# Energy Levels of Light Nuclei $A=11$ 

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#### Abstract

An evaluation of $A=11-12$ was published in Nuclear Physics A506 (1990), p. 1. This version of $A=11$ differs from the published version in that we have corrected some errors discovered after the article went to press. Figures and introductory tables have been omitted from this manuscript. Also, reference key numbers have been changed to the NNDC/TUNL format.


(References closed June 1, 1989)

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Table 11.1: Energy Levels of ${ }^{11} \mathrm{Li}^{\text {a }}$

| $E_{\mathrm{x}}(\mathrm{MeV})$ | $J^{\pi} ; T$ | $\tau_{1 / 2}(\mathrm{~ms})$ | Decay | Reaction |
| :---: | :---: | :---: | :---: | :---: |
| g.s. | $\frac{3}{2}^{-}, \frac{5}{2}$ | $8.5 \pm 0.2$ | $\beta^{-}$ | 1 |

$\quad$ a Excited states are calculated at $E_{\mathrm{x}}=2.68,3.13$ and 3.62 MeV , with $J^{\pi}=\frac{5}{2}^{+}$,
$\frac{3}{2}^{+}$and $\frac{9}{2}$
$\frac{3}{2}^{-}$and $\frac{5}{2}^{-}[(0+1) \hbar \omega$ model space $]$ and at $4.58,21.69$ and 23.22 MeV , with $J^{\pi}=\frac{1}{2}^{-}$,
${ }^{11} \mathrm{He}$
(not illustrated)
${ }^{11} \mathrm{He}$ has not been reported: see (80AJ01). The ground state of ${ }^{11} \mathrm{He}$ is predicted to have $J^{\pi}=\frac{5}{2}^{+}(85 \mathrm{PO} 10)$.

$$
\begin{gathered}
{ }^{11} \mathrm{Li} \\
\text { (Figs. } 1 \text { and } 4 \text { ) }
\end{gathered}
$$

GENERAL (See also (85AJ01).)
The mass excess is $40.94 \pm 0.08 \mathrm{MeV}$ (75TH08), $40.78 \pm 0.12 \mathrm{MeV}$ (88WO09). A.H. Wapstra suggests (private communication) $40.85 \pm 0.08 \mathrm{MeV}$ and we adopt this value. ${ }^{11} \mathrm{Li}$ is then bound with respect to ${ }^{9} \mathrm{Li}+2 \mathrm{n}$ by $247 \pm 80 \mathrm{keV}$ and with respect to ${ }^{10} \mathrm{Li}+\mathrm{n}$ by $1050 \pm 260 \mathrm{keV}$ [see (88AJ01) for the masses of ${ }^{9} \mathrm{Li}$ and $\left.{ }^{10} \mathrm{Li}\right]$.

The magnetic moment of ${ }^{11} \mathrm{Li}$ is $\mu=3.6673 \pm 0.0025 \mathrm{~nm}$ (87AR22). This value requires $J=\frac{3}{2}$ (87AR22). Negative parity is certain from systematics.

The interaction nuclear radius of ${ }^{11} \mathrm{Li}$ is $3.16 \pm 0.11 \mathrm{fm}$ (88TA10, 85TA18), $E=790$ $\mathrm{MeV} / A$; [see also for derived nuclear matter, charge and neutron matter r.m.s. radii]. ${ }^{11} \mathrm{Li}$ has a much larger radius than other neighboring nuclei suggesting either a large deformation and/or a long tail in the matter distribution in ${ }^{11} \mathrm{Li}$ (85TA18). See (88SA2P) and (87MI1A, 87TA1F, 89TA1K). Charge radius and matter radius calculations in the $0 \hbar \omega$ and $(0+2) \hbar \omega$ model spaces predict a gradual increase in matter radii with increasing $A$ and do not support the idea of a neutron halo in ${ }^{11} \mathrm{Li}$ (88PO1E; prelim.). See, however, (88TA1A).

Fragmentation cross sections of ${ }^{11} \mathrm{Li}$ into ${ }^{9} \mathrm{Li},{ }^{8} \mathrm{Li},{ }^{8} \mathrm{He},{ }^{7} \mathrm{Li},{ }^{6} \mathrm{Li}$ and ${ }^{6} \mathrm{He}$ have been studied by (88KO10) [see for a discussion of neutron halos]. See also (88TA1A, 89KO1P).

See also (86DU11, 88ST06), (86AN07, 88HA1Q, 88TA1C, 89AJ1A) and (85SA32, 86EL1A, 86SA30, 87HA30, 87SH1K, 88BE09, 88BE1O, 88JO1C, 88LO1C, 88UC03, 89BA1T, 89BE03; theor.).

1. ${ }^{11} \mathrm{Li}\left(\beta^{-}\right){ }^{11} \mathrm{Be}$
$Q_{\mathrm{m}}=20.68$

Reported half-life measurements are $8.5 \pm 0.2 \mathrm{~ms}$ (74RO31), $8.83 \pm 0.12 \mathrm{~ms}$ (81BJ01), $7.7 \pm 0.6 \mathrm{~ms}$ (86CU01). We adopt $8.5 \pm 0.2 \mathrm{~ms}$. The $\beta$-decay is complex and the evidence is not unambiguous. It involves delayed $\mathrm{n}, \mathrm{t}$ and $\alpha$ emission. Most of the decay ( $\approx 97 \%$ ) takes place to low-lying states in ${ }^{11} \mathrm{Be}$ [but it is not clear which are involved]. All but ${ }^{11} \mathrm{Be}^{*}(0$, 0.32 ) are unstable with respect to neutron emission: see (85AJ01). A $2.9 \%$ branch is reported to ${ }^{11} \mathrm{Be}^{*}(10.59)$ which then decays by neutron emission (possibly to ${ }^{10} \mathrm{Be}^{*}(9.4)$ ) and then delayed $\alpha$-particles $[(0.90 \pm 0.05) \%]$ are reported to ${ }^{6} \mathrm{He}$ or the decay is via $2 \mathrm{n}+2 \alpha[(2.0 \pm 0.6) \%]$ (81LA11). A $(0.30 \pm 0.05) \%$ branch is reported to a state in ${ }^{11} \mathrm{Be}$ at $\approx 18.5 \mathrm{MeV}(\Gamma \approx 0.5$ $\mathrm{MeV})$ which has three modes of decay: triton emission to ${ }^{8} \mathrm{Li}^{*}(0,0.98)$ [(0.010 $\left.\pm 0.004\right) \%$ ], $(\alpha+\mathrm{n})$-emission to ${ }^{6} \mathrm{He}[(0.10 \pm 0.03) \%]$ and 3 n emission $\left[(0.20 \pm 0.05 \%]\right.$ involving ${ }^{10} \mathrm{Be}^{*}(11.76)$ (81LA11, 84LA27). [Comment: In view of the importance of understanding very neutron rich light nuclei it is necessary to determine the parameters (and the location) of the excited states of ${ }^{11} \mathrm{Be}$ with $E_{\mathrm{x}} \lesssim 12 \mathrm{MeV}$. One could then hope to unravel the $\beta^{-}$decay evidence.] See also ${ }^{6} \mathrm{He},{ }^{8} \mathrm{Li}$ and ${ }^{10} \mathrm{Be}$ in (AJ88), Table 11.2 in (85AJ01), (85HA1T, 85HA1K) and (84LI1N, $88 \mathrm{JO1C}$; theor.).

$$
{ }^{11} \mathrm{Be}
$$

(Figs. 1 and 4)
GENERAL: See also (85AJ01).
Model calculations: (84MI1H, 84VA06, 86WI04).
Electromagnetic transitions: (84MI1H, 84VA06, 87HO1L).
Complex reactions involving ${ }^{11} \mathrm{Be}:(85 \mathrm{BO} 1 \mathrm{~A}, 86 \mathrm{AV} 1 \mathrm{~B}, ~ 87 \mathrm{TR} 05,87 \mathrm{WA} 09,88 \mathrm{BA} 53,88 \mathrm{RU} 01$, $88 \mathrm{TA} 1 \mathrm{~N}, 88 \mathrm{TR} 03,89 \mathrm{SA} 10)$.

Muon and neutrino capture and reactions: (84KO24).
Hypernuclei: (85IK1A, 86ME1F).
Other topics: (84MI1H, 85AN28, 86AN07).
Ground-state properties of ${ }^{11} \mathrm{Be}:(84 \mathrm{FR} 13,87 \mathrm{SA} 15,87 \mathrm{VA} 26,89 \mathrm{BE} 03)$.
The interaction matter radius of ${ }^{11} \mathrm{Be}$ is $2.86 \pm 0.04 \mathrm{fm}$ (88TA10). See also (?, 89TA1K).

