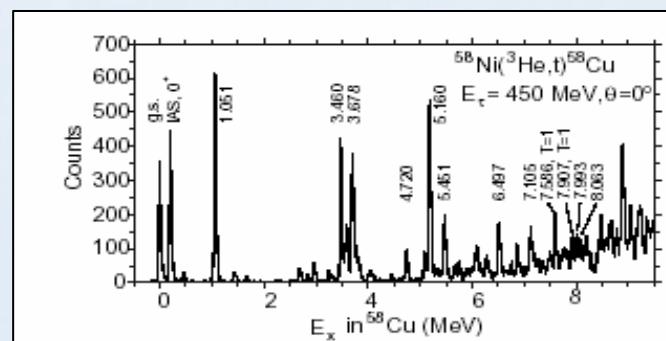
Fujita et al., Phys. Lett. B 365, 29 (1996).

Fujita et al., Eur. Phys. J A13, 411 (2002).

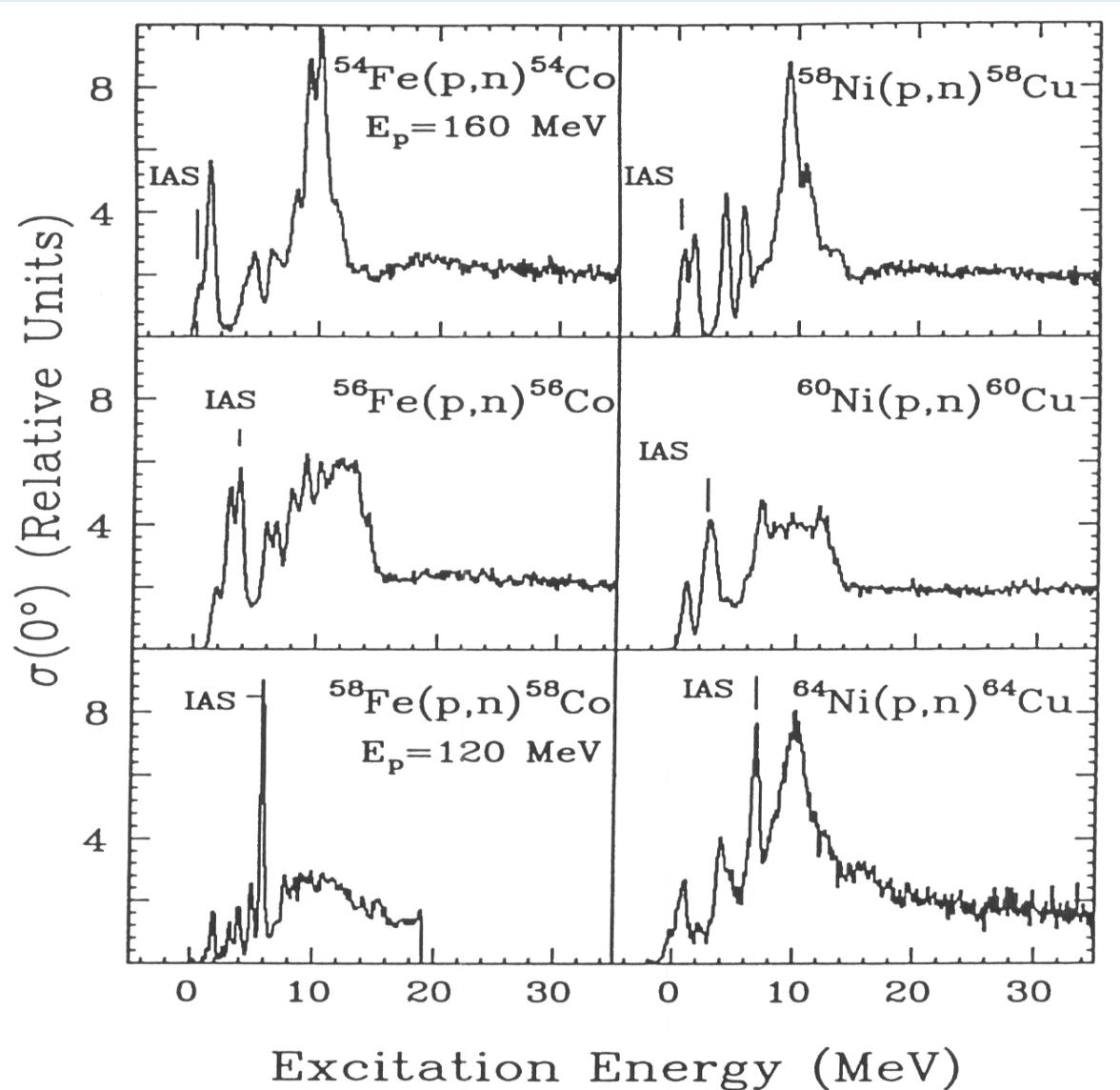


$$\sigma(T=1) : \sigma(T=2) = 1 : 1$$

$$\sigma(T=1) : \sigma(T=2) = 3 : 1$$

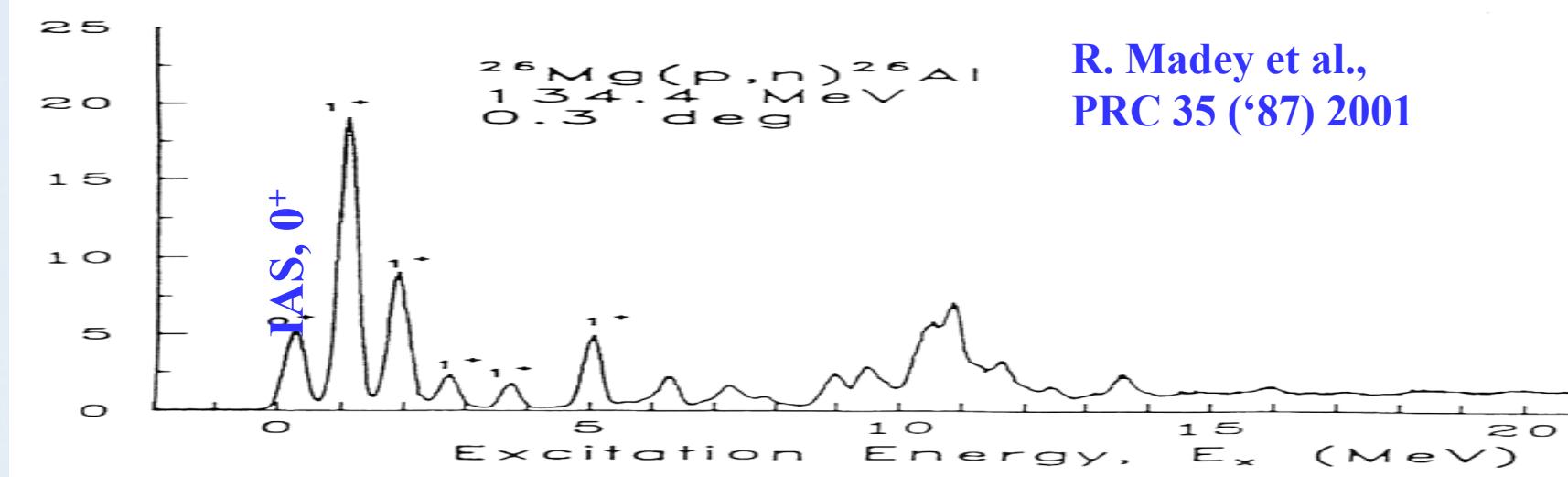


(p, n) spectra for Fe and Ni Isotopes

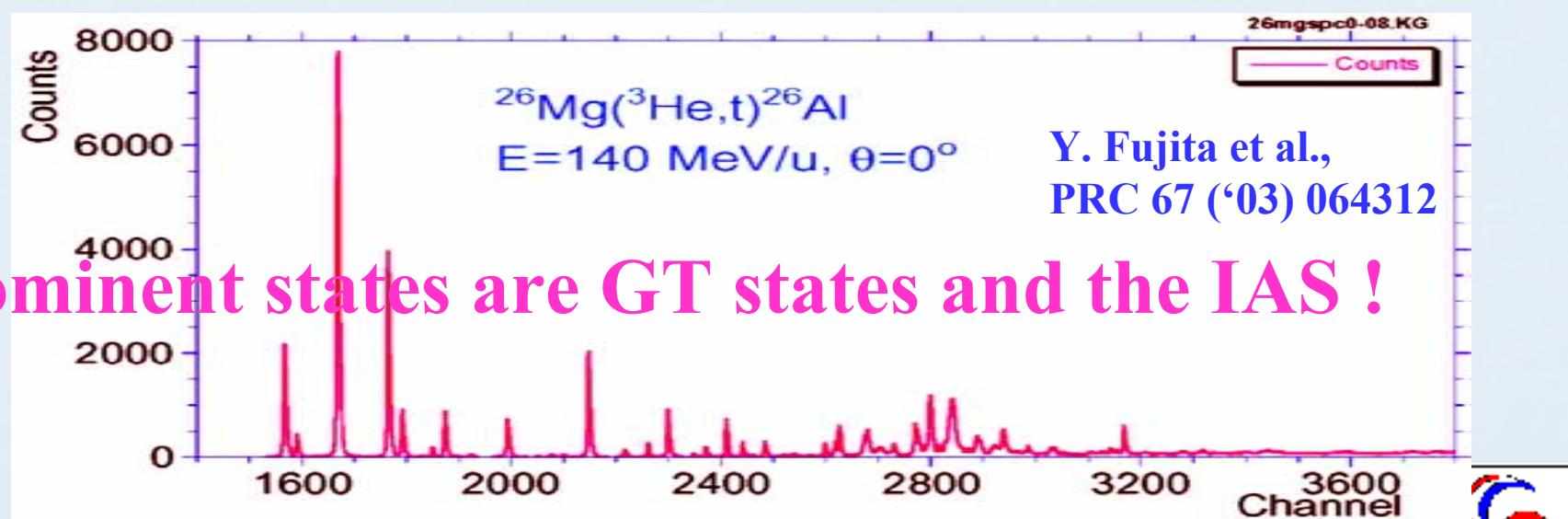


Rapaport
&
Sugarbaker
Rev. Mod. Phys.
('94)

$^{26}\text{Mg}(\text{p}, \text{n})^{26}\text{Al}$ & $^{26}\text{Mg}({}^3\text{He}, \text{t})^{26}\text{Al}$ spectra

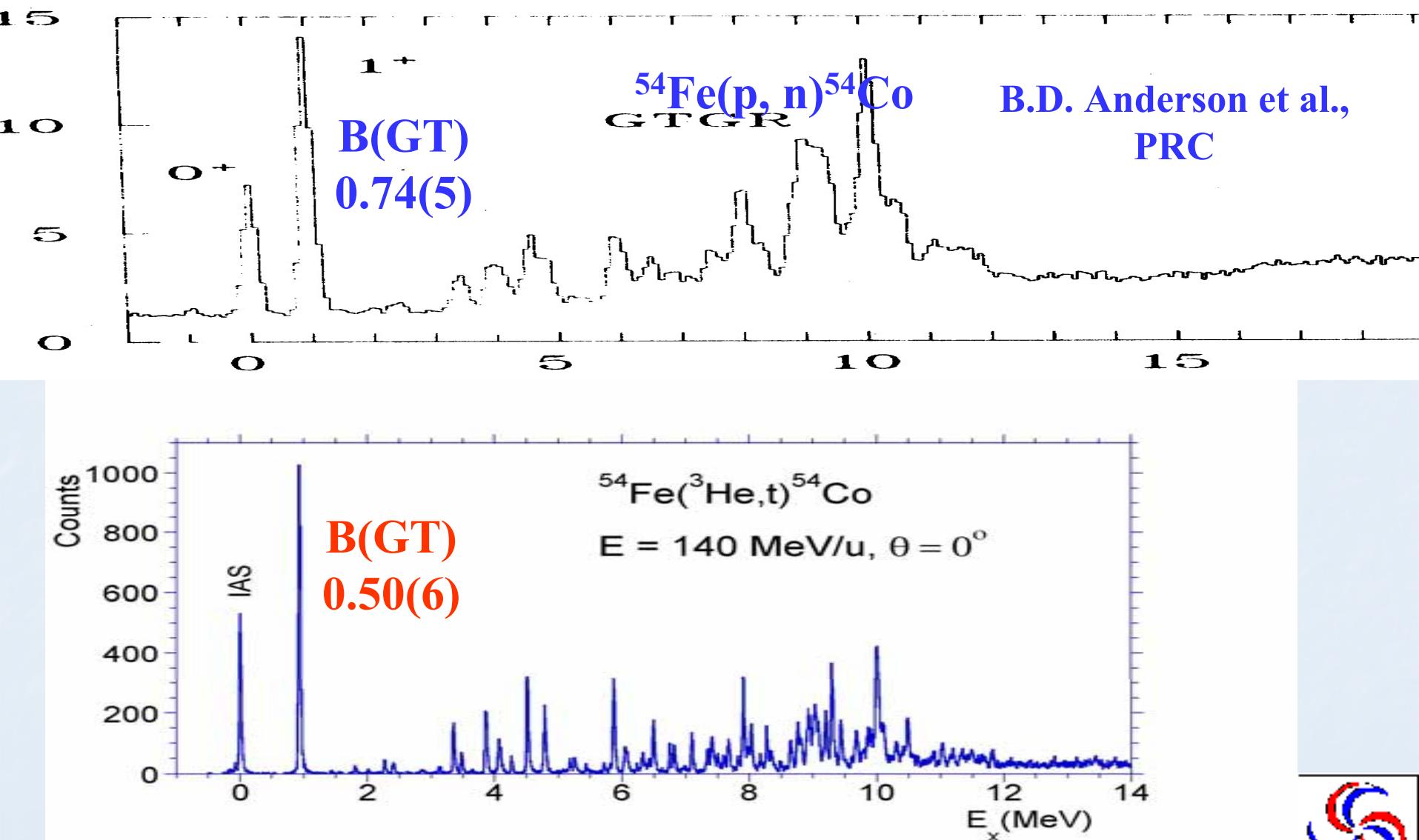


R. Madey et al.,
PRC 35 ('87) 2001



Prominent states are GT states and the IAS !