1. ${ }^{11} \mathrm{Be}\left(\beta^{-}\right){ }^{11} \mathrm{~B} \quad Q_{\mathrm{m}}=11.506$

The decay is complex: see reaction 26 in ${ }^{11} \mathrm{~B}$ and Table 11.12 . The half-life is $13.81 \pm 0.08 \mathrm{~s}$ (70AL21). See also (80AJJ01).
2. ${ }^{9} \mathrm{Be}(\mathrm{t}, \mathrm{p}){ }^{11} \mathrm{Be}$
$Q_{\mathrm{m}}=-1.165$

Table 11.2: Energy Levels of ${ }^{11} \mathrm{Be}$

| $E_{\mathrm{x}}(\mathrm{MeV} \pm \mathrm{keV})$ | $J^{\pi} ; T$ | $\tau$ or $\Gamma_{\text {c.m. }}(\mathrm{keV})$ | Decay |  |
| :---: | :---: | :---: | :---: | :--- |
|  |  |  |  | Reactions |
| 0 | $\frac{1}{2}^{+} ; \frac{3}{2}$ | $\tau_{1 / 2}=13.81 \pm 0.08 \mathrm{~s}$ | $\beta^{-}$ | $1,2,3,5,7$ |
| $0.32004 \pm 0.1$ | $\frac{1}{2}^{-}$ | $\mathrm{fs} \tau_{\mathrm{m}}=166 \pm 15 \mathrm{fs}$ | $\gamma$ | $2,3,4,5,7,8$ |
| $1.778 \pm 12$ | $\left(\frac{5}{2}, \frac{3}{2}\right)^{+}$ | $\Gamma=100 \pm 20$ | $(\mathrm{n})$ | $2,3,6,7$ |
| $2.69 \pm 20$ | $\left(\frac{1}{2}, \frac{3}{2}, \frac{5}{2}^{+}\right)$ | $200 \pm 20$ | $(\mathrm{n})$ | 2,8 |
| $3.41 \pm 20$ | $\left(\frac{1}{2}, \frac{3}{2}, \frac{5}{2}^{+}\right)$ | $125 \pm 20$ | $(\mathrm{n})$ | 2,7 |
| $3.887 \pm 15$ | $\geq \frac{7}{2}^{-}$ | $<10$ | $(\mathrm{n})$ | 2 |
| $3.956 \pm 15$ | $\frac{3}{2}^{-}$ | $15 \pm 5$ | $(\mathrm{n})$ | 2,8 |
| $5.240 \pm 21$ |  | $45 \pm 10$ | $(\mathrm{n})$ | 2 |
| $(5.86)$ |  | $\approx 300$ | $(\mathrm{n})$ | 2 |
| $6.51 \pm 50$ |  | $120 \pm 50$ | (n) | 2 |
| $6.705 \pm 21$ |  | $40 \pm 20$ | $(\mathrm{n})$ | 2 |
| $7.03 \pm 50$ |  | $300 \pm 100$ | $(\mathrm{n})$ | 2 |
| $8.816 \pm 32$ |  | $200 \pm 50$ | $(\mathrm{n})$ | 2,4 |
| $10.59 \pm 50$ |  | $210 \pm 40$ | $\mathrm{n}, \alpha$ | 2,4 |
| $(\approx 18.5)$ |  | $\approx 500$ | $\mathrm{n}, \mathrm{t}, \alpha$ | 4 |

Proton groups have been observed to the states displayed in Table 11.2. $\tau_{\mathrm{m}}$ for the first excited state is $166 \pm 15 \mathrm{fs}$, corresponding to a very large E1 transition strength of $0.36 \pm 0.03$ W.u.; $E_{\gamma}=320.04 \pm 0.10 \mathrm{keV}$. The $J^{\pi}$ of ${ }^{11} \mathrm{Be}^{*}(0.32)$ is $\frac{1}{2}^{-}$, as determined by a study of the yield of $320 \mathrm{keV} \gamma$-rays as a function of time in $\mu^{-}$capture by ${ }^{11} \mathrm{~B}$. The strength of the E1 transition fixes $J^{\pi}$ of ${ }^{11} \mathrm{Be}\left(\right.$ g.s. ) to be $\frac{1}{2}^{+}$or $\frac{3}{2}^{+}$, using the parity information obtained from the nature of the $\beta^{-}$decay of the ground state [see reaction 26 in ${ }^{11} \mathrm{~B}$ ]. ${ }^{11} \mathrm{Be}^{*}(5.24$, $6.71,8.82$ ) are strongly populated at $E_{\mathrm{t}}=20 \mathrm{MeV}$ indicating that these states have a large overlap with ${ }^{9} \mathrm{Be}_{\text {g.s. }}+2 \mathrm{n}$. See (80AJ01, 85AJ01) for references.
3. ${ }^{10} \mathrm{Be}(\mathrm{d}, \mathrm{p}){ }^{11} \mathrm{Be}$
$Q_{\mathrm{m}}=-1.720$

Angular distributions of the $\mathrm{p}_{0}$ and $\mathrm{p}_{1}$ groups have been measured at $E_{\mathrm{d}}=6 \mathrm{MeV}$ and $12 \mathrm{MeV}: l_{\mathrm{n}}=0\left[\right.$ and therefore $J^{\pi}=\frac{1}{2}^{+}$for $\left.{ }^{11} \mathrm{Be}(0)\right]$ and $\mathrm{l}, S=0.73 \pm 0.06$ and $0.63 \pm 0.15$, respectively. At $E_{\mathrm{d}}=25 \mathrm{MeV}^{11} \mathrm{Be}^{*}(0,0.32,1.78)$ are strongly populated: $S=0.77,0.96$, and 0.50 , respectively, $J^{\pi}=\left(\frac{5}{2}, \frac{3}{2}\right)^{+}$for ${ }^{11} \mathrm{Be}^{*}(1.78)\left[l_{\mathrm{n}}=2\right]$. See (80AJ01) for references.
4. ${ }^{11} \mathrm{Li}\left(\beta^{-}\right){ }^{11} \mathrm{Be}$
$Q_{\mathrm{m}}=20.68$

See ${ }^{11} \mathrm{Li}$.
5. ${ }^{11} \mathrm{~B}\left(\pi^{-}, \gamma\right){ }^{11} \mathrm{Be}$
$Q_{\mathrm{m}}=128.063$

The photon spectrum from stopped pions includes a peak corresponding to ${ }^{11} \mathrm{Be} *(0 \pm 0.32)$ (86PE05).
6. ${ }^{12} \mathrm{C}\left(\pi^{-}, \mathrm{p}\right){ }^{11} \mathrm{Be} \quad Q_{\mathrm{m}}=112.105$

See (87BL07; $\left.E_{\pi^{-}}=145 \mathrm{MeV}\right)$.
7. ${ }^{12} \mathrm{C}\left({ }^{7} \mathrm{Li},{ }^{8} \mathrm{~B}\right){ }^{11} \mathrm{Be}$
$Q_{\mathrm{m}}=-28.187$

At $E\left({ }^{7} \mathrm{Li}\right)=82 \mathrm{MeV}{ }^{11} \mathrm{Be}^{*}(0+0.32,1.8,3.4)$ are populated (85AL1G).
8. ${ }^{13} \mathrm{C}\left({ }^{6} \mathrm{Li},{ }^{8} \mathrm{~B}\right){ }^{11} \mathrm{Be}$
$Q_{\mathrm{m}}=-25.884$

At $E\left({ }^{6} \mathrm{Li}\right)=80 \mathrm{MeV},{ }^{11} \mathrm{Be} e^{*}(0.32)$ is strongly populated and the angular distribution to this state has been measured. ${ }^{11} \mathrm{Be}^{*}(2.69,4.0)$ are also observed: see (80AJ01).

$$
\begin{gathered}
{ }^{11} \mathrm{~B} \\
\text { (Figs. } 2 \text { and } 4 \text { ) }
\end{gathered}
$$

GENERAL: See also (85AJJ01).
Nuclear models: (84ZW1A, 85KW02, 87KI1C, 88OR1C, 88WO04).
Special states: (84ZW1A, 85CH27, 85GO1A, 85HA1J, 85SH24, 87KI1C, 88KW1A, 88ZH1B, 89BA60, 89OR02).

Electromagnetic transitions and giant resonances: (83GM1A, 84MO1D, 84VA06, 85GO1A, 86ER1A, 87KI1C, 89BA60).

Astrophysical questions: (82AU1A, 82CA1A, 84TR1C, 85DW1A, 85WA1K, 87AR1J, 87AU1A, 87DW1A, 87MA2C, 87RO1D, 87WE1E, 88AP1A, 88BA1H, 88FE1A, 88RE1B, 89BO1F, 89BO1M, 89GU1Q, 89JI1A).

Complex reactions involving ${ }^{11} \mathrm{~B}$ : (84AI1A, 84FI17, 84HO23, 84RE1A, 84SI15, 84XI1B, 85AG1A, 85BE40, 85BH02, 85JA1B, 85MC03, 85MO08, 85PO11, 85SH1G, 85SI19, 85WA1F, 85WA22, 86AV1B, 86BA69, 86BI1A, 86BO1B, 86CH2G, 86CS1A, 86HA1B, 86MA19, 86ME06, 86MO15, 86PO06, 86RE13, 86SA30, 86SH2B, 86UT01, 86WA1H, 86WE1C, 87AN1A, 87AR19, 87BA1G, 87BA38, 87BE58, 87BE55, 87BO1K, 87BU07, 87DE37, 87FE1A, 87GR1O, 87JA06, 87KI05, 87KO15, 87LY04, 87MA2F, 87MU1D, 87NA01, 87OS1E, 87PA01, 87PO1I, 87SH23, 87SI1C, 87ST01, 87TE1D 87TR05, 87VI02, 87WA09, 87WE1D, 87YA16, 88BA53, 88BL09, 88CA06, 88FE1A, 88FO03, 88GA12, 88KA1L, 88KH1G, 88KI05, 88KI06, 88MI28, 88MO1K, 88PAZS, 88RA10, 88RU01, 88SA19, 88TE03, 88UT02, 89BL1D, 89CEZZ, 89HA1L, 89PA06, 89PO06, 89SA10, 89SE03, 89ST1G, 89YO02).