$^{54}\text{Fe}(\text{p},\text{n})$ & $^{54}\text{Fe}({}^3\text{He},\text{t})$

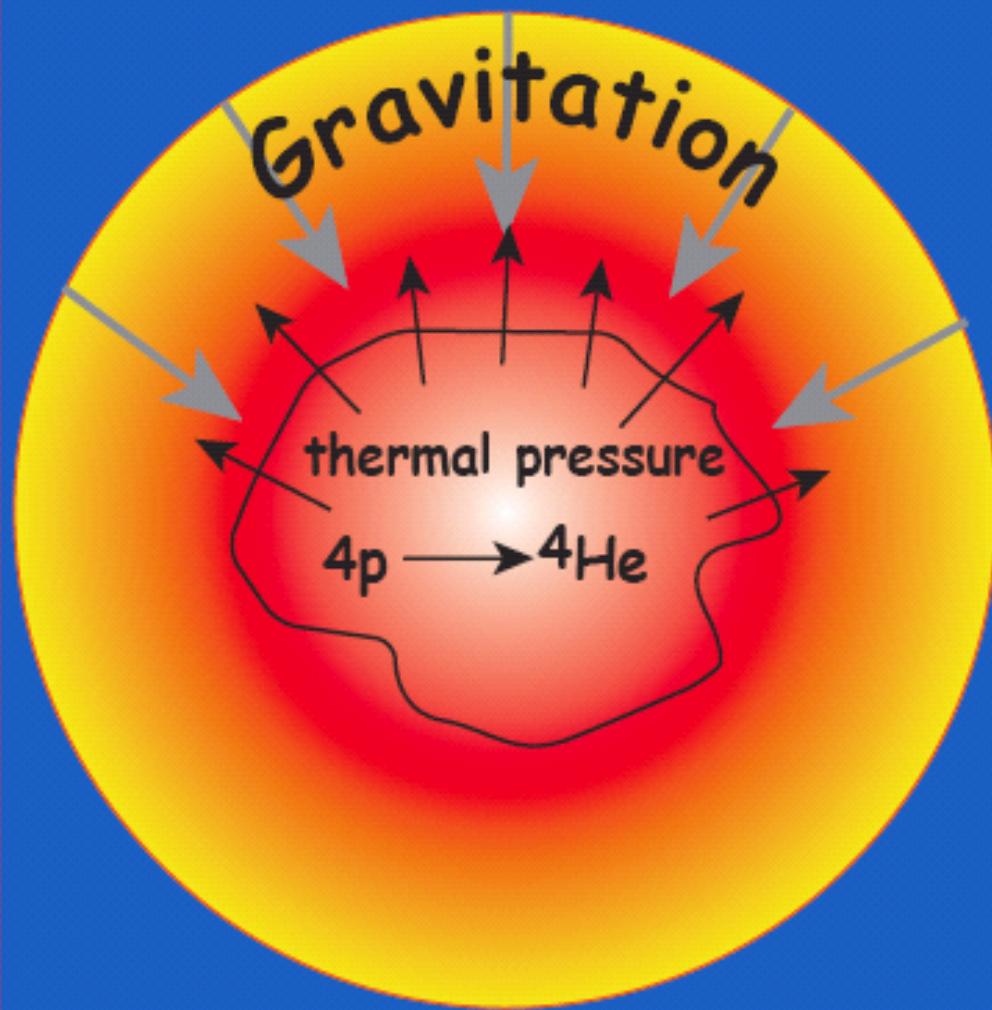


Determination of GT⁺ Strength and its Astrophysical Implications



Supernovae
Cassiopeia A
Chandra

Nuclear processes and energy household of supernovae



initial condition:

$$M > 10 M_{\odot}$$

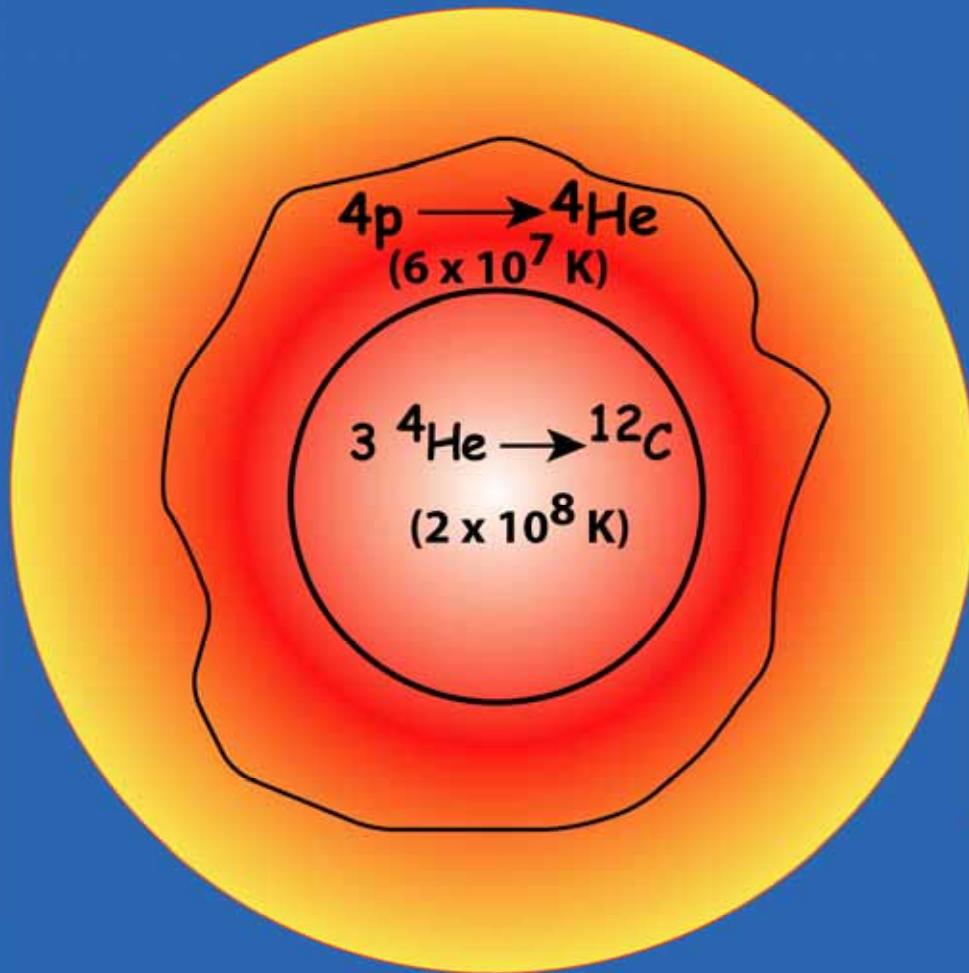
energy:



at: $T \sim 10^7 - 10^8 \text{ K}$

lifetime: $10^6 - 10^7 \text{ y}$

after 10^6 - 10^7 y



end of H-burning

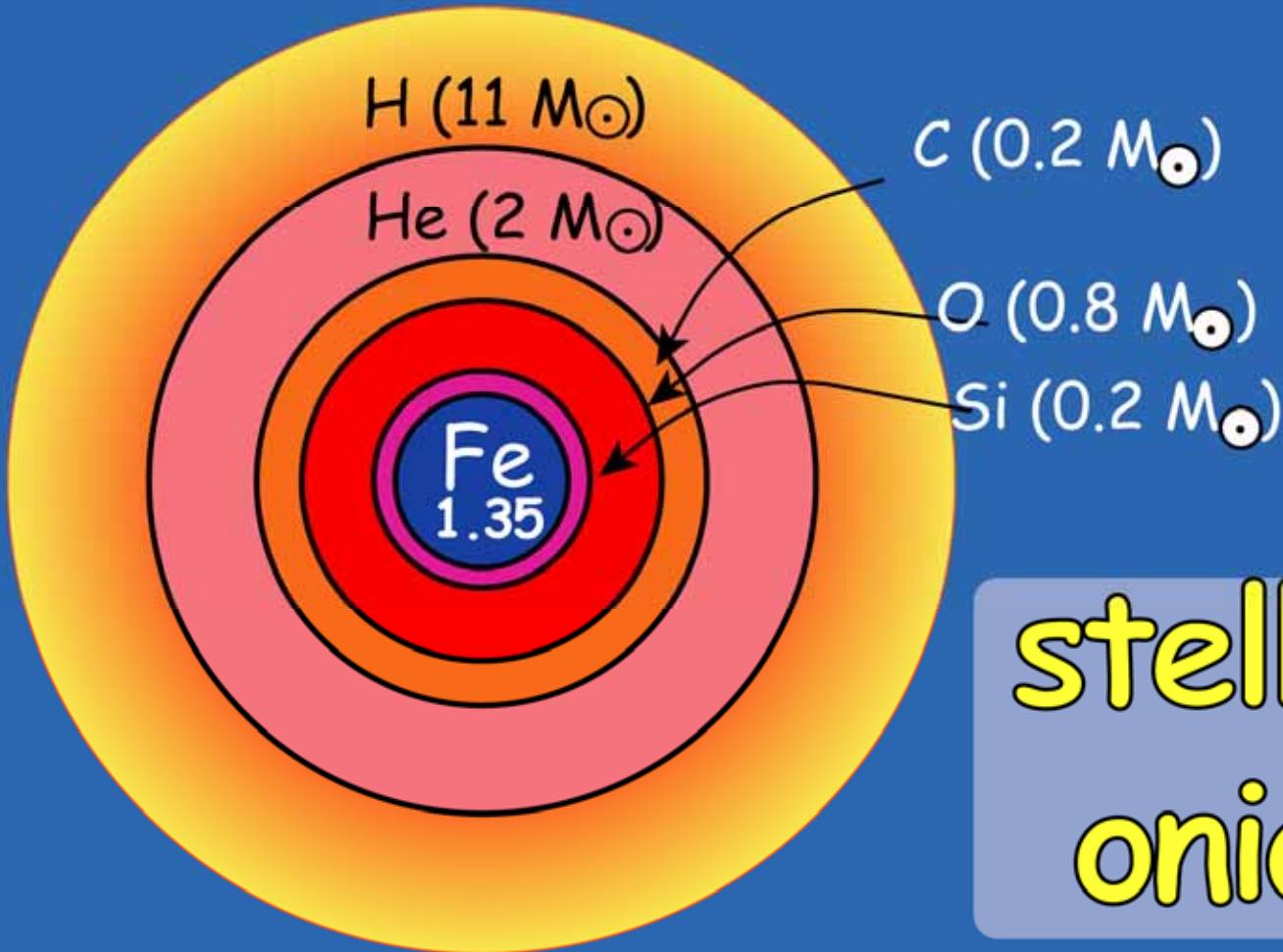
contraction of star

temperature increase

Red Giant (Super-Giant)

lifetime: 5×10^5 y

end of stellar evolution $M_{\text{star}} \sim 15 M_{\odot}$

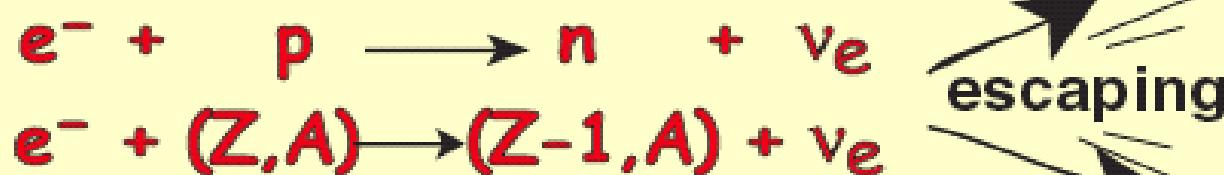


stellar
onion

SN-explosion scenario

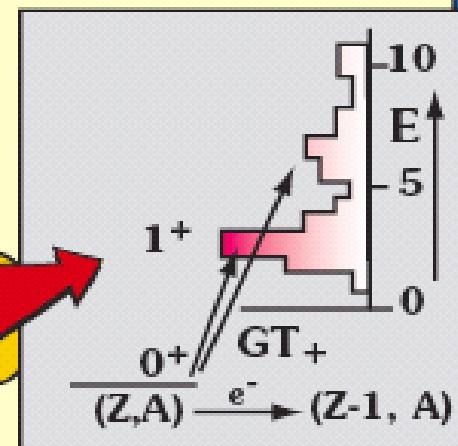
- interior of star gets enriched with Fe
- gravitational pressure increases
balanced by degenerate electron gas up to
Chandrasekhar limit: $M_{\text{ch}} = 1.44 (2Y_e)^2 M_{\odot}$

- start of collapse at
 $T = 10^9 \text{ K}$ and $\rho = 3 \times 10^7 \text{ g/cm}^3$
accelerated by neutronization (de-leptonization)



- loss of pressure
- accelerated collapse
- reduction of Y_e
and loss of energy!!

rate determined
by GT-strength
($\Delta S = 1, \Delta T = 1, \Delta L = 0$)



Y_e at freeze-out determines the explosive energy!!

Electron capture in *fp*-shell

- In supernova explosions, electron capture (EC) on *fp*-shell nuclei plays a dominant role during the last few days of a heavy star [presupernova stage; deleptonization \Rightarrow core collapse \Rightarrow subsequent type IIa Supernova (SN) explosion]
Bethe *et al.* (1979)
- The rate for EC is governed by the GT⁺ strength distribution at low excitation energy; not accessible to β -decay.
Fuller, Fowler and Newman (FFN) (1982-1985); estimates of stellar rates in stellar environments using s.p. model.
Caurier *et al.*, Martínez-Pinedo & Langanke (1999), Otsuka *et al.* \Rightarrow Large shell-model calculations \Rightarrow marked deviations from FFN EC rate; generally smaller EC rates.
- Experiments and theory relied on (*n,p*) data (TRIUMF) which have a rather poor energy resolution

Stellar EC rates

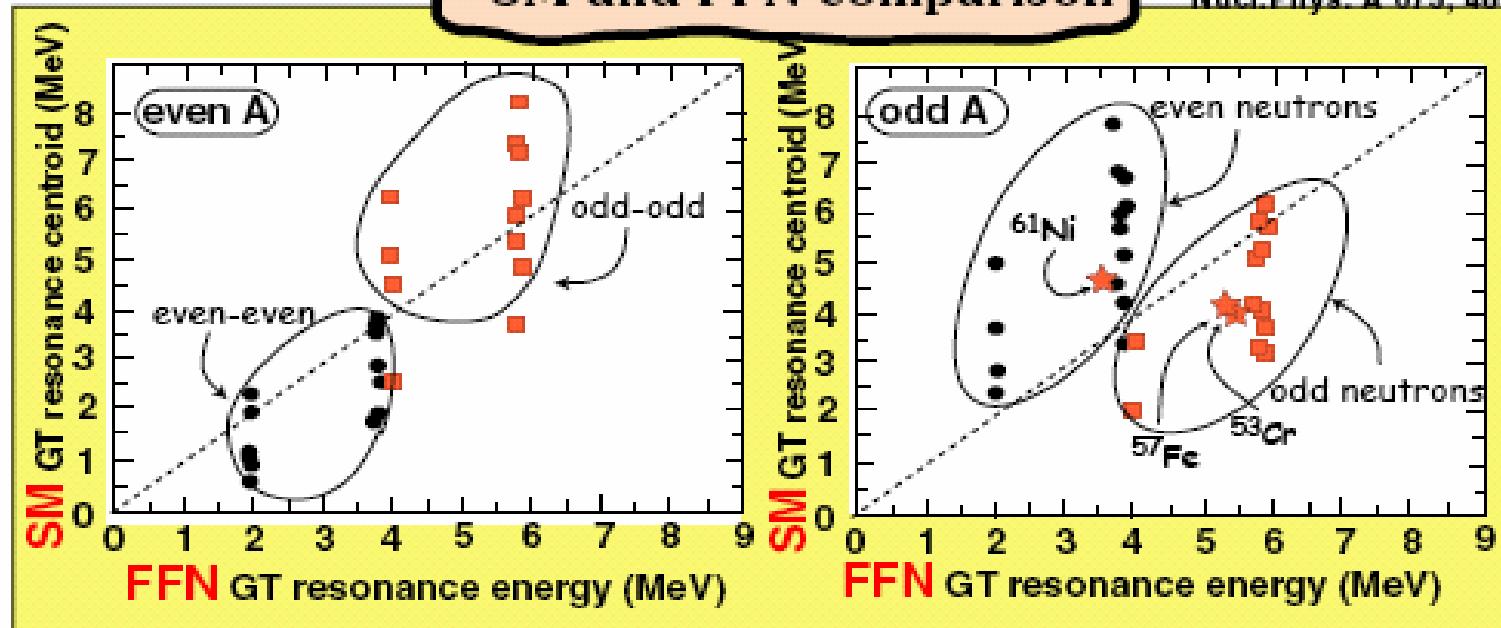
at a given temperature

Important parameters

- location of GT resonance (most important)
- level of quenching
- fragmentation over excitation energy

SM and FFN comparison

from:
Langanke, Martinez-Pinedo
Nucl.Phys. A 673, 481 (2000)



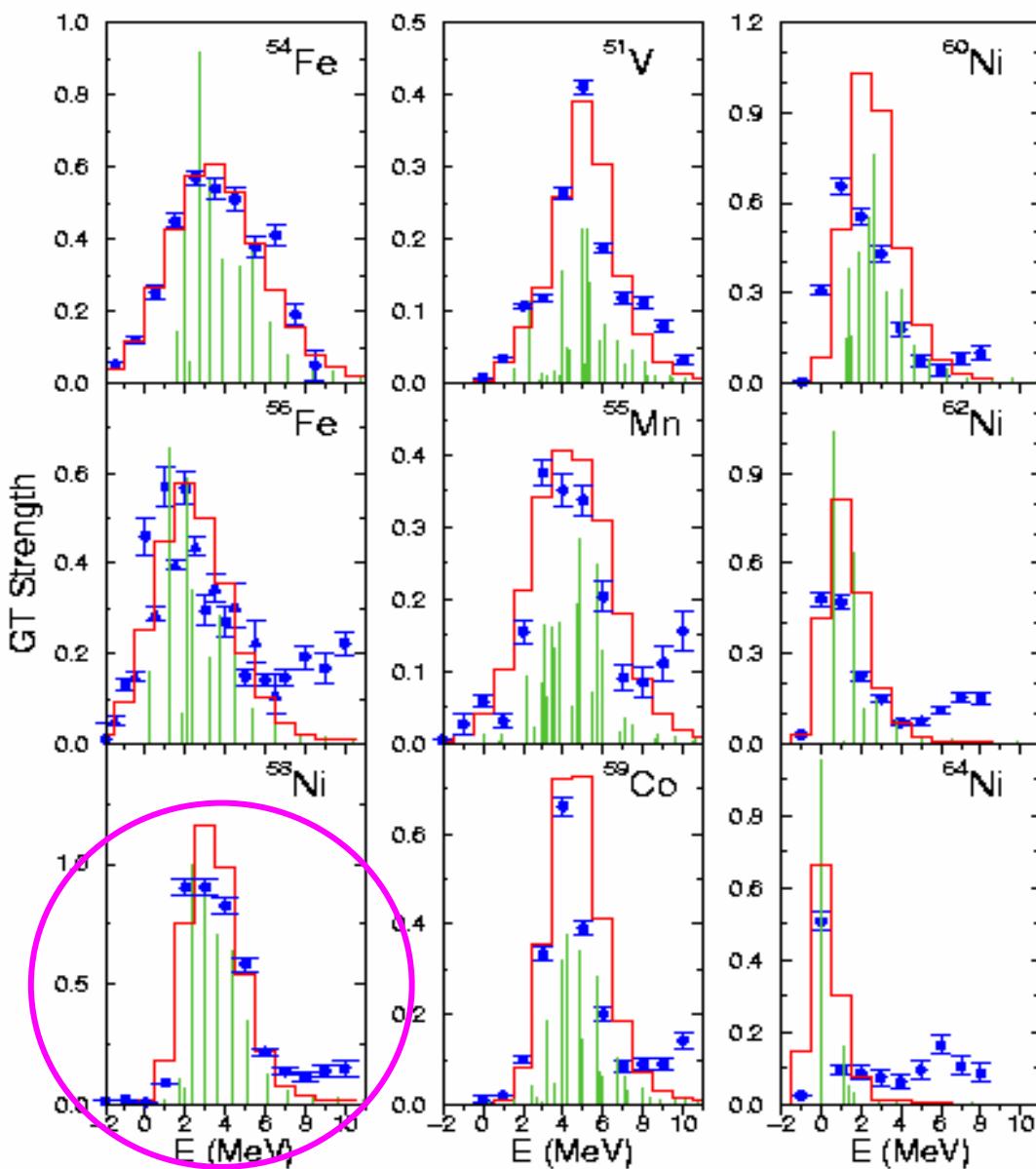
- no exp. data on odd-N nuclei (usually rare!!)
- no exp. data on odd-odd nuclei (^{50}V is the only stable one)

fp-shell nuclei: large scale shell model calculations

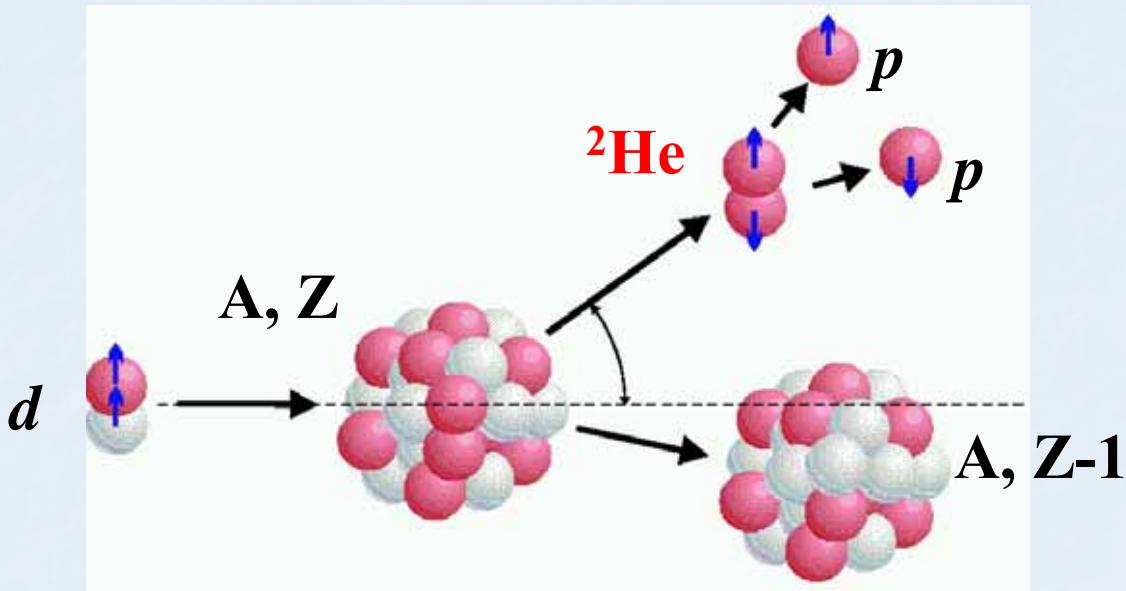
E. Caurier *et al.*
NPA 653 (1999) 439

- Stellar weak reaction rates with improved reliability
- Large scale shell model (SM) calculations
- Tuned to reproduce GT⁺ strength measured in (n,p)
- (n,p) data from TRIUMF
- GT⁺ strength from SM
- Folded with energy resolution

Case study: ^{58}Ni



Exclusive excitations $\Delta S = \Delta T = 1$: (d , ${}^2\text{He}$)

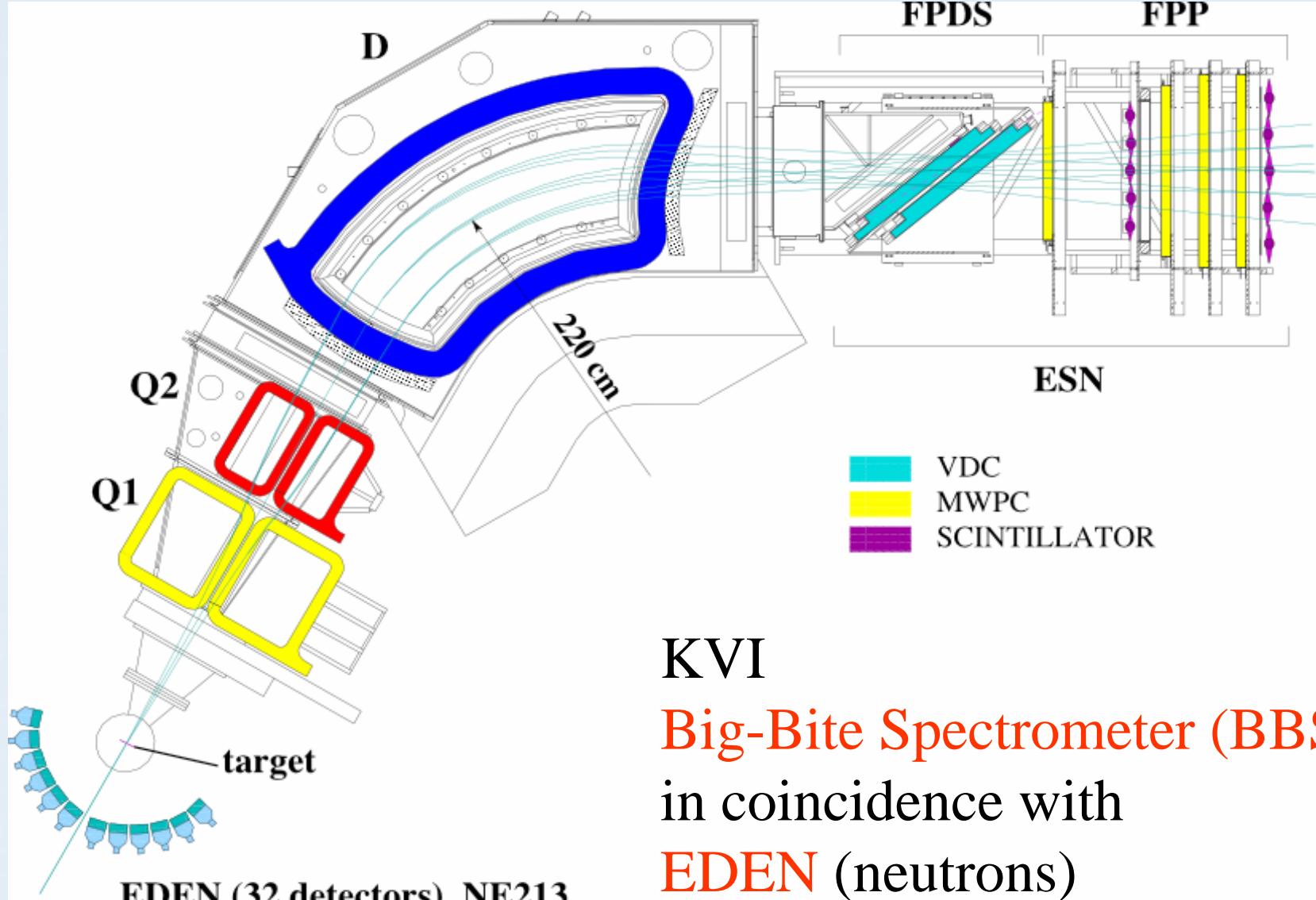


${}^3\text{S}_1$ deuteron $\Rightarrow {}^1\text{S}_0$ di-proton (${}^2\text{He}$)

${}^1\text{S}_0$ dominates if (relative) proton kinetic energy $\varepsilon < 1$ MeV

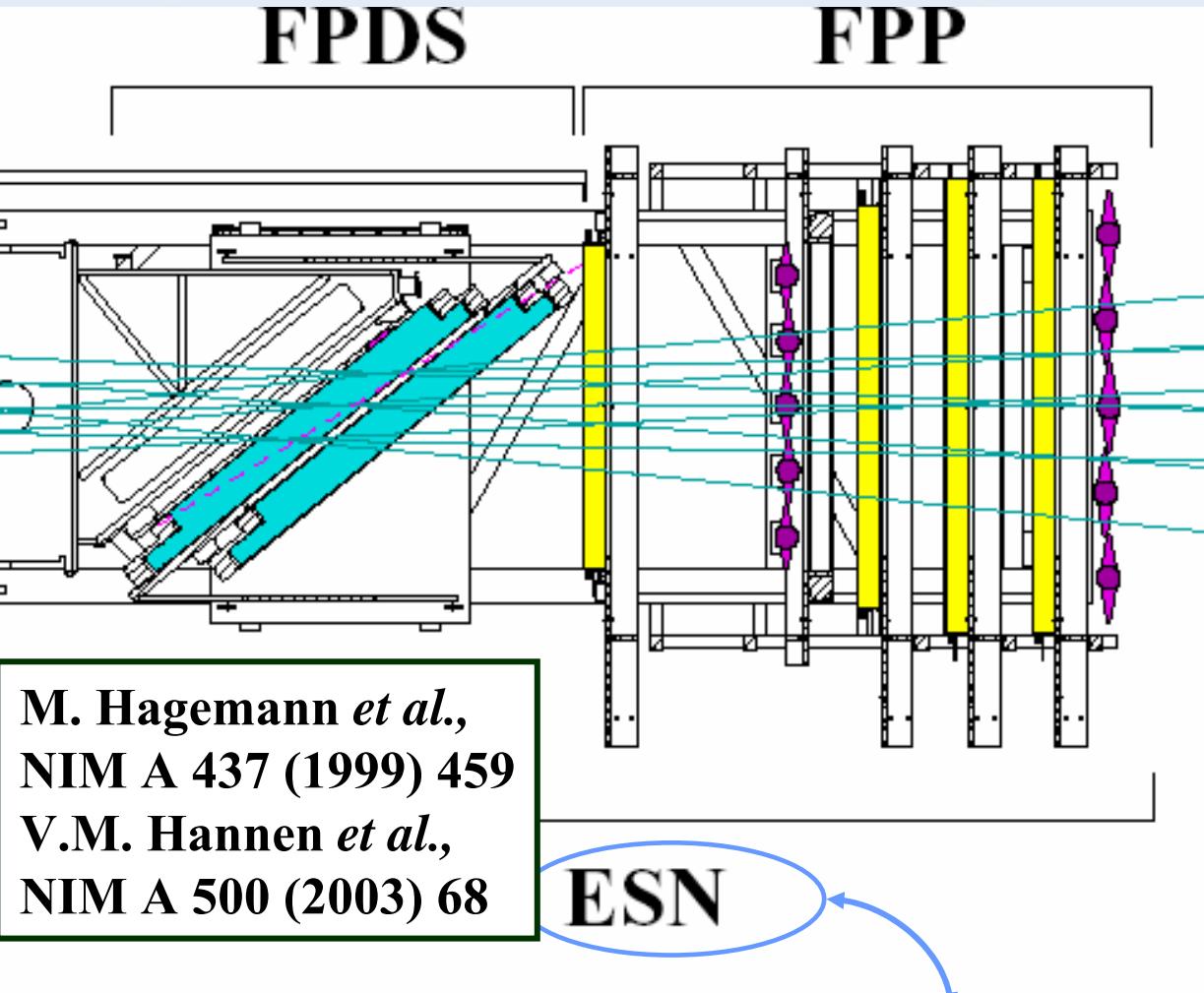
(n,p)-type probe with exclusive $\Delta S=1$ character (GT⁺ transitions)

But near 0°: tremendous background from d -breakup



KVI
Big-Bite Spectrometer (BBS)
in coincidence with
EDEN (neutrons)
SiLi ball (protons)

Setup: ESN detector



Focal-Plane Detector:
(FPDS): 2 VDCs

Focal-Plane Polarimeter:
**(FPP): 4 MWPCs &
graphite analyzer**

features a.o.:
fast readout
VDC readout pipeline
TDC's
**VDC decoding using
imaging techniques**
DSP based online analysis

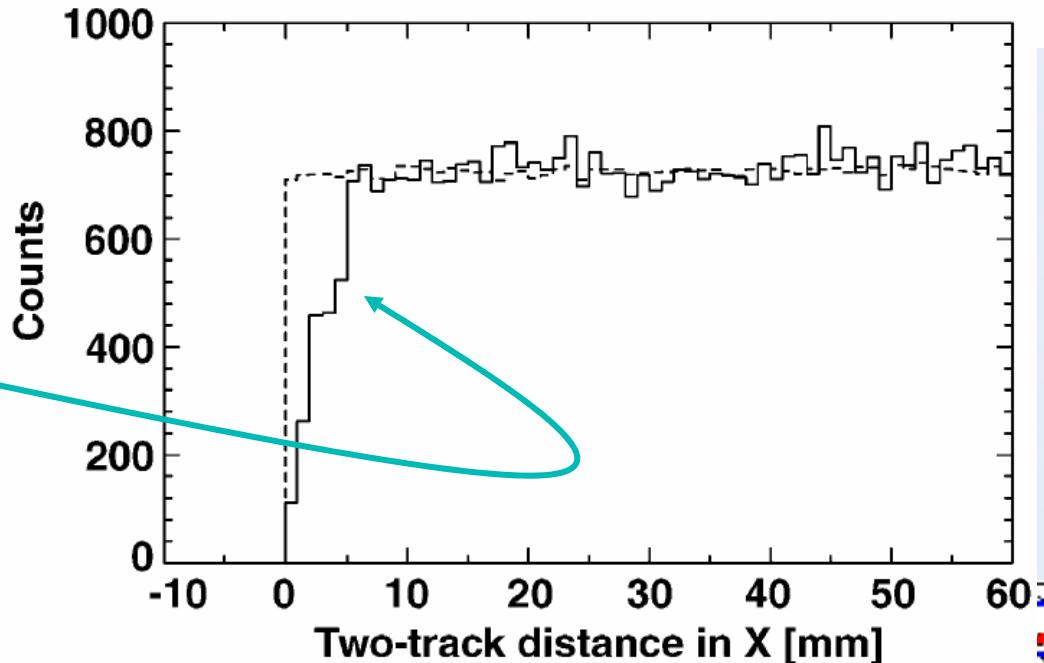
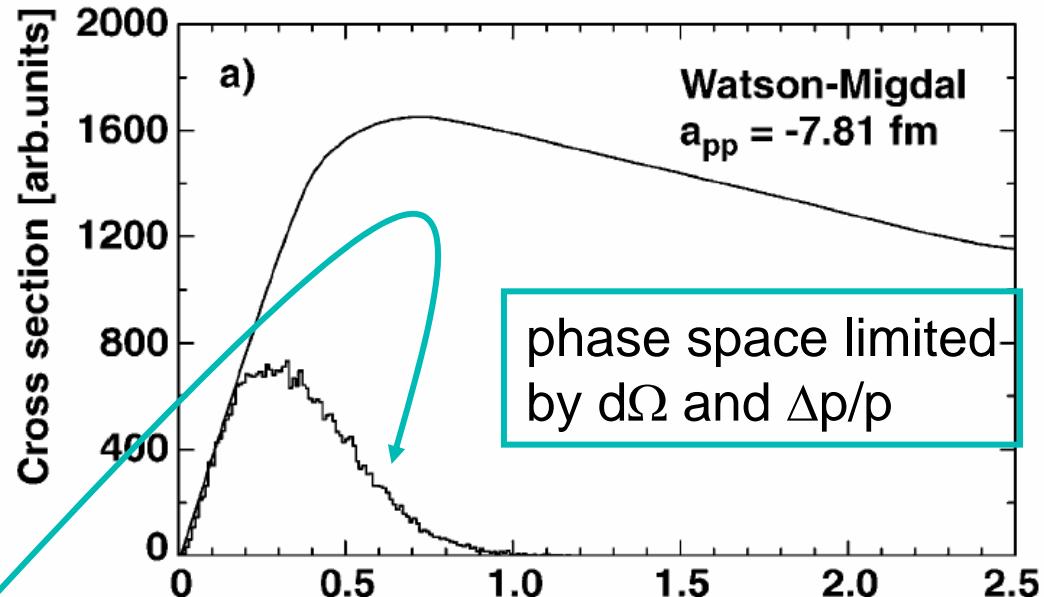
Bari, Darmstadt, Gent, Iserlohn, KVI, Milano, Münster, TRIUMF