Applications: (84CA1D, 86NO1C, 88XI1B).
Muon and neutrino capture and reactions: (83GM1A, 84KO24, 85MI1D, 86KE1Q, 87KU23, 87SU06, 87WE1E, 88RA1E, 89MI1G).

Pion and kaon capture and reactions (see also reactions 20, 30, and 48): (83GE1C, 83GM1A, 84BA1T, 84BA1U, 85CO16, 86PE05, 86RO03, 87AB1E, 87BO1X, 88GIZU).

Antinucleon interactions: (85BA51).
Hypernuclei: (83SH1E, 84CH1G, 84SH1J, 84ZH1B, 85AH1A, 85GA1E, 85GR10, 86AN1R, 86BA3L, 86BI1G, 86DA1H, 86DA1G, 86DA1B, 86DU1P, 86FR1J, 86GA1J, 86GA1H, 86KI1K, 86KO1A, $86 \mathrm{ME} 1 \mathrm{~F}, 86 \mathrm{PO} 1 \mathrm{H}, ~ 86 \mathrm{SZ1A}, ~ 86 \mathrm{YA} 1 \mathrm{~F}, ~ 87 \mathrm{MI} 1 \mathrm{~A}, ~ 87 \mathrm{PO} 1 \mathrm{H}, 88 \mathrm{MA1G}, 88 \mathrm{MO} 1 \mathrm{~L}$, 88TA1B, 88TA14, 89MI30).

Other topics: (84P()11, 85AN28, 85SH24, 86AN07, 88KW1A, 880R1C, 89BA60, 890R02).

Table 11.3: Energy Levels of ${ }^{11} \mathrm{~B}$

| $E_{\text {x }}$ | $\begin{gathered} J^{\pi} ; T \\ (\mathrm{MeV} \pm \mathrm{keV}) \\ \hline \end{gathered}$ | $\begin{gathered} \tau_{\mathrm{m}}(\mathrm{fs}) \text { or } \\ \Gamma_{\mathrm{c} . \mathrm{m} .}(\mathrm{keV}) \end{gathered}$ | Decay | Reactions |
| :---: | :---: | :---: | :---: | :---: |
| 0 | $\frac{3}{2}^{-} ; \frac{1}{2}$ | stable |  | $1,2,6,7,9,13$, $14,15,16,17,19$, $23,24,25,26,27$, $28,29,30,31,32$, $33,34,35,36,37$, $38,39,40,41,42$, $43,44,45,46,47$, $48,49,50,51,52$, $53,54,55,56,57$, $58,59,60,61,62$, 63,64 |
| $2.124693 \pm 0.027$ | $\frac{1}{2}^{-}$ | $\tau_{\mathrm{m}}=5.5 \pm 0.4$ | $\gamma$ | $\begin{aligned} & 1,6,7,9,13,14, \\ & 15,16,17,23,24, \\ & 25,26,27,29,30, \\ & 32,33,36,37,38, \\ & 40,47,48,49,51, \\ & 52,53,55,58,59, \\ & 60,61,62,63,64 \end{aligned}$ |
| $4.44489 \pm 0.50$ | $\frac{5}{2}^{-}$ | $1.18 \pm 0.04$ | $\gamma$ | $1,2,6,7,9,13$, $14,15,19,23,24$, 25, 26, 27, 29, 30, $32,33,36,37,38$, 40, 47, 49, 51, 53, 59, 60, 61 |
| $5.02031 \pm 0.30$ | $\frac{3}{2}^{-}$ | $0.34 \pm 0.01$ | $\gamma$ | $1,6,7,9,14,15$, $23,24,25,26,27$, 29, 30, 32, 33, 36, 37, 38, 47, 48, 51, $52,53,55,59,60$, 61 |
| $6.7429 \pm 1.8$ | $\frac{7}{2}^{-}$ | $22 \pm 5$ | $\gamma$ | $1,2,6,14,15,19$, <br> $23,24,25,26,29$, <br> $33,36,37,38,47$, <br> $48,53,55,59,60$, 61 |

Table 11.3: Energy Levels of ${ }^{11} \mathrm{~B}$ (continued)

| $E_{\text {x }}$ | $\begin{gathered} J^{\pi} ; T \\ (\mathrm{MeV} \pm \mathrm{keV}) \end{gathered}$ | $\begin{gathered} \tau_{\mathrm{m}}(\mathrm{fs}) \text { or } \\ \Gamma_{\text {c.m. }}(\mathrm{keV}) \\ \hline \end{gathered}$ | Decay | Reactions |
| :---: | :---: | :---: | :---: | :---: |
| $6.79180 \pm 0.30$ | $\frac{1}{2}^{+}$ | $1.7 \pm 0.2$ | $\gamma$ | $\begin{aligned} & 1,2,6,14,15,23, \\ & 24,25,27,29,33, \\ & 37,40,47,48,51, \\ & 55,60 \end{aligned}$ |
| $7.28551 \pm 0.43$ | $\frac{5}{2}^{+}$ | $0.57 \pm 0.04$ | $\gamma$ | $1,2,6,13,14,15$, 23, 24, 25, 27, 29, 33, 38, 48, 53 |
| $7.97784 \pm 0.42$ | $\frac{3}{2}+$ | $0.57 \pm 0.06$ | $\gamma$ | $\begin{aligned} & 1,2,14,23,24 \\ & 27,29,33,48,53 \end{aligned}$ |
| $8.5603 \pm 1.8$ | $\left(\frac{3}{2}^{-}\right)$ | $0.70 \pm 0.07$ | $\gamma$ | 1, 13, 14, 23, 24, 29, 30, 33, 48, 53, 60, 61 |
| $8.9202 \pm 2.0$ | $\frac{5}{2}^{-}$ | $\Gamma=4.37 \pm 0.02 \mathrm{eV}$ | $\gamma, \alpha$ | $\begin{aligned} & 1,2,13,14,19, \\ & 23,24,26,29,30, \\ & 33,38,55,59,60, \\ & 61 \end{aligned}$ |
| $9.1850 \pm 2.0$ | $\frac{7}{2}^{+}$ | $1.9{ }_{-1.1}^{+1.5} \mathrm{eV}$ | $\gamma, \alpha$ | $\begin{aligned} & 1,2,14,23,24, \\ & 26,33,62 \end{aligned}$ |
| $9.2744 \pm 2$ | $\frac{5}{2}^{+}$ | 4 | $\gamma, \alpha$ | $\begin{aligned} & 1,2,14,23,24, \\ & 33,62 \end{aligned}$ |
| $9.82 \pm 25$ | $\left(\frac{1}{2}{ }^{+}\right)$ |  |  | 48 |
| $9.876 \pm 8$ | $\frac{3}{2}+$ | $110 \pm 15$ | $\alpha$ | 5,14, 27 |
| $10.26 \pm 15$ | $\frac{3}{2}^{-}$ | $150 \pm 25$ | $\gamma, \alpha$ | 2, 5, 14, 61 |
| $10.33 \pm 11$ | $\frac{5}{2}^{-}$ | $110 \pm 20$ | $\gamma, \alpha$ | 2, 5, 14, 24, 61 |
| $10.597 \pm 9$ | $\frac{7}{2}+$ | $100 \pm 20$ | $\gamma, \alpha$ | 2, 5, 14, 20, 22 |
| $10.96 \pm 50$ | $\frac{5}{2}^{-}$ | 4500 | $\alpha$ | 5 |
| $11.265 \pm 17$ | $\frac{9}{2}+$ | $110 \pm 20$ | $\alpha$ | 5, 14 |
| $11.444 \pm 19$ |  | $103 \pm 20$ | $\alpha$ | 5, 14 |
| $11.600 \pm 30$ | $\frac{5}{2}^{+}$ | $170 \pm 30$ | n, $\alpha$ | $\begin{aligned} & 3,5,14,20,22, \\ & 33,61 \end{aligned}$ |
| $11.886 \pm 17$ | $\frac{5}{2}^{-}$ | $200 \pm 20$ | n, $\alpha$ | 3, 5, 14, 20, 22 |
| $12.0 \pm 200$ | $\frac{7}{2}{ }^{+}$ | $\approx 1000$ | n, $\alpha$ | 5, 20, 22 |
| $12.557 \pm 16$ | $\frac{1}{2}^{+}\left(\frac{3}{2}+{ }^{+}\right) ; \frac{3}{2}$ | $210 \pm 20$ | $\gamma, \mathrm{p}, \alpha$ | 5, 14, 17, 18, 36 |