- Good double tracking
- Use VDC information
- Good phase space coverage for small relative proton energies

S. Rakers *et al.*

NIM A481 (2002) 253

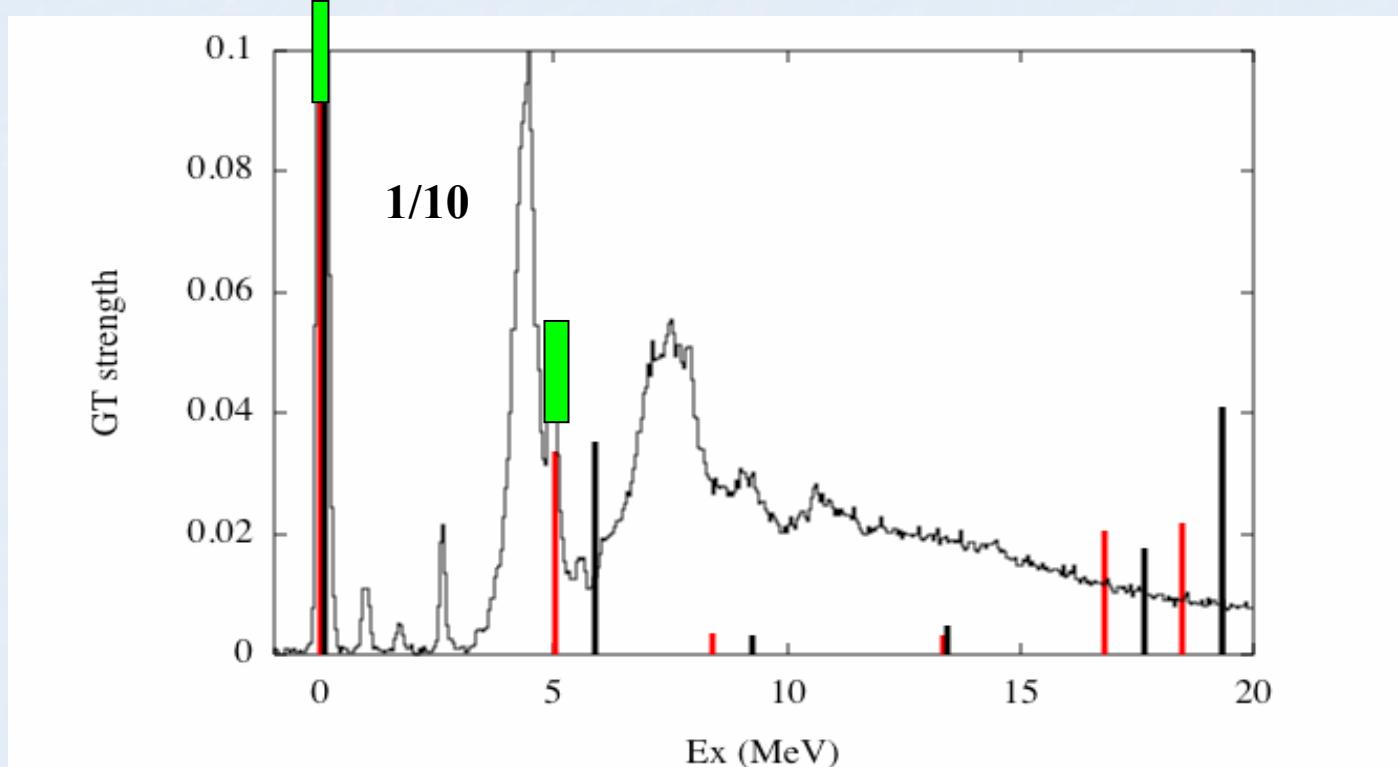
measured



Exclusive measurement of $\Delta S = \Delta T = 1$ strength:

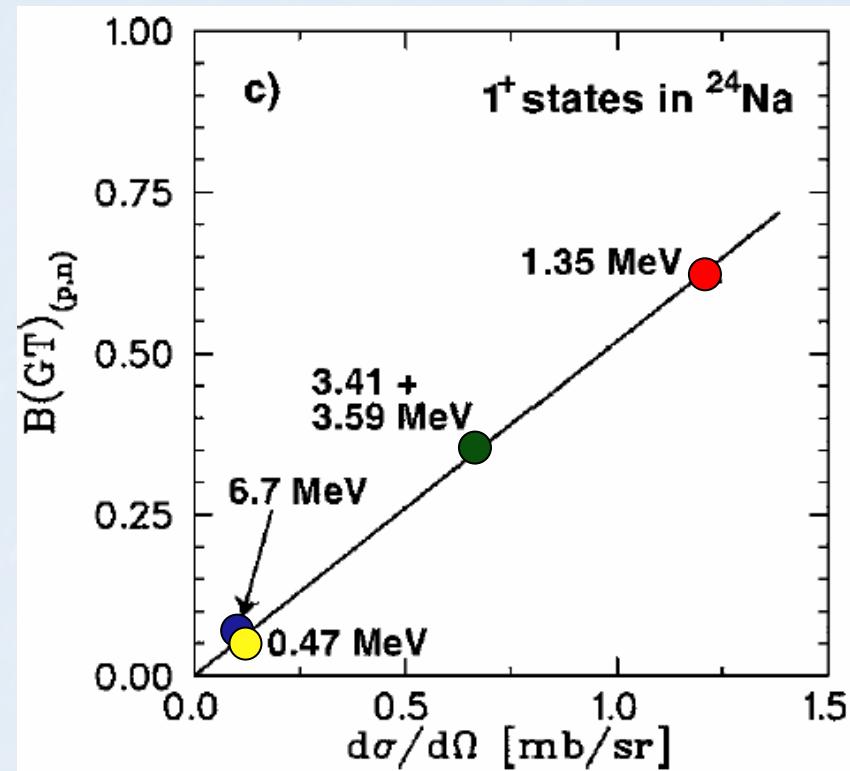
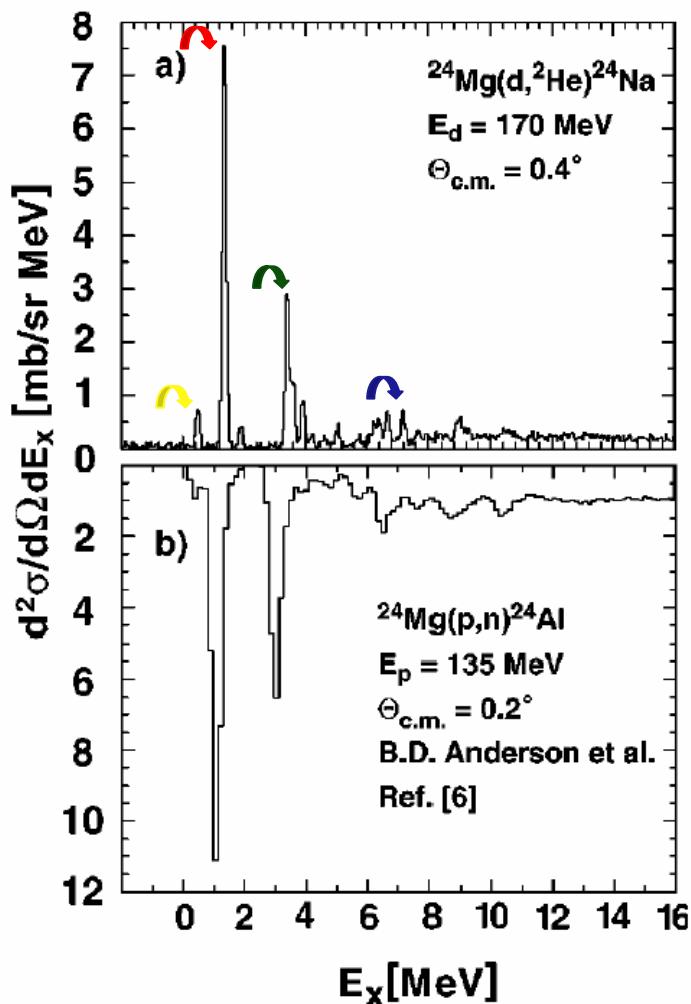


$E_0 = 171 \text{ MeV}, \theta = 0^\circ$



- shell model calculations $4 \hbar\omega$ & $6 \hbar\omega$ (G. Martinez-Pinedo)
- B (GT^+) (S. Rakers) ■■■

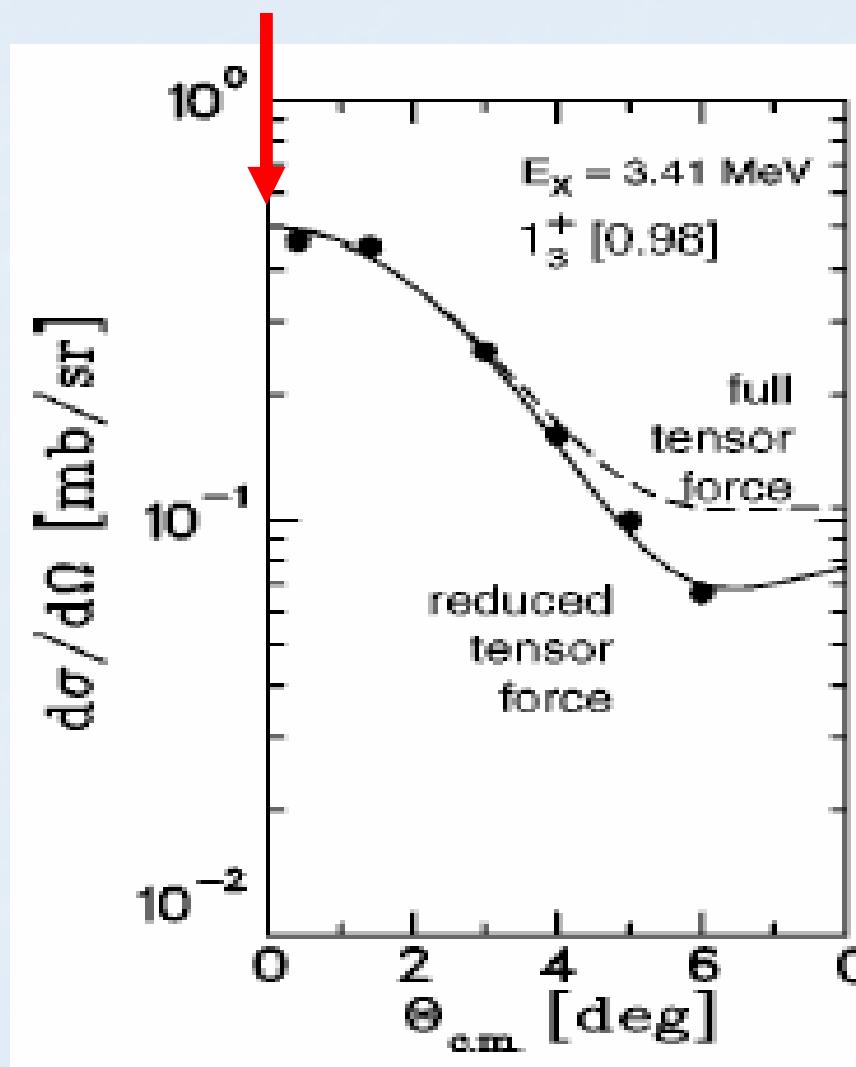
(p,n) vs $(d,{}^2\text{He})$: calibration



Self-conjugate ${}^{24}\text{Mg}$

S. Rakers *et al.*
PRC 65 (2002) 044323

Experimental cross section and GT strength



$$B_{\text{exp}}(\text{GT}+) =$$

$$\frac{d\sigma(q=0)}{d\Omega} \cdot \left[\frac{d\sigma(GT)}{d\Omega} \right]^{-1}$$

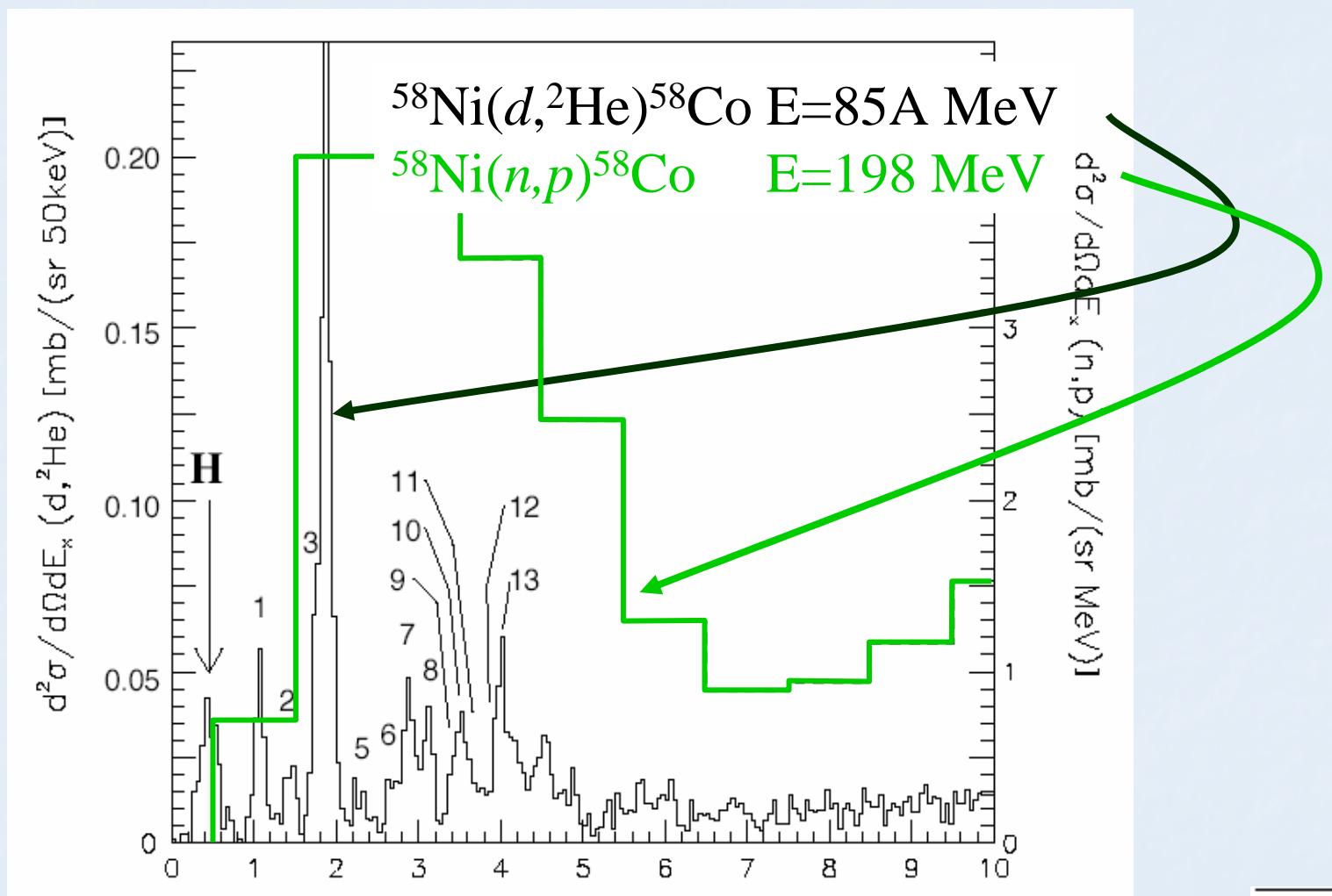
extrapolated
(DWBA)

unit cross section

GT Strength in ^{12}B and ^{24}Na from (d, ^2He) reaction

Target	Reference data				Present data		
	E_x	$B(\text{GT}_-)$	E_x	$d\sigma/d\Omega(q=0)$	$\sigma(L=0)/\sigma(\tau_0\tau)$	$B(\text{GT}_+)$	
	[MeV]		[MeV]	[mb/sr]	(q=0)		(C=0.267)
^{12}B	0.00	0.998	0.00	2.580 ± 0.138	0.988	0.930 ± 0.050	
			5.00	0.138 ± 0.010	0.976	0.050 ± 0.004	
^{24}Na	0.44	0.050	0.47	0.138 ± 0.012	0.821	0.049 ± 0.004	
	1.07	0.613	1.35	1.563 ± 0.085	0.948	0.654 ± 0.035	
	1.58	0.020	1.89	0.087 ± 0.026	0.649	0.025 ± 0.008	
	2.98	0.362	3.41	0.667 ± 0.039	0.980	0.290 ± 0.016	
			3.59	0.266 ± 0.018	0.806	0.095 ± 0.006	
	3.33	0.059	3.92	0.193 ± 0.058	0.809	0.070 ± 0.022	
	4.69	0.015	5.06	0.093 ± 0.027	0.561	0.024 ± 0.007	
			6.24	0.086 ± 0.026	0.818	0.031 ± 0.010	
	6.46	0.068	6.70	0.161 ± 0.012	0.972	0.071 ± 0.005	
	6.87	0.029	7.20	0.173 ± 0.013	0.642	0.050 ± 0.004	