Table 11.3: Energy Levels of ${ }^{11} \mathrm{~B}$ (continued)

| $E_{\text {x }}$ | $\begin{gathered} J^{\pi} ; T \\ (\mathrm{MeV} \pm \mathrm{keV}) \end{gathered}$ | $\begin{gathered} \tau_{\mathrm{m}}(\mathrm{fs}) \text { or } \\ \Gamma_{\mathrm{c} . \mathrm{m} .}(\mathrm{keV}) \end{gathered}$ | Decay | Reactions |
| :---: | :---: | :---: | :---: | :---: |
| $12.916 \pm 12$ | $\frac{1}{2}^{-} ; \frac{3}{2}$ | $200 \pm 25$ | $\gamma, \mathrm{p}, \alpha$ | $\begin{aligned} & 5,14,17,18,33, \\ & 59,61 \end{aligned}$ |
| $13.137 \pm 40$ | $\frac{9}{}{ }^{-}$ | $426 \pm 40$ | $\mathrm{nt}, \alpha$ | 3, 14, 20, 21, 22 |
| 13.16 | $\frac{5}{2}^{+} ; \frac{7}{2}^{+}$ | 430 | n, $\alpha$ | 20, 22 |
| $14.04 \pm 100$ | $\frac{11}{2}^{+}$ | $500 \pm 200$ | n, $\alpha$ | 3, 20, 22 |
| $14.34 \pm 20$ | $\frac{5}{2}^{+} ; \frac{3}{2}$ | $254 \pm 18$ | $\gamma, \mathrm{p}$ | 14, 17, 36 |
| $14.565 \pm 15$ |  | $\leq 30$ | $\mathrm{n}, \mathrm{t}, \alpha$ | $\begin{aligned} & 3,14,20,21,22, \\ & 36,61 \end{aligned}$ |
| $15.29 \pm 25$ | $\left(\frac{3}{2}, \frac{5}{2}, \frac{7}{2}\right)^{+} ;\left(\frac{3}{2}\right)$ | $250 \pm 50$ | $\gamma, \mathrm{p}, \mathrm{n}, \alpha$ | 20, 22, 33, 61 |
| $16.437 \pm 20$ | $T=\frac{3}{2}$ | $\leq 30$ | $\mathrm{p}, \mathrm{d}, \alpha$ | $\begin{aligned} & 11,14,22,30,33, \\ & 61 \end{aligned}$ |
| 17.33 |  | $\approx 1000$ | $\mathrm{n}, \mathrm{d}, \mathrm{t}, \alpha$ | 11, 21, 22 |
| $17.43 \pm 50$ | $T=\frac{3}{2}$ | $100 \pm 30$ | $\gamma, \mathrm{n}, \mathrm{p}, \mathrm{d}, \alpha$ | 3, 9, 11, 14 |
| 18.0 | $T=\frac{3}{2}$ | $870 \pm 100$ |  | 14 |
| $18.37 \pm 50$ | $\left(\frac{1}{2}, \frac{3}{2}, \frac{5}{2}\right)^{+}$ | $260 \pm 80$ | $\gamma, \mathrm{d}$ | 9 |
| $19.13 \pm 30$ | ( $\pi=+$ ) ; $\frac{3}{2}$ | $115 \pm 25$ |  | 14, 61 |
| 19.7 | $\left(\frac{1}{2}^{+}\right)$ | broad | $\gamma, \mathrm{d}$ | 9, 28 |
| $21.27 \pm 50$ | $T=\frac{3}{2}$ | $300 \pm 30$ |  | 14 |
| 23.7 | $\left(\frac{1}{2}, \frac{3}{2}, \frac{5}{2}\right)^{+}$ |  | $\gamma, \mathrm{d}$ | 9 |
| 26.5 |  | broad | $\gamma, \mathrm{n}$ | 28 |

Ground-state properties of ${ }^{11} \mathrm{~B}$ : (84AN1B, 84ZI04, 85AN28, 85GO1A, 85HA18, 85FA01, 85ZI05, 86DO1E, 86GL1A, 86RO03, 86WI04, 87AB03, 87FU06, 87KI1C, 88AR1I, 88BI1A, 88VA03, 88WA08, 88WO04, 89SA10).

$$
\begin{gathered}
\mu=+2.688637(2) \mathrm{nm}(78 \mathrm{LEZA}), \\
Q=40.65(26) \mathrm{mb}[\mathrm{see}(80 \mathrm{AJ} 01)], \\
B\left(\mathrm{E} 2 ; \frac{3}{2}^{-} \rightarrow \frac{1}{2}^{-}\right)=2.6 \pm 0.4 e^{2} \cdot \mathrm{fm}^{4}(80 \mathrm{FE} 07) .
\end{gathered}
$$

Mass of ${ }^{11} \mathrm{~B}$ : The mass excess of ${ }^{11} \mathrm{~B}$ has been measured to be $9303.09 \pm 1.30 \mu \mathrm{u}$ (84EL05) [mass spectrometer]. The mass excess listed by (88WA18) is $8668.2 \pm 0.3 \mathrm{keV}$, and we adopt it.

Table 11.4: Electromagnetic transitions in ${ }^{11} \mathrm{~B}{ }^{\text {a }}$

| Initial <br> state |  | $J^{\pi}$ | $\begin{gathered} \hline \Gamma_{\gamma}(\text { total }) \\ (\mathrm{eV}) \end{gathered}$ | Branching ratios (\%) to final state |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | g.s. |  | 2.12 | 4.44 | 5.02 | 6.74 | 6.79 | 7.29 |
|  | $2.12{ }^{\text {b }}$ |  | $\frac{1}{2}^{-}$ | $0.120 \pm 0.009$ | 100 |  |  |  |  |  |  |
|  | $4.44{ }^{\text {b }}$ | $\frac{5}{2}$ | $0.56 \pm 0.02$ | $100{ }^{\text {c }}$ |  |  |  |  |  |  |
|  | $5.02{ }^{\text {b }}$ | $\frac{3}{2}$ | $1.963 \pm 0.067$ | $85.6 \pm 0.6{ }^{\text {d }}$ | $14.4 \pm 0.6{ }^{\text {e }}$ |  |  |  |  |  |
|  | $6.74{ }^{\text {b }}$ | $\frac{7^{-}}{}{ }^{-}$ | $0.030 \pm 0.007$ | $70 \pm 2^{\text {f }}$ | <3 | $30 \pm 2$ | $<1$ |  |  |  |
|  | $6.79{ }^{\text {b }}$ | $\frac{1}{2}+$ | $0.385 \pm 0.044$ | $67.5 \pm 1.1$ | $28.5 \pm 1.1$ | < 0.04 | $4.0 \pm 0.3$ |  |  |  |
|  | $7.29{ }^{\text {b }}$ | $\frac{5}{2}+$ | $1.149 \pm 0.080$ | $87.0 \pm 2.0$ | $<1$ | $5.5 \pm 1$ | $7.5 \pm 1$ |  |  |  |
| $\checkmark$ | $7.98{ }^{\text {b }}$ | $\frac{3}{2}+$ | $1.15 \pm 0.15$ | $46.2 \pm 1.1$ | $53.2 \pm 1.2$ | < 0.06 | <0.09 |  | $<0.10$ | $0.85 \pm 0.04$ |
| $\bigcirc$ | $8.56{ }^{\text {b }}$ | $\left(\frac{3}{2}^{-}\right)^{\mathrm{g}}$ | $0.946 \pm 0.090$ | $56 \pm 2$ | $30 \pm 2$ | $5 \pm 1$ | $9 \pm 1$ |  |  |  |
|  | $8.92{ }^{\text {b }}$ | $\frac{5}{2}^{-}$ | $4.368 \pm 0.021$ | $95 \pm 1^{\text {h }}$ | $<1$ | $4.5 \pm 0.5$ | $<1$ | $<1$ | $<1$ |  |
|  | $9.19{ }^{\text {i }}$ | $\frac{7}{2}+$ | $0.17_{-0.03}^{+0.06}$ | $0.9 \pm 0.3$ |  | $86.6 \pm 2.3$ |  | $12.5 \pm 1.1$ | <1.3 |  |
|  | $9.27^{\text {i }}$ | $\frac{5}{2}^{+}$ | $1.15 \pm 0.16$ | $18.4 \pm 0.9$ |  | $69.7 \pm 1.4$ |  | $11.9 \pm 0.6$ | < 0.6 |  |