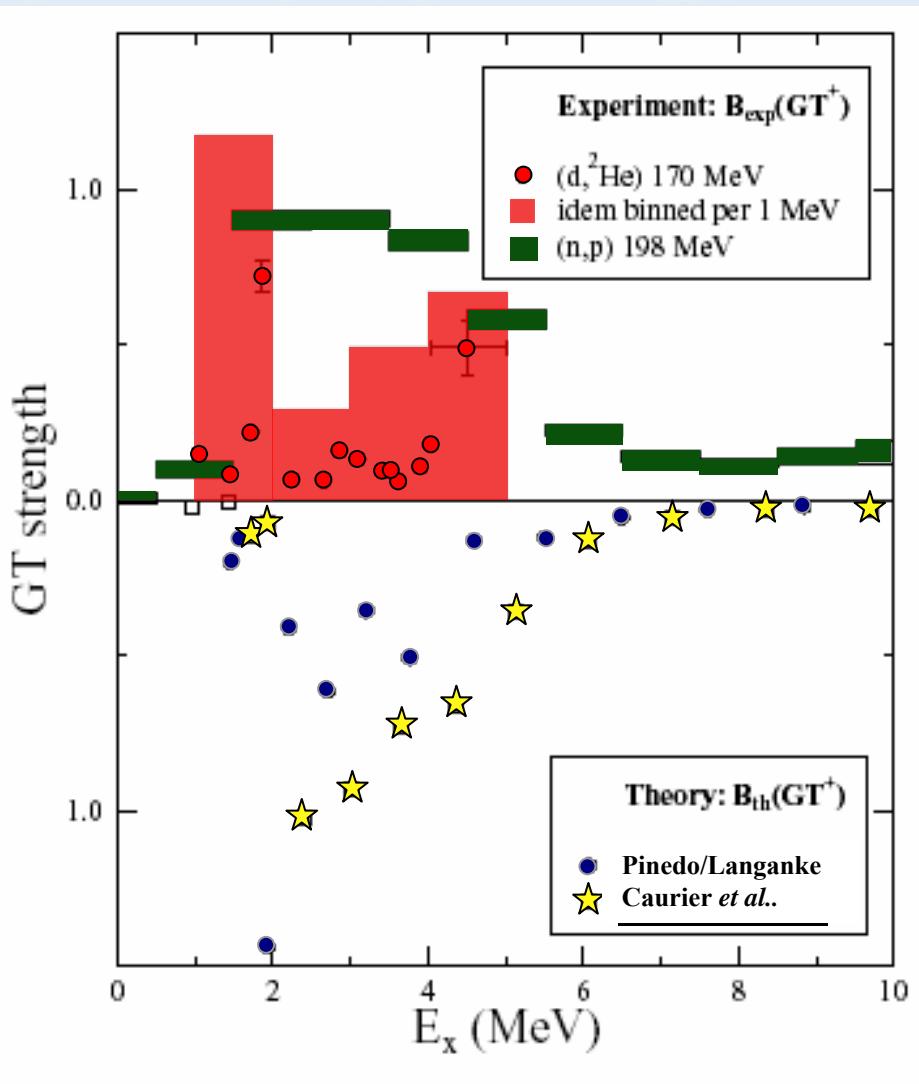
$(d,^2\text{He})$ as GT⁺ probe in *fp* shell nuclei



GT Strength in ^{58}Co from (d, ^2He) reaction

E_x	$d\sigma/d\Omega(0.5^\circ)$	$\sigma(L=0)/\sigma(\tau o \tau)$	$B(GT+)$
[MeV]	[mb/sr]		
1.050	0.159±0.009	0.88	0.15±0.01
1.435	0.078±0.006	1.00	0.09±0.01
1.729	0.148±0.014	1.00	0.16±0.02
1.868	0.648±0.020	1.00	0.72±0.05
2.249	0.047±0.004	1.00	0.05±0.01
2.660	0.057±0.005	0.96	0.06±0.01
2.860	0.145±0.009	0.99	0.17±0.01
3.100	0.126±0.008	0.99	0.15±0.01
3.410	0.065±0.007	0.96	0.07±0.01
3.520	0.080±0.009	0.95	0.09±0.01
3.625	0.067±0.007	0.87	0.07±0.01
3.900	0.062±0.006	0.97	0.07±0.01
4.030	0.155±0.010	1.00	0.19±0.01
4.05-5.00	0.381±0.061		0.49±0.09

GT⁺ strength: comparison (*n,p*), (*d,2He*) & theory



Up to 4 MeV excitation:

13 GT transitions measured (*d,2He*)

Strength rebinned in 1 MeV bins

Significant differences

Updated shell model calculations by
Pinedo/Langanke

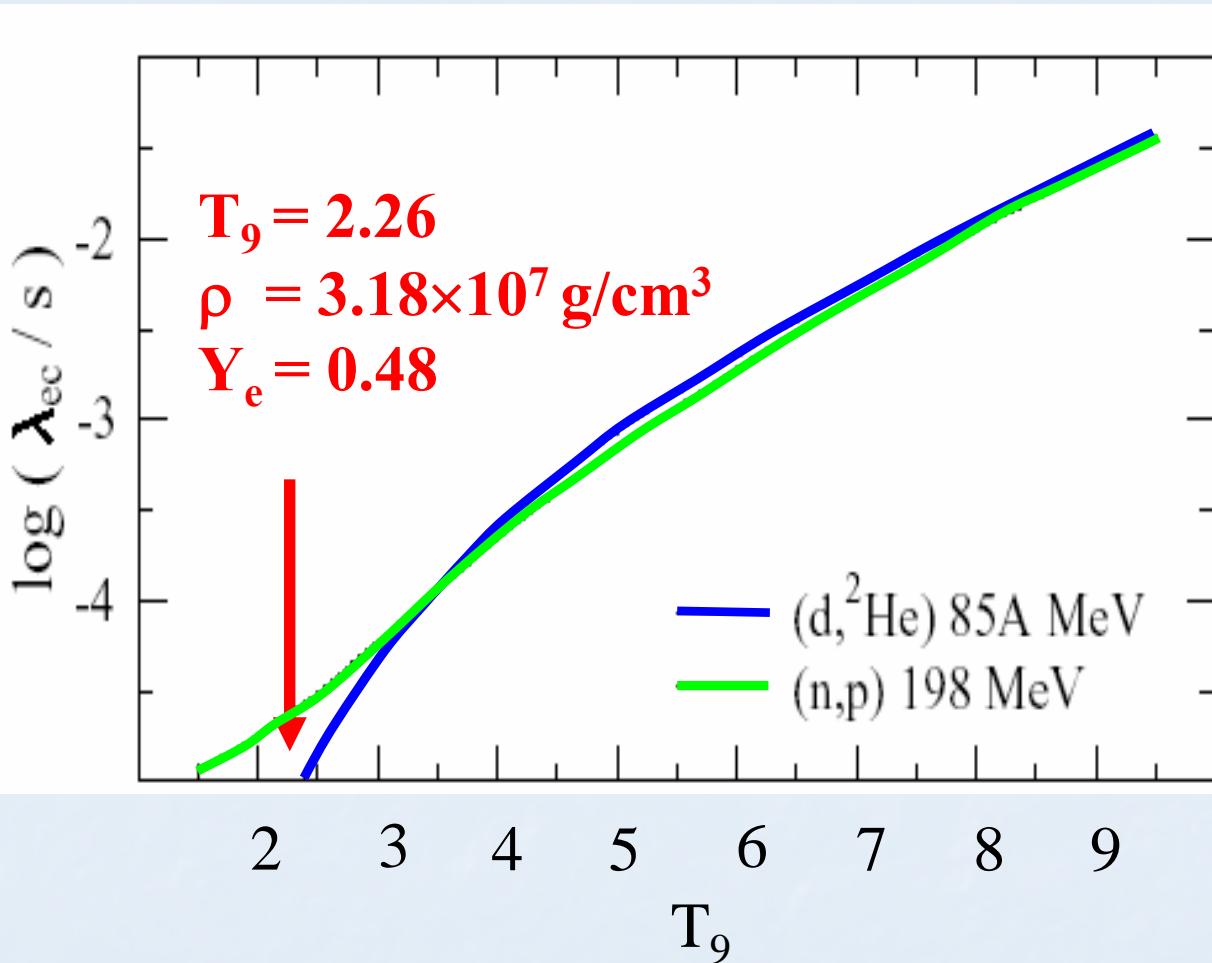
Electron capture rate

$$\lambda_{ec} \approx \sum_i B_i(GT) \int_{\omega_l}^{\infty} \omega p \left(Q_i + \omega \right)^2 F(Z, \omega) S_e(\omega, T) d\omega$$

With

- $B_i(GT)$ Gamow-Teller strength distribution
- ω and p energy and momentum electrons
- $S_e(\omega, T)$ Fermi-Dirac distribution electron gas at temperature T

e^- -capture rates using experimental strengths (Pinedo, Langanke)



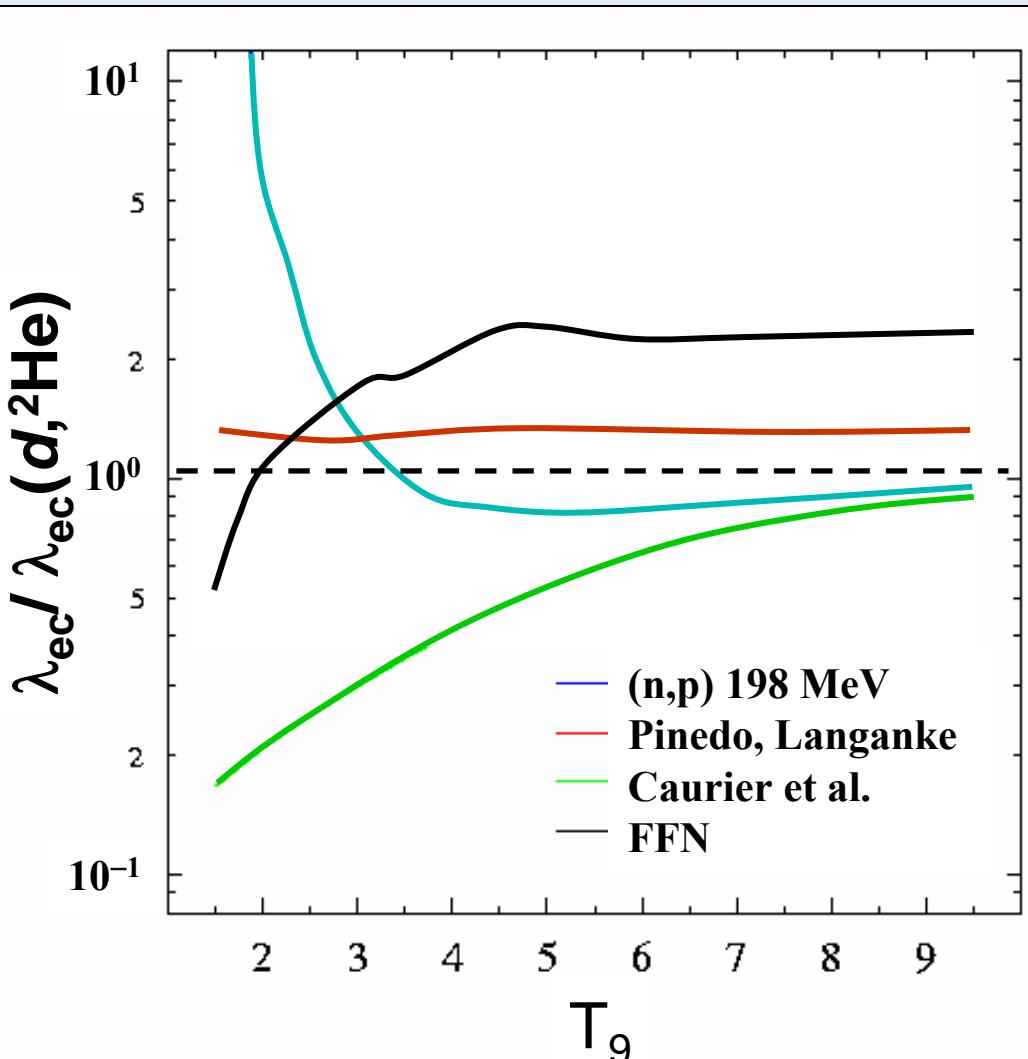
Evolution of core of
 $25 M_{\odot}$ star. Conditions
following silicon
depletion.
[Heger *et al.*, *Astrophys.
J.* 560 (2001) 307]

$T_9 = 4.05$
 $\rho = 3.18 \times 10^7 \text{ g/cm}^3$
 $Y_e = 0.48$

Calculate EC rates for
g.s. as a function of T_9

Strength deviations at low excitation: rates deviation at low T

^{58}Ni : comparison of e-capture rates theory/experiment

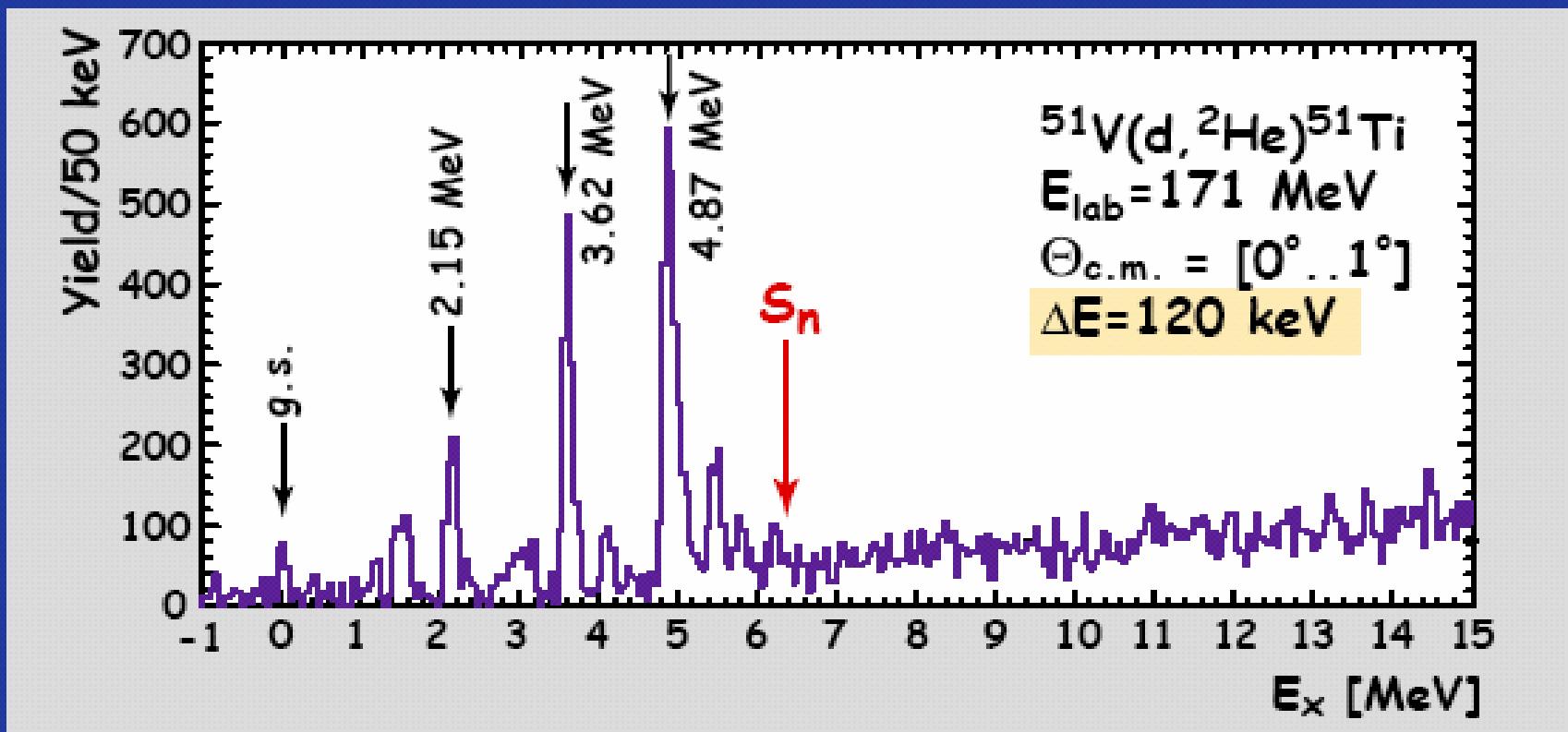


- Influence of GT strength distribution on calculated capture rate is dramatic, especially at low temperatures
- rates vary up to a factor 5-6
- FFN not too far off
- large scale shell-model calculations fail at low T
- calculations with improved residual interaction in reasonable agreement

$^{51}\text{V}(d, ^2\text{He})^{51}\text{Ti}$: B(GT⁺) for proton-odd fp-shell nucleus

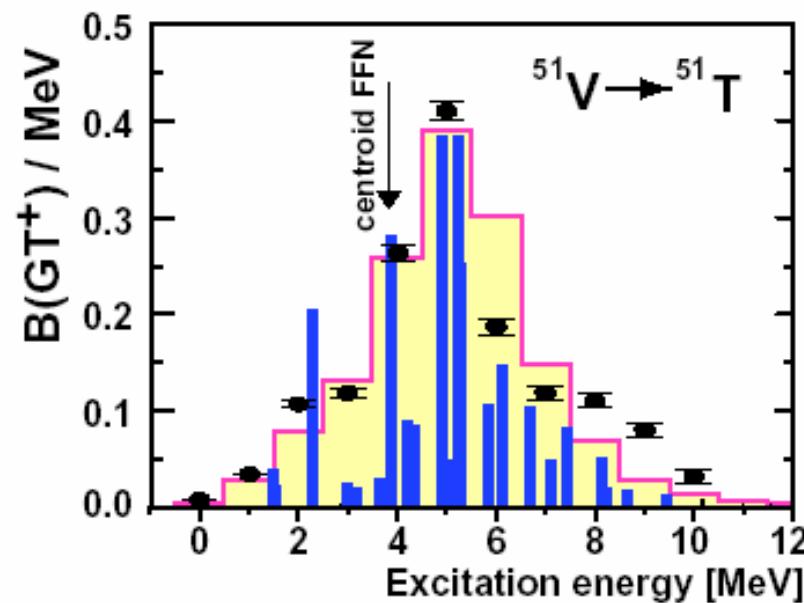
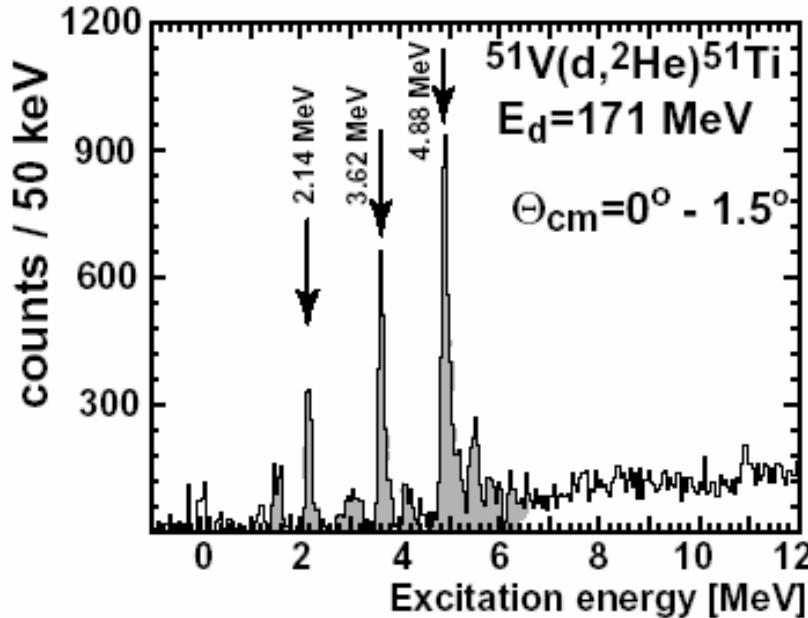
^{51}V g.s. ($J^\pi = 7/2^-$, T=5/2) \Rightarrow ^{51}Ti ($J^\pi = 5/2^-, 7/2^-, 9/2^-, \text{T}=7/2$)

Independent single-particle model (FFN): $E_x(\text{GTR}) = 3.83$ MeV

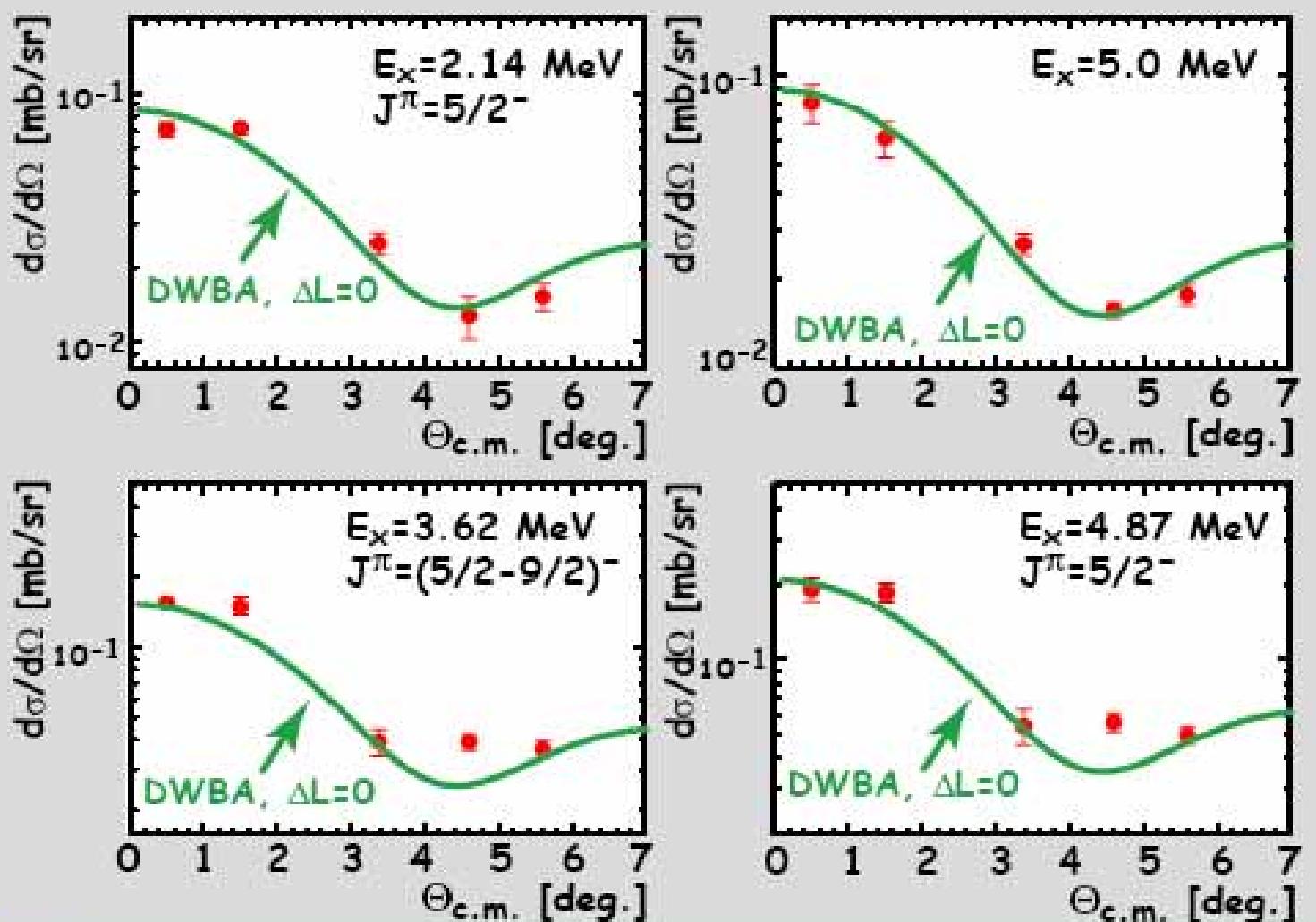


C. Bäumer *et al.*, PRC 68, 031303(R) (2003)

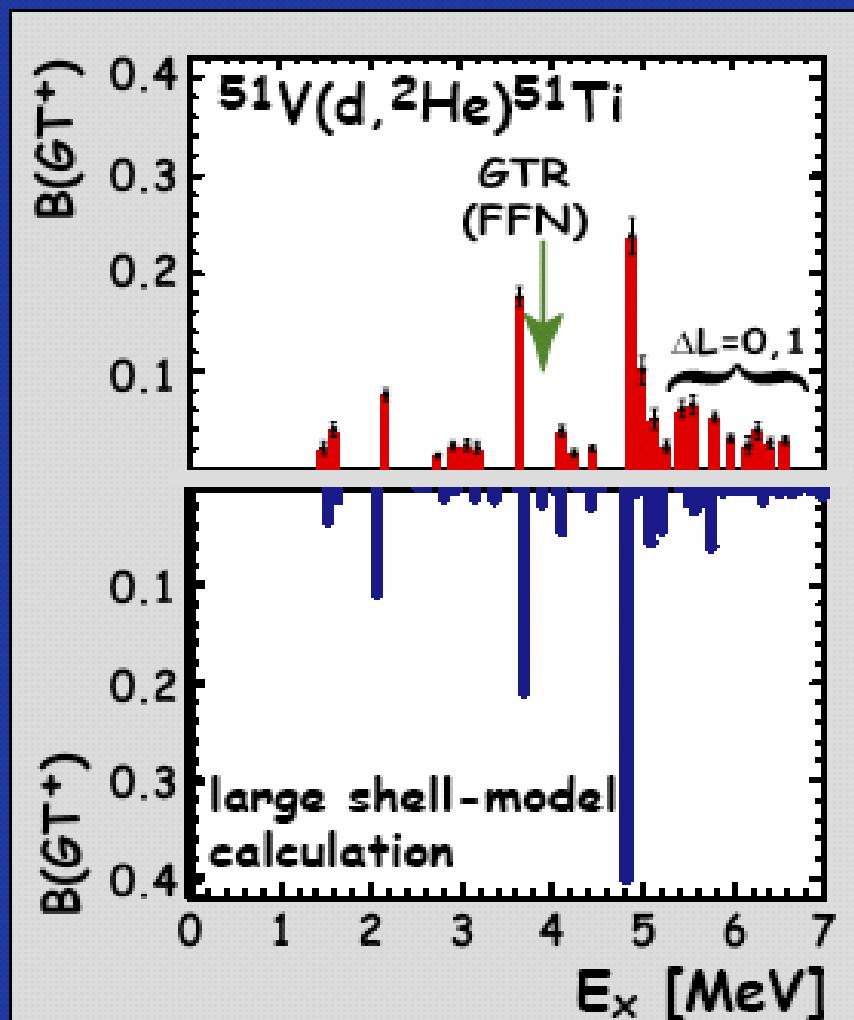
C. Bäumer *et al.*,
PRC **68**, 031303(R)
(2003)



$^{51}\text{V}(d,^2\text{He})$: Angular distributions of $d\sigma/d\Omega$



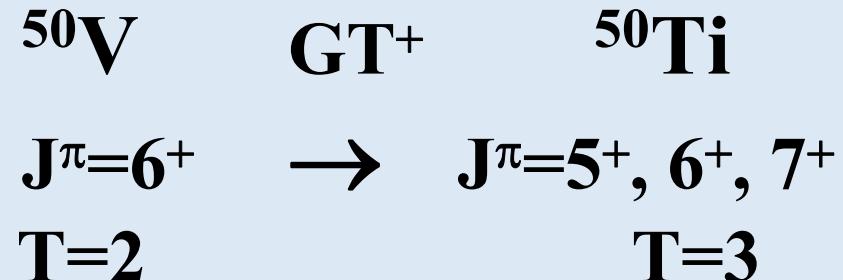
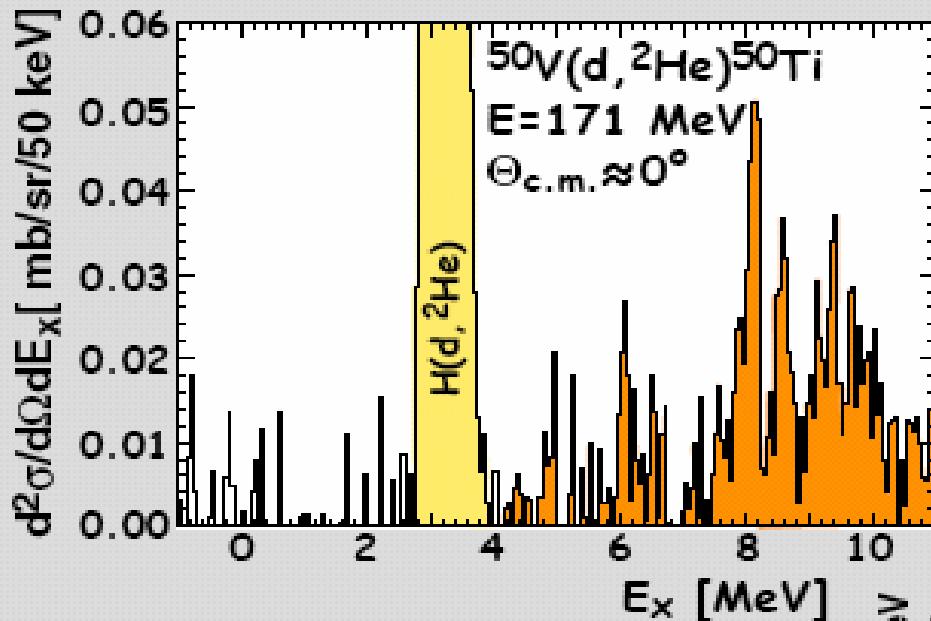
$^{51}\text{V}(d, ^2\text{He})$: Comparison with shell-model calculations