[^0]Comments [mainly from (65OL03, 62GR07)]
(1) $4.44 \mathrm{MeV} .9 .28 \rightarrow 4.44 \rightarrow 0$ angular distribution fixes $J=\frac{5}{2}$. Odd parity determined from direct interaction assignments.
(2) 5.02 MeV . Internal pair correlation permit M1, E2 for the g.s. transition: $J^{\pi} \leq \frac{7}{2}^{-}$(parity from $l$-assignments). $\tau_{\mathrm{m}}$ excludes $\frac{7}{2}$, branch to 2.12 , $\frac{5}{2}$. Angular correlation fixes $\frac{3}{2}^{-}$.
(3) 6.74 MeV . Internal pairs indicate practically pure $\mathrm{E} 2 \mathrm{~g} . \mathrm{s}$. radiation. Angular distributions and branching ratios (and l-assignments) all lead to $\frac{7}{2}^{-}$.
(4) 6.79 MeV . The allowed $\beta$-decay from ${ }^{11} \mathrm{Be}\left[J^{\pi}=\frac{1}{2}^{+}\right]$requires $J^{\pi} \leq \frac{3}{2}^{+}$. The relatively strong $\gamma$-branch to ${ }^{11} \mathrm{~B}^{*}(2.12)$ favors $\frac{1}{2}^{+}, \frac{3}{2}^{+}$. All $\gamma^{\text {'s }}$ from this level are isotropic, suggesting $J^{\pi}=\frac{1}{2}^{+}$, but not excluding $\frac{3}{2}^{+}$.
(5) 7.29 MeV . The g.s. transition is mainly E1, so $J^{\pi} \leq \frac{5}{2}^{+}$. The assignment $\frac{1}{2}^{+}$is excluded by the strength of $(7.29 \rightarrow 4.44) . J^{\pi}=\frac{5}{2}^{+}$is consistent with $\log f t>8.04$ in the ${ }^{11} \mathrm{Be} \beta$-decay.
(6) 7.98 MeV. Transitions to ${ }^{11} \mathrm{~B}(0,2.12)$ are predominantly E1; thus ${ }^{11} \mathrm{~B}^{*}(7.98)$ has even parity, and the odd parity of ${ }^{11} \mathrm{~B}^{*}(2.12)$ is confirmed. The transition to ${ }^{11} \mathrm{~B}^{*}(2.12)$ is not isotropic, so $J^{\pi}=\frac{3}{2}^{+}$.
$\not{ }^{\leftharpoondown} \quad(7) 8.56 \mathrm{MeV}$. Correlation of internal pairs indicate that the g.s. transition is $\mathrm{M} 1+\mathrm{E} 2$ or $\mathrm{E} 1+\mathrm{M} 2, \mathrm{~J}^{\pi}=\leq \frac{5}{2}^{+}$or $\leq \frac{7}{2}^{+}$; the lifetime to ${ }^{11} \mathrm{~B}^{*}(2.12)$ excludes $\frac{7}{2}^{-}$. If the level has even parity, the required M2 admixture is excessive. $J^{\pi} \leq \frac{5}{2}^{-}$is favored. See also footnote ${ }^{\mathrm{i}}$ in Table 11.4.
(8) 8.92 MeV . From ${ }^{7} \operatorname{Li}(\alpha, \gamma)^{11} \mathrm{~B}, J^{\pi}=\frac{3}{2}^{+}, \frac{5}{2}^{+}, \frac{5}{2}^{-}$. The internal pair correlation confirms $\frac{5}{2}^{-}$. For higher states see comments under individual reactions and (68AJ02).

Isotopic abundance: $(80.1 \pm 0.2) \%$ ( 84 DE 1 A$)$.

1. ${ }^{6} \mathrm{Li}\left({ }^{6} \mathrm{Li}, \mathrm{p}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=12.215$

Angular distributions have been measured for the proton groups to the first eight states of ${ }^{11} \mathrm{~B}$ at $E\left({ }^{6} \mathrm{Li}\right)=2$ to 16 MeV (87DO05). For the earlier work see (80AJ01). For excitation functions see ${ }^{12} \mathrm{C}$. See also (87DO07).
2. ${ }^{7} \operatorname{Li}(\alpha, \gamma){ }^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=8.6637
$$

Resonances for capture radiation are displayed in Table 11.5. See also (84YA1A, 85CA41, 88BU01, 88CA26; astrophys.).
3. ${ }^{7} \operatorname{Li}(\alpha, \mathrm{n}){ }^{10} \mathrm{~B}$
$Q_{\mathrm{m}}=-2.7905$
$E_{\mathrm{b}}=8.6637$

The total cross section has been measured from threshold to $E_{\alpha}=5.67 \mathrm{MeV}$ [see also reaction 21]: a broad maximum at $E_{\alpha} \approx 5.1 \mathrm{MeV}\left(\sigma_{\max }=40 \mathrm{mb}\right)$ is observed (84OL05). For the earlier work see Tables 11.7 in (80AJ01) and (85AJ01). See also (85CA41; astrophys.).
4. ${ }^{7} \mathrm{Li}(\alpha, \mathrm{t})^{8} \mathrm{Be}$
$Q_{\mathrm{m}}=-2.5597$
$E_{\mathrm{b}}=8.6637$

Excitation functions have been measured for $E_{\alpha}=14$ to $25 \mathrm{MeV}\left(\mathrm{t}_{0}\right)$ and 18 to 25 MeV $\left(\mathrm{t}_{1}\right)$ : see (80AJ01). See also ${ }^{8}$ Be in (88AJ01) and (87DM1C).

## 5. ${ }^{7} \mathrm{Li}(\alpha, \alpha)^{7} \mathrm{Li}$

$$
E_{\mathrm{b}}=8.6637
$$

The elastic scattering and the scattering to ${ }^{7} \mathrm{Li}^{*}(0.48)$ have been studied at many energies to $E_{\alpha}=22.5 \mathrm{MeV}$ : see (75AJ02, 80AJ01, 85AJ01). Observed resonances are displayed in Table 11.6. For $\alpha-{ }^{7} \mathrm{Li}$ correlations see (87PO03) and the "General" section. See also (87BU27), (87EL1B; applied) and (85CH27; theor.).
6. ${ }^{7} \mathrm{Li}\left({ }^{6} \mathrm{Li}, \mathrm{d}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=7.189$

Angular distributions have been measured for $E\left({ }^{7} \mathrm{Li}\right)=3.3$ to 5.95 MeV : see (75AJ02).

Table 11.5: Resonances in ${ }^{7} \operatorname{Li}(\alpha, \gamma)^{11} \mathrm{~B}^{\text {a }}$

|  | $\begin{gathered} E_{\text {res }} \\ (\mathrm{keV}) \end{gathered}$ | $\begin{aligned} & \Gamma_{\text {c.m. }} \\ & (\mathrm{keV}) \end{aligned}$ | $\begin{gathered} { }^{11} \mathrm{~B}^{*} \\ (\mathrm{MeV}) \end{gathered}$ | $J^{\pi}$ | $\begin{gathered} \omega \gamma \\ (\mathrm{eV}) \\ \hline \end{gathered}$ | $\begin{aligned} & \Gamma_{\gamma_{0}} \\ & (\mathrm{eV}) \end{aligned}$ | Percentage decay to ${ }^{11} \mathrm{~B}^{*}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  | 0 | 4.44 | 6.74 | 6.79 |
|  | $401 \pm 3{ }^{\text {b }}$ | $4.37 \pm 0.02 \mathrm{eV}$ | 8.919 | $\frac{5}{2}^{-}$ | $(8.8 \pm 1.4) \times 10^{-3}$ | $4.15 \pm 0.02^{\text {c }}$ | $95 \pm 1$ | $4.5 \pm 0.5$ |  |  |
|  | $814 \pm 2^{\text {b }}$ | $1.8{ }_{-1.1}^{+1.5} \mathrm{eV}$ | 9.182 | $\frac{7^{+}}{}{ }^{+}$ | $0.310 \pm 0.047$ | $0.17_{-0.01}^{+0.05}{ }^{\text {d }}$ | $0.9 \pm 0.3$ | $90.8 \pm 4.0$ | $8.3 \pm 1.0$ | $<1.3$ |
|  | $953 \pm 2^{\text {b }}$ | 4 | 9.271 | $\frac{5}{2}^{+}$ | $1.72 \pm 0.24$ | $0.20 \pm 0.03{ }^{\text {c }}$ | $17.1 \pm 1.0$ | $71.7 \pm 1.8$ | $11.2 \pm 0.6$ | $<0.6{ }^{\text {e }}$ |
| $\leftharpoondown$ | $2500 \pm 20$ | 433 | 10.26 |  |  | 17 | f |  |  |  |
| $\omega$ | $2620 \pm 20$ | 100 | 10.33 |  |  | 1.0 | f |  |  |  |
|  | $2800 \pm 50$ | $\approx 140$ | 10.45 |  |  | $10 /(2 J+1)$ |  |  |  |  |
|  | (3040) | 90 | (10.60) |  |  | $<0.2$ | f |  |  |  |

a See Table 11.6 in (80AJ01) for comments and references.
${ }^{\mathrm{b}} \Gamma_{\alpha}(\mathrm{c} . \mathrm{m})=.(5.9 \pm 0.9) \times 10^{-3}, 1.6_{-1.1}^{+1.5}$, and $4 \times 10^{3} \mathrm{eV}$ for ${ }^{11} \mathrm{~B}^{*}(8.92,9.19,9.27)(84 \mathrm{HA} 13)$. See also Table 11.4
c See Table 11.4
${ }^{\mathrm{d}} \Gamma_{\gamma}$, not $\Gamma_{\gamma 0}$. See also Table 11.4.
${ }^{\text {e }}$ The decay to ${ }^{11} \mathrm{~B}^{*}(7.29,7.98)\left[J^{\pi}=\frac{5}{2}^{+}, \frac{3}{2}^{+}\right]$is also observed: $\approx 1 \%$ and $\approx 0.03 \%$ respectively.
f $<10 \%$ to ${ }^{11} \mathrm{~B}^{*}(2.12)$.
7. ${ }^{7} \mathrm{Li}\left({ }^{7} \mathrm{Li}, \mathrm{t}\right){ }^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=6.196
$$