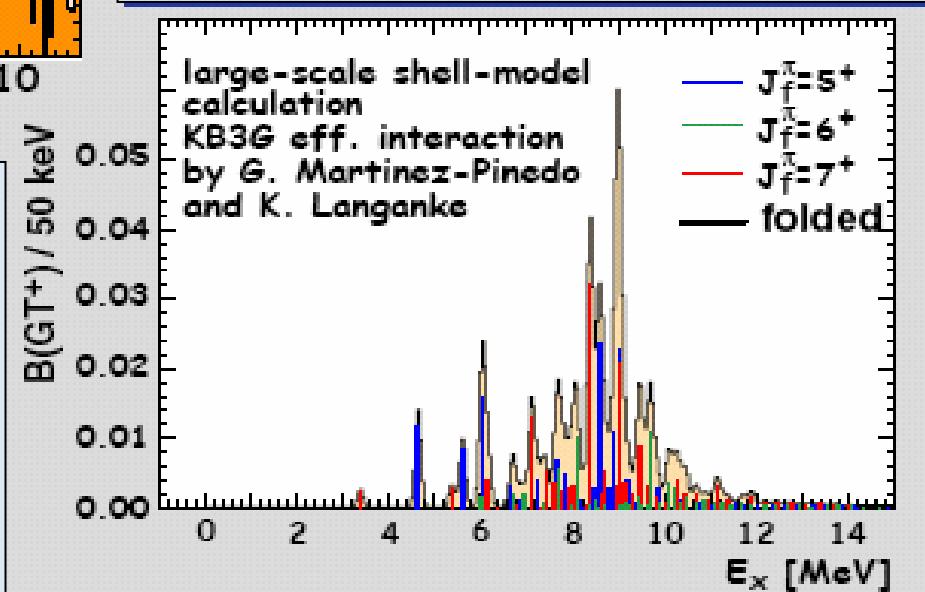
← Experimental result

← Calculations
G. Martínez-Pinedo,
K. Langanke

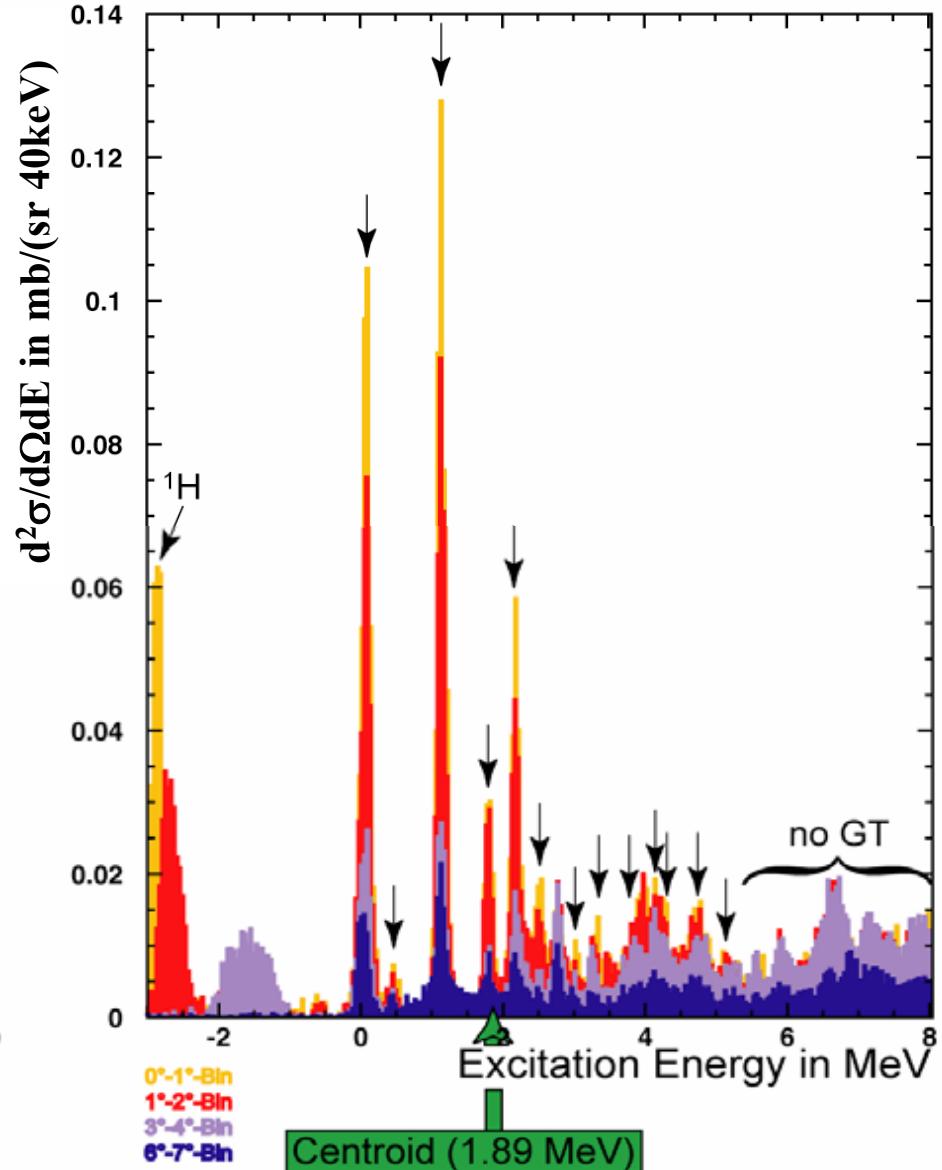
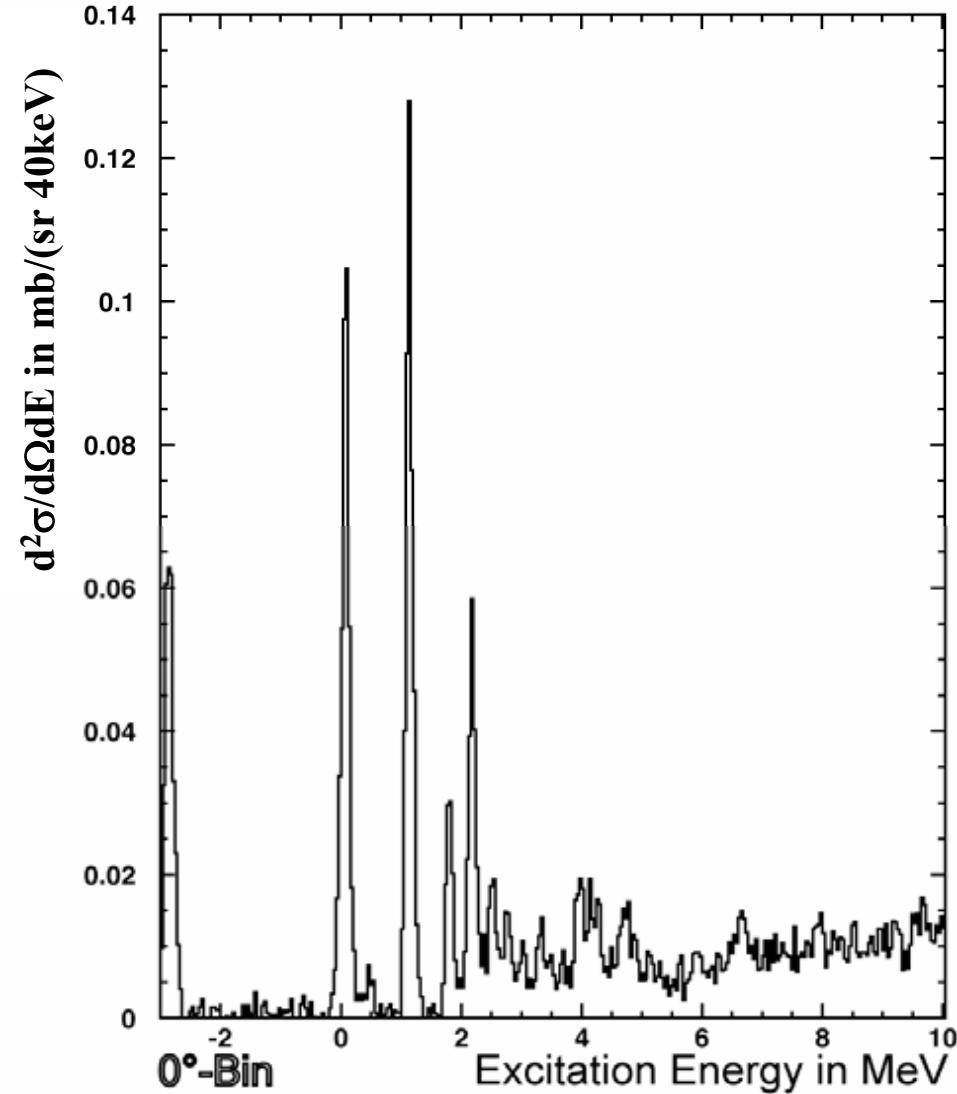
$^{50}\text{V}(d, ^2\text{He})$: GT⁺ transitions from odd-odd nucleus



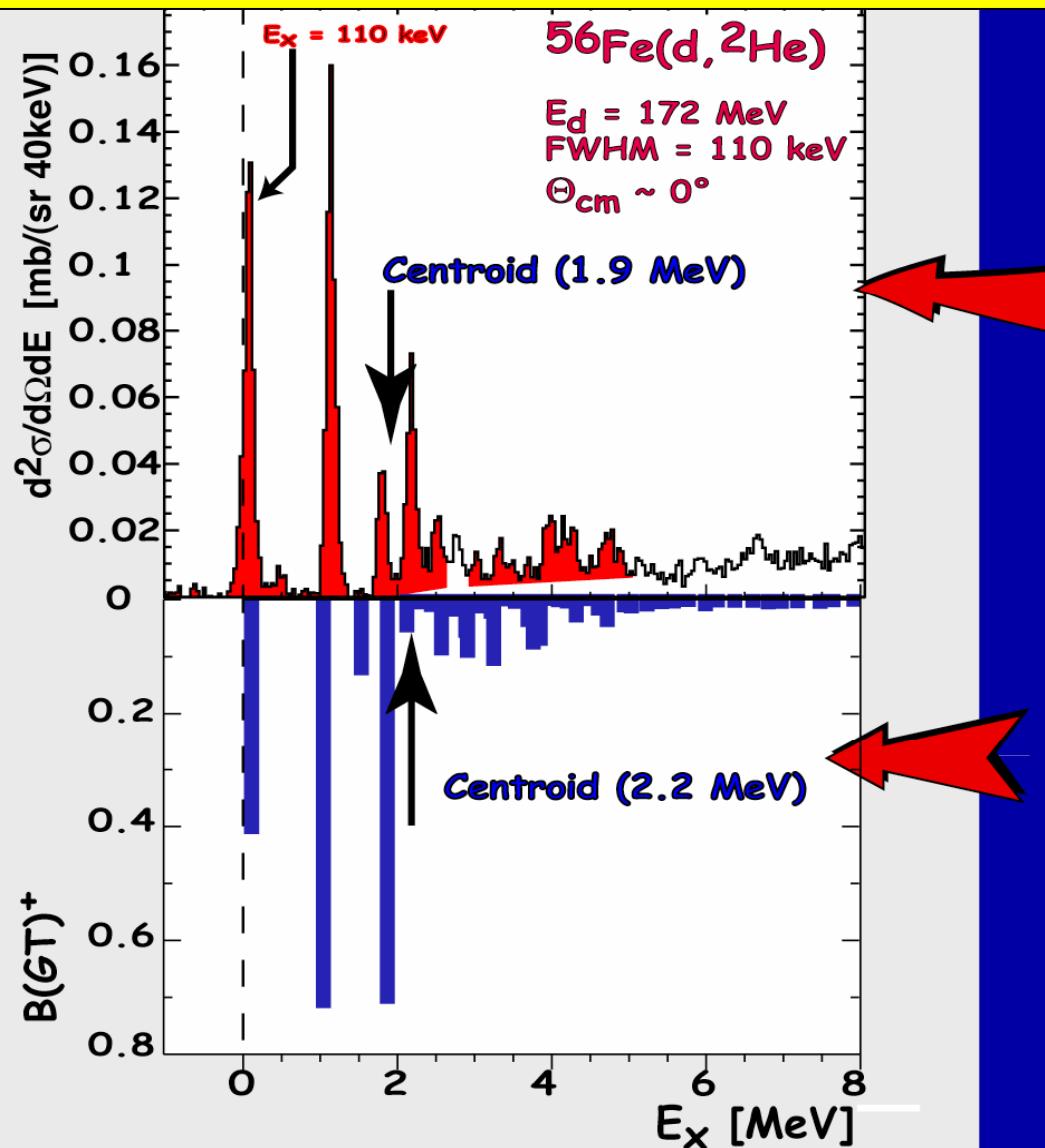
GT-centroid located
at ~ 9 MeV



Spectrum of the $^{56}\text{Fe}(\text{d},^2\text{He})^{56}\text{Mn}$ -Reaction



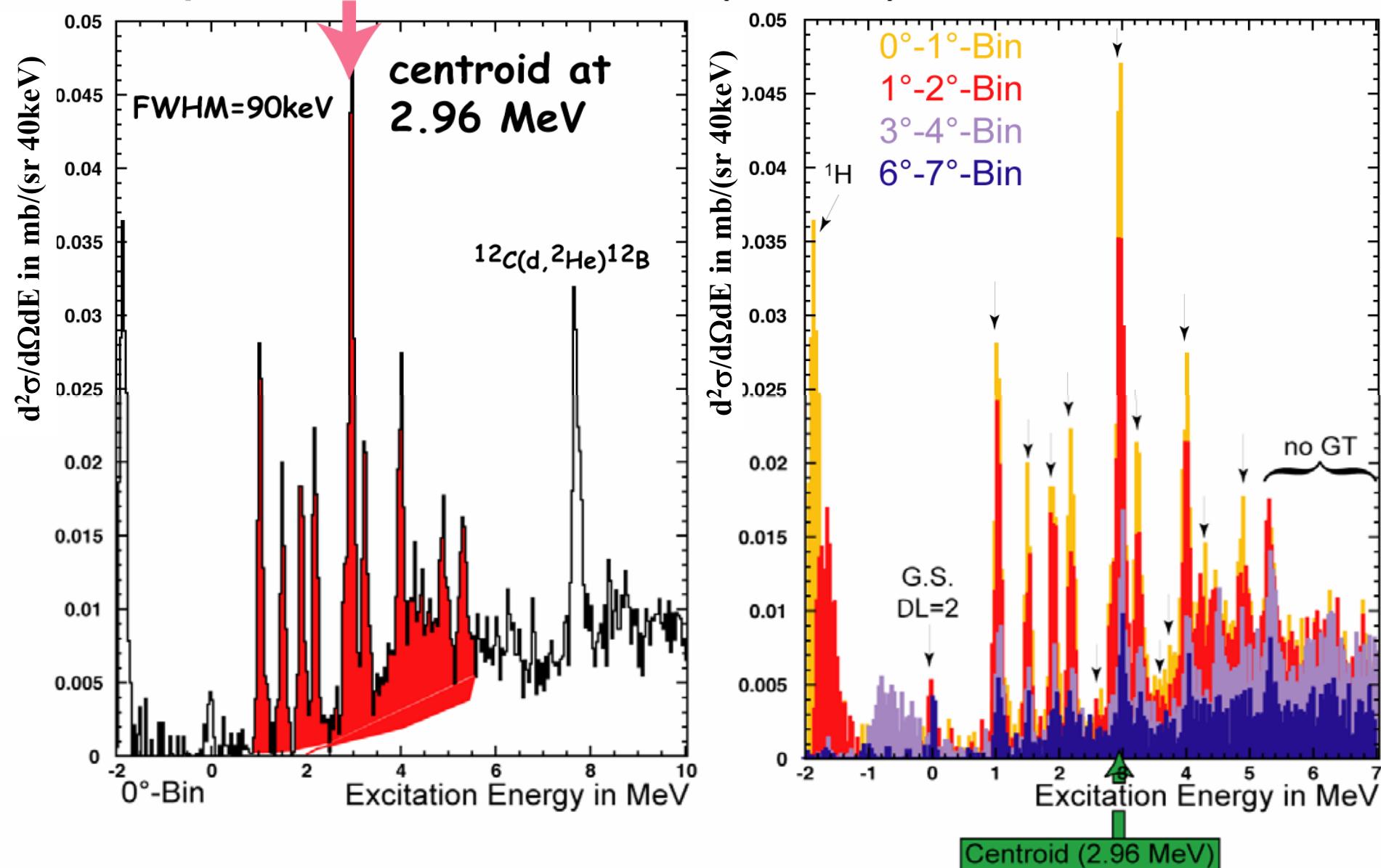
$^{56}\text{Fe}(d,^2\text{He})$: Comparison with shell-model calculations