Angular distributions have been measured at $E\left({ }^{7} \mathrm{Li}\right)=2.10$ to 5.75 MeV . At $E\left({ }^{7} \mathrm{Li}\right)=$ 79.6 MeV transitions are observed to several ${ }^{11} \mathrm{~B}$ states. ${ }^{11} \mathrm{~B}_{\text {g.s. }}$ is particularly strongly populated. See (75AJ02) for references.
8. ${ }^{8} \mathrm{Li}(\alpha, \mathrm{n}){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=6.6309$

See (88MA1U; astrophysics). See also (88SA2Q, 89BO1K).
9. ${ }^{9} \mathrm{Be}(\mathrm{d}, \gamma){ }^{11} \mathrm{~B} \quad Q_{\mathrm{m}}=15.8153$

The $90^{\circ} \gamma_{0}$ differential cross section has been measured for $E_{\mathrm{d}}=0.5$ to 11.9 MeV : see (75AJ02). The behavior of the $\gamma_{0}, \gamma_{1}$, and $\gamma_{2+3}$ total cross sections and of the angular distributions of these $\gamma$-rays indicate two resonances at $E_{\mathrm{d}}=1.98 \pm 0.05$ and $3.12 \pm 0.05 \mathrm{MeV}$ with $\Gamma_{\text {lab }}=225 \pm 50$ and $320 \pm 100 \mathrm{keV}$, corresponding to ${ }^{11} \mathrm{~B}^{*}(17.43,18.37)$. The higher resonance was not observable in the $\gamma_{2}+\gamma_{3}$ cross section which was not measured beyond $E_{\mathrm{d}}=2.5 \mathrm{MeV}$. The maximum $\gamma_{0}$ cross section observed is $10.1 \pm 3.5 \mu \mathrm{~b}$ at $E_{\mathrm{d}} \approx 0.96 \mathrm{MeV}$. Resonant behavior is observed in the $90^{\circ} \gamma_{0}$ cross section at $E_{\mathrm{d}} \approx 3.4$ and $9.65 \mathrm{MeV}\left({ }^{11} \mathrm{~B}^{*}(18.6,23.7)\right)$ in addition to a wide structure at $4.7 \mathrm{MeV}\left({ }^{11} \mathrm{~B}^{*}(19.7)\right)$. The angular distributions of $\gamma_{0}$ from ${ }^{11} \mathrm{~B}^{*}(18.6,23.7)$ are typical of E 1 transitions. The $\left(\mathrm{d}, \gamma_{0}\right)$ reaction appears to proceed via excitation of the $T=\frac{1}{2}$ component of the giant dipole resonance in ${ }^{11} \mathrm{~B}$.
10. ${ }^{9} \mathrm{Be}(\mathrm{d}, \mathrm{n}){ }^{10} \mathrm{~B} \quad Q_{\mathrm{m}}=4.3612 \quad E_{\mathrm{b}}=15.8153$

The cross section follows the Gamow function for $E_{\mathrm{d}}=70$ to 110 keV . The fast neutron and $\gamma$-yield rise smoothly to $E_{\mathrm{d}}=1.8 \mathrm{MeV}$ except for a possible "resonance" at $E_{\mathrm{d}} \approx 0.94 \mathrm{MeV}$. The fast neutron yield then remains approximately constant to 3 MeV : see (68AJ02) for references. The excitation functions for $n_{0} \rightarrow n_{4}$, and $n$ to ${ }^{10} B^{*}(5.1,6.57)$ have been measured for $E_{\mathrm{d}}=14$ to 16 MeV ; no strong fluctuations are observed: see (75AJ02). Thick target yields for $\gamma$-rays have been measured at $E_{\mathrm{d}}=48$ to 170 keV : see ( 85 AJ 01 ). Thick target yields are also reported at $E_{\mathrm{d}}=14.8,18.0$ and 23.0 MeV : see (80AJ01). Polarization measurements have been carried out at $E_{\mathrm{d}}=0.4$ to 5.5 MeV [see (75AJ02, 80AJ01)] and at $E_{\overrightarrow{\mathrm{d}}}=12.3 \mathrm{MeV}$ : see (85AJ01). See also ${ }^{10} \mathrm{~B}$ in (88AJ01), (85SM08, 86BA40) and (88ZVZZ; theor.).

Table 11.6: Structure in ${ }^{7} \mathrm{Li}(\alpha, \alpha)^{7} \mathrm{Li}$ and ${ }^{7} \mathrm{Li}\left(\alpha, \alpha^{\prime}\right)^{7} \mathrm{Li}^{\text {a }}$

| $E_{\alpha}{ }^{\mathrm{b}}$ <br> $(\mathrm{keV})$ | $E_{\alpha}{ }^{\mathrm{c}}$ <br> $(\mathrm{keV})$ | $\Gamma_{\text {c.m. }}$ <br> $(\mathrm{keV})$ | $E_{\mathrm{x}}$ <br> $(\mathrm{MeV} \pm \mathrm{keV})$ | $J^{\pi}$ |
| :---: | :---: | :---: | :---: | :---: |
| $1900 \pm 10$ |  | $130 \pm 30$ | $9.873 \pm 10$ | $\frac{3}{2}^{{ }^{+}}$ |
| $2480 \pm 50$ |  | $150 \pm 40$ | $10.24 \pm 50$ | $\frac{3}{2}^{(-)}, \frac{1}{2}$ |
| $3040 \pm 10$ | $3630 \pm 30$ | $80 \pm 30$ | $10.34 \pm 30$ | $\frac{5}{2}^{-}, \frac{7}{2}^{2}$ |
| $3600 \pm 50$ |  | $70 \pm 10$ | $10.599 \pm 10$ | $\frac{7}{2}^{+}$ |
|  | $4120 \pm 30$ | $90 \pm 50$ | $10.96 \pm 50$ | $\frac{5}{2}^{-}$ |
| $4430 \pm 50$ | 4430 |  | $11.29 \pm 30$ | $\frac{9}{2}^{+}$ |
| $4600 \pm 50$ |  | $150 \pm 50$ | $11.59 \pm 50$ |  |
| $5050 \pm 30$ |  | $150 \pm 50$ | $11.88 \pm 30$ |  |
|  | $5300 \pm 200$ | $\approx 1000$ | $12.0 \pm 200$ |  |
| $6100 \pm 30$ | $5500 \pm 100$ | $60 \pm 50$ | $(12.17 \pm 100)^{\mathrm{d}}$ |  |
| $6850 \pm 60$ |  | $150 \pm 50$ | $12.55 \pm 30$ |  |
| $(7200 \pm 50)^{\mathrm{e}}$ |  | $270 \pm 50$ | $13.03 \pm 60$ |  |
| $(8450 \pm 200)^{\mathrm{f}}$ | $7800 \pm 100$ | $500 \pm 200$ | $(13.25 \pm 50)^{\mathrm{d}}$ |  |
| $(9450 \pm 200)^{\mathrm{f}}$ |  | $500 \pm 200$ | $(13.63 \pm 100)^{\mathrm{d}}$ |  |
|  |  | $\leq 250$ | $(14.0 \pm 200)$ |  |
| $(11200 \pm 200)^{\mathrm{f}}$ | $9950 \pm 20$ | $500 \pm 200$ | $(15.00 \pm 20)^{\mathrm{d}}$ |  |
|  |  | $(15.8 \pm 200)$ |  |  |

[^1]11. (a) ${ }^{9} \operatorname{Be}(d, p){ }^{10} \mathrm{Be}$
$Q_{\mathrm{m}}=4.5874 \quad E_{\mathrm{b}}=15.8153$
(b) ${ }^{9} \mathrm{Be}(\mathrm{d}, \alpha)^{7} \mathrm{Li}$
$Q_{\mathrm{m}}=7.152$
(c) ${ }^{9} \mathrm{Be}(\mathrm{d}, \mathrm{t})^{8} \mathrm{Be}$
$Q_{\mathrm{m}}=4.5919$
Measurements of proton yields have been carried out at $E_{\mathrm{d}}$ up to 6.0 MeV for $\mathrm{p}_{0}$ and $\mathrm{p}_{1}$ [see (75AJ02, 80AJ01, 85AJ01)]. The $\mathrm{p}_{0}$ and $\mathrm{p}_{1}$ yields show a resonance at $E_{\mathrm{d}}=750 \pm 15 \mathrm{keV}$ $\left[{ }^{11} \mathrm{~B}^{*}(16.43), \Gamma \approx 40 \mathrm{keV}\right]$ and the $\mathrm{p}_{1}$ yield resonates at $1.85 \mathrm{MeV}\left[{ }^{11} \mathrm{~B}^{*}(17.33), \Gamma_{\mathrm{c} . \mathrm{m} .} \approx\right.$ $1.0 \mathrm{MeV}]$ and $2.3 \mathrm{MeV}\left[{ }^{11} \mathrm{~B}^{*}(17.70)\right.$, sharp]. See also (75AJ02, 85AJ01) for other possible structures. Polarization of the protons has been measured at $E_{\mathrm{d}}=1$ to 21 MeV [see (75AJ02, 80AJ01, 85AJ01)] and at $E_{\overrightarrow{\mathrm{d}}}=2.0$ to 2.8 MeV (84DE46; VAP; p $\mathrm{p}_{0}, \mathrm{p}_{1}$ ). See also ${ }^{10} \mathrm{Be}$ in (88AJ01) and (84AN1D).