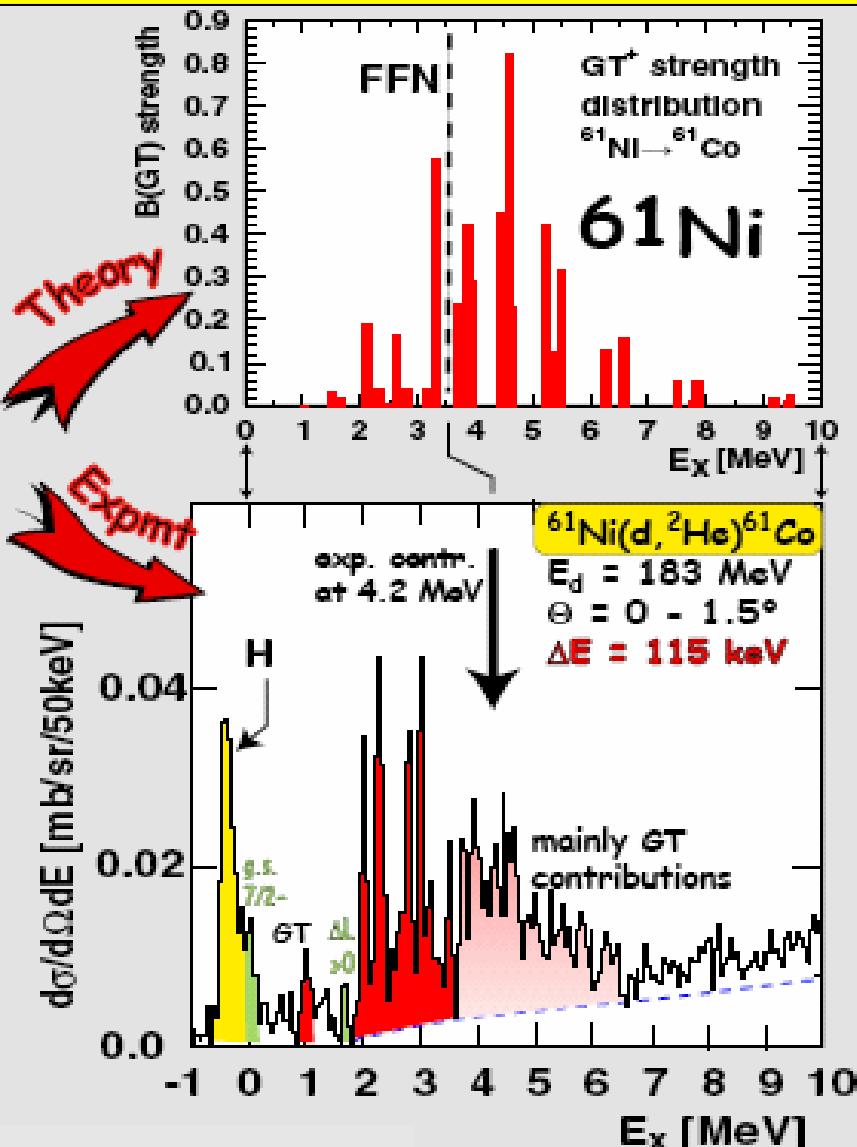
Experiment

Full fp -shell model
calculations (KG3G)
(G. Martínez-Pinedo)

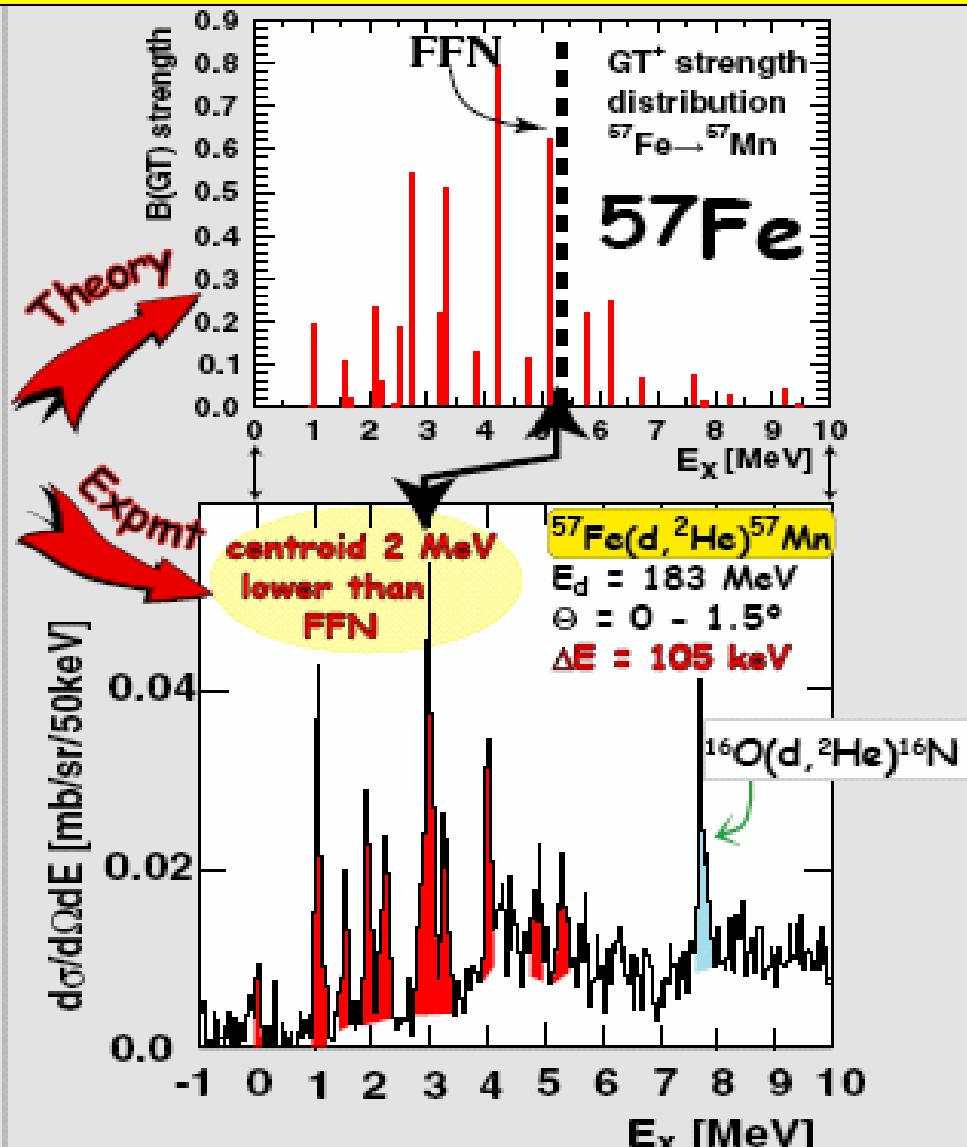
Spectrum of the $^{57}\text{Fe}(\text{d},^2\text{He})^{57}\text{Mn}$ -Reaction



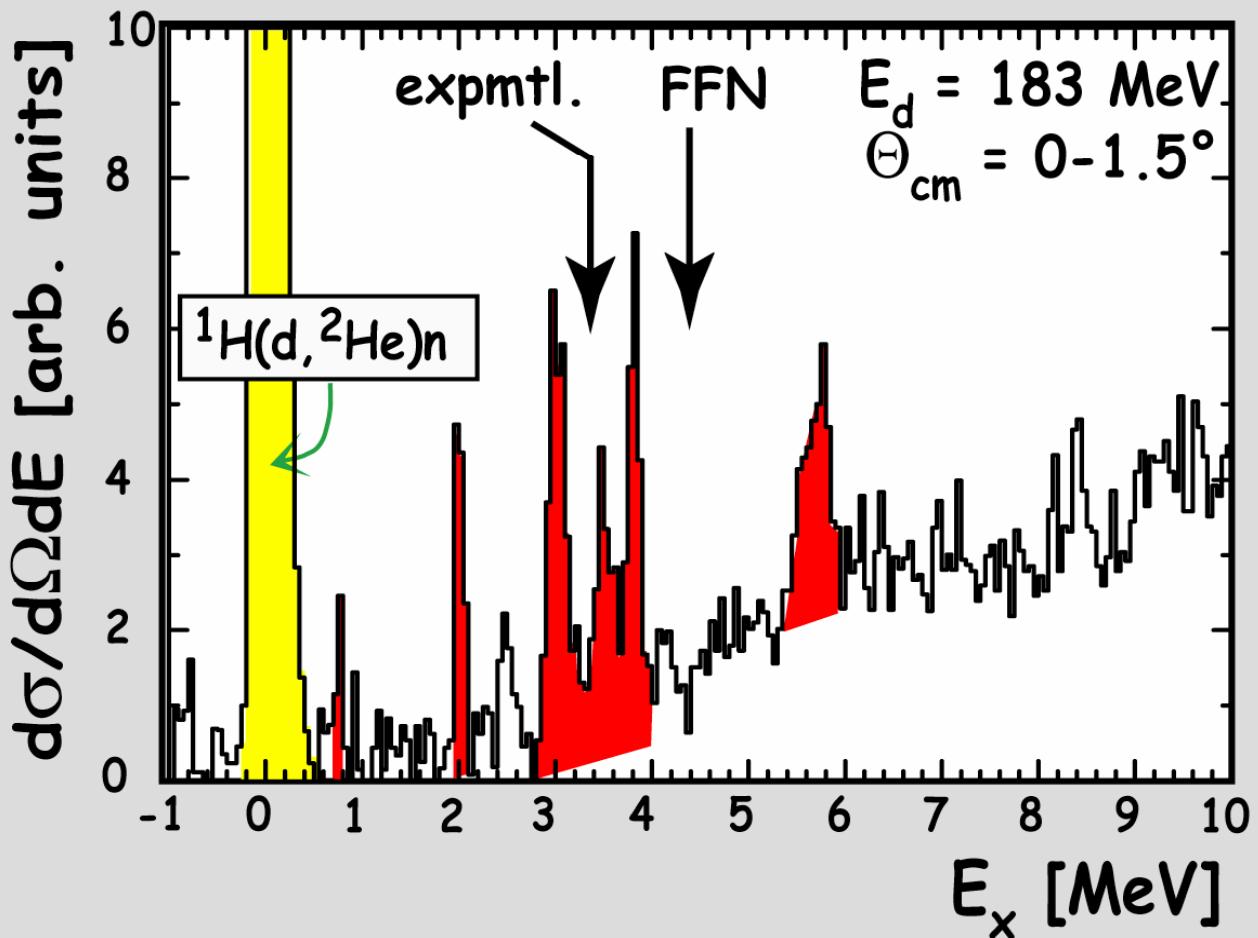
$^{61}\text{Ni}(d,^2\text{He})^{61}\text{Co}$: GT⁺ distribution



$^{57}\text{Fe}(d,^2\text{He})^{57}\text{Mn}$: GT⁺ distribution



$^{67}\text{Zn}(d,^2\text{He})^{67}\text{Cu}$: GT⁺ distribution



GT⁺ centroid comparison

		FFN	SM	Exp.
even-even	Fe-56	3.8	2.2	1.9
	Ni-58	3.8	3.6	3.4
odd-A odd-p	V-51	3.8	4.7	4.1
odd-A odd-n	Fe-57	5.3	4.1	2.9
	Ni-61	3.5	4.6	4.2
	Zn-67	4.4	--	3.4
odd-odd	V-50	9.7	8.5	8.8



Conclusions

- Presupernova models depend sensitively on EC rates.
- GT⁺ transitions in *fp*-shell nuclei play a decisive role in determining EC rates and thus provide input into modeling of explosion dynamics of massive stars.
- Large shell-model calculations are needed especially as function of T. (Caurier *et al.*; Martínez-Pinedo & Langanke [KB3G]; Otsuka *et al.* [GXPF]) \Rightarrow smaller EC rates for A=45-60 than FFN \Rightarrow Larger Y_e (electron to baryon ratio) and smaller iron core mass (Heger *et al.*)
- New high resolution ($d, ^2\text{He}$) experiments provide essential tests for shell model calculations at 0 T.

Outlook

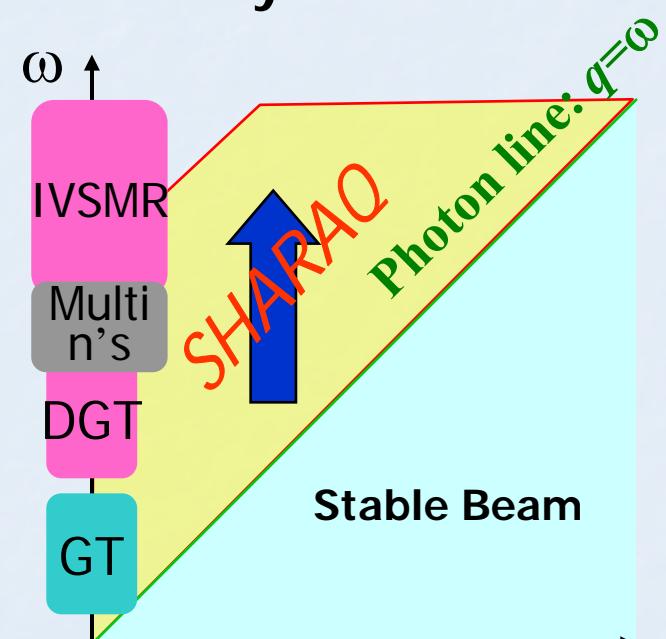
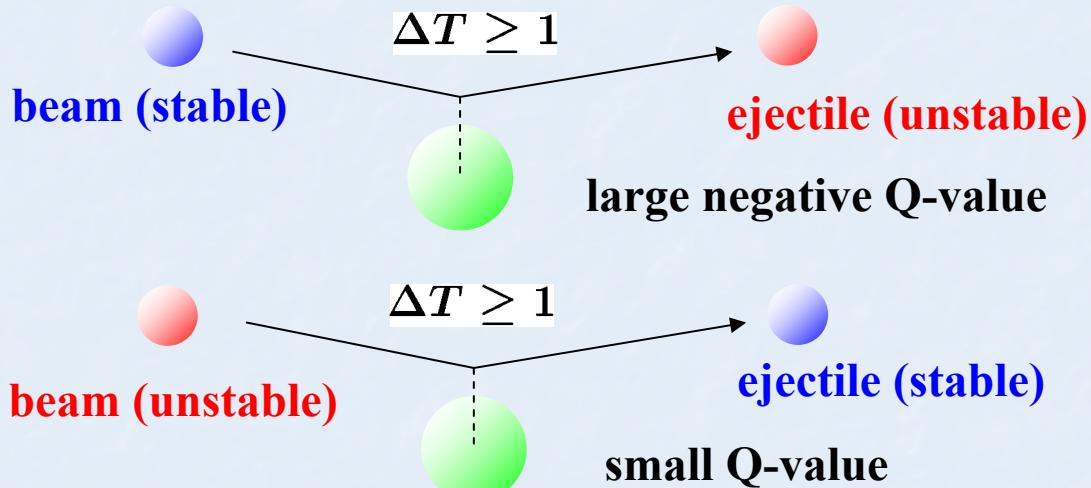
Radioactive ion beams will be available at energies where it will be possible to study GRs (RIKEN, SPIRAL2, NSCL, FAIR, RIA, EURISOL)

- Multipole strength functions in unstable n-rich nuclei soft modes, pygmy resonances, low-lying non-collective strength, modification of effective NN interaction but also various Giant multipole resonances
- ISGMR in a long chain of isotopes determine K_{sym}
- Determine GT strength in unstable *sd* & *fp* shell nuclei
- Use GRs as tools to determine n-skin [IV(S)GDR]
- Exotic excitations such as Double GT



RI Beam ($E = 150\text{-}400 \text{ MeV/A}$) as a new PROBE to nuclear systems

- Large Isospin iso-tensor excitations
- Large internal energy (q, ω) inaccessible by stable beams



Exothermal Charge Exchange Reactions

END