The yield of $\alpha$-particles (reaction (b)) has been measured for $E_{\mathrm{d}}=0.3$ to 14.43 MeV [see (75AJ02, 80AJ01, 85AJ01)]. The 0.75 MeV resonance, observed in reaction (a), is weakly populated in the $\alpha_{0}$ yield. For polarization measurements see (85AJ01) and (84AN1D: $E_{\mathrm{d}}=2.0$ to $\left.2.8 \mathrm{MeV} ; \alpha_{0+1} ; \mathrm{VAP}\right)$. See also ${ }^{7} \mathrm{Li}$ in (88AJ01).

The cross section for reaction (c) has been measured for $E_{\mathrm{d}}=0.15$ to 19 MeV : see (68AJ02, 75AJ02, 80AJ01). Polarization meaurements are reported at $E_{\overrightarrow{\mathrm{d}}}=12$ and 15 MeV [see (80AJ01)] and at $E_{\overrightarrow{\mathrm{d}}}=2.0$ to 2.8 MeV (84AN1D; $\mathrm{t}_{0}$ ). There is no clear evidence of resonance structure. See also ${ }^{8} \mathrm{Be}$ in (88AJ01).
12. ${ }^{9} \mathrm{Be}(\mathrm{d}, \mathrm{d}){ }^{9} \mathrm{Be}$

$$
E_{\mathrm{b}}=15.8153
$$

Excitation functions for elastically scattered deuterons have been measured for $E_{\mathrm{d}}=0.4$ to 7.0 MeV and for 12.17 to 14.43 MeV (also $\mathrm{d}_{1}, \mathrm{~d}_{2}$ ) [see (75AJ02, 80AJ01)]. Polarization measurements have been reported at $E_{\overrightarrow{\mathrm{d}}}=6.3$ to 15 MeV [see (75AJ02, 80AJ01)] and at $E_{\mathrm{d}}=2.0$ to 2.8 MeV (83DE50; $\left.\mathrm{d}_{0} ; \mathrm{VAP}\right)$. See also ${ }^{9}$ Be in (88AJ01).
13. ${ }^{9} \mathrm{Be}(\mathrm{t}, \mathrm{n}){ }^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=9.5580
$$

Angular distributions have been measured at $E_{\mathrm{t}}=1.1$ to $1.7 \mathrm{MeV}\left(\mathrm{n}_{0}, \mathrm{n}_{1}, \mathrm{n}_{2}, \mathrm{n}_{6}, \mathrm{n}_{8}\right.$, $\mathrm{n}_{9}$ ): see (80AJ01).
14. ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=10.3218$

Observed proton groups are displayed in Table 11.7. Angular distributions have been obtained at a number of energies in the range $E\left({ }^{3} \mathrm{He}\right)=1.0$ to 38 MeV [see (80AJ01, 85AJ01)] and at 3 to 6 MeV (81LI1C; nine groups; DWBA). It is suggested that the $T=\frac{1}{2}$ strength is strongly fragmented (82ZW02). See also (85AJ01), ${ }^{12} \mathrm{C}, ~(85 \mathrm{MC1C}$; applied) and (88KH11; theor.).

Table 11.7: Energy levels of ${ }^{11} \mathrm{~B}$ from ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$

| $\begin{gathered} E_{\mathrm{x}}{ }^{\mathrm{a}} \\ \mathrm{MeV} \pm \mathrm{keV}) \end{gathered}$ | $\begin{gathered} E_{\mathrm{x}}{ }^{\mathrm{b}} \\ \mathrm{MeV} \pm \mathrm{keV}) \end{gathered}$ | $\begin{aligned} & \Gamma_{\text {c.m. }}{ }^{b}{ }^{b}(\mathrm{keV}) \end{aligned}$ | $L$ |
| :---: | :---: | :---: | :---: |
| 0 |  |  | 0 |
| $2.1243 \pm 0.9$ |  |  | 0 |
| $4.4434 \pm 1.8$ |  |  | 0 |
| $5.0187 \pm 2.3$ |  |  | 0 |
| $6.7411 \pm 3.0$ |  |  |  |
| $6.7909 \pm 3.1$ |  |  | 1 |
| $7.285 \pm 10$ |  |  |  |
| $7.975 \pm 10$ |  |  |  |
| $8.553 \pm 10$ |  |  | 0 |
| $8.909 \pm 10$ | $8.934 \pm 15$ |  | $0+2$ |
| $9.175 \pm 10$ | $9.183 \pm 15$ |  | (1) +3 |
| $9.264 \pm 10$ | $9.265 \pm 15$ | $10 \pm 10$ | $1+3$ |
| $9.86 \pm 20$ | $9.887 \pm 15$ | $104 \pm 15$ | 1 |
|  | $10.265 \pm 25$ | $168 \pm 25$ | 2 |
|  | $10.337 \pm 20$ | $123 \pm 20$ | $0+2$ |
|  | $10.580 \pm 20$ | $122 \pm 20$ | $1+3$ |
|  | $11.254 \pm 20$ | $110 \pm 20$ | 3 |
|  | $11.437 \pm 20$ | $103 \pm 20$ | $(0+2)$ |
|  | $11.588 \pm 30$ | $180 \pm 30$ | $1+3$ |
|  | $11.889 \pm 20$ | $204 \pm 20$ | $0+2$ |
|  | $12.563 \pm 20^{\text {c }}$ | $202 \pm 25$ | 1 |
|  | $12.920 \pm 20^{\text {c }}$ | $155 \pm 25$ | 2 |
|  | $13.137 \pm 40$ | $426 \pm 40$ | $1+3$ |
|  | $\equiv 14.40{ }^{\text {d }}$ | $261 \pm 25$ | $1+3$ |
|  | $14.565 \pm 15$ | $\leq 30$ | (1) |
|  | $16.437 \pm 20^{\text {c,e }}$ | $\leq 30$ |  |
|  | $\equiv 17.69{ }^{\text {c,e }}$ | $91 \pm 25$ | $(0+2)$ |
|  | $18.0 \pm 100^{\text {c,e }}$ | $870 \pm 100$ | $(1+3)$ |
|  | $19.146 \pm 30^{\mathrm{c}, \mathrm{e}}$ | $115 \pm 25$ | 3 |
|  | $21.27 \pm 50^{\text {c }}$ | $300 \pm 30$ | $(1+3)$ |

${ }^{\text {a }}$ See Table 11.9 in (80AJ01) for references and Table 11.14 here.
${ }^{\text {b }} E\left({ }^{3} \mathrm{He}\right)=38 \mathrm{MeV}$; DWBA analysis.
${ }^{\text {c }} T=\frac{3}{2}$ state.
${ }^{\mathrm{d}}$ This state may have mixed isospin $\left(T=\frac{1}{2}+T=\frac{3}{2}\right)$.
e Not observed in ${ }^{9} \operatorname{Be}(\alpha, d){ }^{11} B$.
15. ${ }^{9} \operatorname{Be}(\alpha, d){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-8.0314$

Angular distributions have been measured at a number of energies in the range $E_{\alpha}=23.4$ to 28.3 MeV [see (80AJ01)] and at 30.2 MeV (84VA07; $\mathrm{d}_{0} \rightarrow \mathrm{~d}_{3}$ ). The predominant $L$ transfers are $L=0,2 ; 0 ; 0$ for ${ }^{11} \mathrm{~B}^{*}(0,2.12,5.02)$. The angular distribution to ${ }^{11} \mathrm{~B}^{*}(4.44)$ is flat at $E_{\alpha}=27 \mathrm{MeV}$. At $E_{\alpha}=48 \mathrm{MeV},{ }^{11} \mathrm{~B}^{*}(16.44,17.69,18.0,19.15)$ are not excited suggesting that these states are rather pure $T=\frac{3}{2}$ states (82ZW02): see Table 11.7.
16. ${ }^{9} \mathrm{Be}\left({ }^{6} \mathrm{Li}, \alpha\right){ }^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=14.3403
$$

Angular distributions have been determined for seven $\alpha$-groups at $E\left({ }^{6} \mathrm{Li}\right)=3$ to 4 MeV , and at 24 MeV to ${ }^{11} \mathrm{~B}^{*}(0,2.12)$ and to a number of unresolved levels with $E_{\mathrm{x}} \leq 13.2 \mathrm{MeV}$ : see (68AJ02, 75AJ02). For the breakup reactions see (75AJ02).
17. ${ }^{10} \mathrm{Be}(\mathrm{p}, \gamma){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=11.2279$

The yield of $\gamma_{0}$ has been measured at $90^{\circ}$ for $E_{\mathrm{p}}=0.6$ to 6.3 MeV . Observed resonances are displayed in Table 11.8. $T=\frac{3}{2}$ assignments are made for the states at $E_{\mathrm{x}}=12.56,12.91$, 14.33 and 15.32 MeV whose energies match those of the first four states of ${ }^{11} \mathrm{Be}$ [compare with the $T=\frac{3}{2}$ states reported in ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ - Table 11.7]. See also Table 11.14. Several known $T=\frac{1}{2}$ states in ${ }^{11} \mathrm{~B}$ are not observed in this reaction: see Table 11.3.
18. ${ }^{10} \mathrm{Be}(\mathrm{p}, \mathrm{n}){ }^{10} \mathrm{~B}$
$Q_{\mathrm{m}}=-0.2262$
$E_{\mathrm{b}}=11.2279$

The reaction cross section has been measured for $E_{\mathrm{p}}=0.89$ to 1.93 MeV : the excitation of ${ }^{11} \mathrm{~B}^{*}(12.56,12.91)$ is reported ( 86 TE 1 A and G.M. Ter-Akopian, private communication; 87ER1D). See also (88DUO6; theor.).

Table 11.8: Levels of ${ }^{11} \mathrm{~B}$ from the ${ }^{10} \mathrm{Be}\left(\mathrm{p}, \gamma_{0}\right)^{11} \mathrm{~B}$ reaction (70GO04)

| $E_{\mathrm{p}}$ <br> $(\mathrm{MeV} \pm \mathrm{keV})$ | $E_{\mathrm{x}}$ <br> $(\mathrm{MeV})$ | $\Gamma_{\text {c.m. }}$ <br> $(\mathrm{keV})$ | $\left(J+\frac{1}{2}\right)$ <br> $\left(\Gamma_{\mathrm{p}} / \Gamma\right) \Gamma_{\gamma_{0}} \mathrm{a}$ <br> $(\mathrm{eV})$ | $\Gamma_{\gamma_{0}}{ }^{\mathrm{a}}$ <br> $(\mathrm{eV})$ | $\Gamma_{\gamma_{1} / \Gamma_{\gamma_{0}}}$ | $J^{\pi}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $(1.05 \pm 40)^{\mathrm{b}}$ | $(12.18)$ | $230 \pm 90$ | $3.1_{-2.0}^{+2.9}$ |  |  |  |
| $1.46 \pm 30$ | 12.56 | $230 \pm 65$ | $10_{-5}^{+7}$ | $10_{-5}^{+7}$ | $0.25 \pm 0.08$ | $\frac{1}{2}^{+}\left(\frac{3}{2}^{+}\right)$ |
| $1.85 \pm 20$ | 12.91 | $235 \pm 27$ | $29 \pm 9$ | $29 \pm 9{ }^{\mathrm{c}}$ | $\leq 0.06$ | $\frac{1}{2}^{-}$ |
| $3.41 \pm 20$ | 14.33 | $255 \pm 36$ | $29 \pm 9$ | $14.5 \pm 4.3$ | $\leq 0.1$ | $\frac{5}{2}^{(+)}\left(\frac{3}{2}^{-}\right)$ |
| $4.5 \pm 100$ | 15.32 | $635 \pm 180$ | $53_{-26}^{+34}{ }^{\mathrm{d}}$ |  |  |  |

${ }^{\text {a }}$ Values reported in (70GO04) are here shown multiplied by 1.7: see (73GO09). See also Table 11.14
${ }^{\mathrm{b}}$ May be due to ${ }^{10} \mathrm{~B}^{*}(0.7)+\mathrm{n}$ threshold.
${ }^{c}$ In the ( $\mathrm{e}, \mathrm{e}^{\prime}$ ) work of (75KA02) a strong group is observed at $E_{\mathrm{x}}=13.0 \pm 0.1 \mathrm{MeV}$. If it corresponds to the excitation of ${ }^{11} \mathrm{~B}^{*}(12.91)$ with $J^{\pi}=\frac{1^{-}}{2} ; T=\frac{3}{2}$, then $\Gamma_{\gamma_{0}}=36 \pm 7 \mathrm{eV}$ (75KA02).
${ }^{\mathrm{d}}$ Assumes that $\sigma_{\text {total }}=4 \pi \mathrm{~d} \sigma / \mathrm{d} \Omega\left(90^{\circ}\right)$.

Table 11.9: Neutron capture $\gamma$-rays from ${ }^{10} \mathrm{~B}+\mathrm{n}{ }^{\text {a }}$

| $E_{\gamma}(\mathrm{keV})$ | $I_{\gamma}{ }^{\mathrm{b}}$ | $I_{\gamma}{ }^{\mathrm{c}}$ | Assignment | $E_{\mathrm{x}}(\mathrm{keV})$ |
| :---: | :---: | :---: | :---: | :---: |
| $11447.35 \pm 0.52$ | $4.6 \pm 0.3$ | $4.7 \pm 0.3$ | capt. $\rightarrow$ g.s. |  |
| $8916.80 \pm 0.27$ | $13 \pm 1$ | $13.4 \pm 0.9$ | $8.92 \rightarrow \mathrm{~g} . \mathrm{s}$. | $8920.44 \pm 0.27$ |
| $6738.34 \pm 0.50$ | $19 \pm 2$ | $19.0 \pm 0.9$ | $6.74 \rightarrow$ g.s. | $6741.76 \pm 0.24$ |
| $4444.03 \pm 0.12$ | $67 \pm 4$ | $65.7 \pm 2.4$ | $4.44 \rightarrow \mathrm{~g} . \mathrm{s}$. | $4444.95 \pm 0.15$ |
| $7006.75 \pm 0.10$ | $56 \pm 2$ | $55.4 \pm 1.7$ | capt. $\rightarrow 4.44$ |  |
| $4711.17 \pm 0.10$ | $28 \pm 2$ | $25.6 \pm 0.9$ | capt. $\rightarrow 6.74$ |  |
| $2533.49 \pm 0.23$ | $12 \pm 4$ | $14.4 \pm 1.8$ | capt. $\rightarrow 8.92$ |  |
| $2296.61 \pm 0.59$ | $7 \pm 4$ | $8.9 \pm 2.4$ | $6.74 \rightarrow 4.44$ |  |

${ }^{\text {a }}$ ( 86 KO 19 ). For the earlier work see Table 11.12 in (75AJ02): $I_{\gamma}$ for $5.02 \rightarrow$ g.s. and $2.12 \rightarrow$ g.s. are $<2$ and $<3$, respectively ( 67 TH 05 ).
${ }^{\mathrm{b}}$ Photons/ 100 captures.
${ }^{\text {c }}$ Adopted: weighted mean of ( 67 TH 05 ) and ( 86 KO 19 ).

Table 11.10: Resonances in ${ }^{10} \mathrm{~B}+\mathrm{n}^{\text {a }}$

| ${ }^{10} \mathrm{~B}\left(\mathrm{n}, \mathrm{n}^{\prime} \gamma\right)^{10} \mathrm{~B}$ |  | ${ }^{10} \mathrm{~B}(\mathrm{n}, \alpha)^{7} \mathrm{Li}$ |  | Yield <br> of | ${ }^{11} \mathrm{~B}^{*}$ <br> $(\mathrm{MeV})$ |
| :---: | :---: | :---: | :---: | :---: | :--- |
| $E_{\text {res }}(\mathrm{MeV})$ | $\Gamma(\mathrm{keV})$ | $E_{\text {res }}(\mathrm{MeV})$ | $\Gamma(\mathrm{keV})$ | $\sigma_{\mathrm{t}}, \alpha$ | 11.66 |
|  |  | $0.23^{\mathrm{b}}$ |  | $\sigma_{0}, \alpha_{1}$ | 11.94 |
| 1.93 | 260 | $0.53^{\mathrm{b}, \mathrm{c}}$ | 140 | 1.86 | 570 |
| $(2.6)$ | broad | 2.79 | 530 | $\sigma_{\mathrm{t}}, \alpha_{0}, \alpha_{1}, \mathrm{t}, \mathrm{n}^{\prime}$ | 13.2 |
| 3.31 | 370 | 3.43 | $<120$ | $\alpha_{0}, \alpha_{1}, \mathrm{n}^{\prime}$ | 14.0 |
| 4.1 |  | 4.1 | 800 | $\sigma_{\mathrm{t}}, \alpha_{0}, \alpha_{1}, \mathrm{n}^{\prime}$ | 14.57 |
| 4.73 |  |  |  | $\mathrm{n}^{\prime}$ | 15.2 |
|  |  | 5.7 | broad | $\alpha_{0}, \mathrm{t}$ | 16.6 |
|  |  | 6.4 | broad | $\alpha_{0}, \mathrm{t}$ | 17.3 |

${ }^{\text {a }}$ See also Table 11.11 For references see Table 11.12 in (80AJ01).
b (84OL05) [see reaction 21] report $E_{\mathrm{R}}=241 \pm 18$ and $493 \pm 4 \mathrm{keV}, \Gamma=166 \pm 40$ and $194 \pm 6 \mathrm{keV}: E_{\mathrm{x}}$ are then 11.673 and 11.902 MeV .
${ }^{c}$ See footnote ${ }^{\mathrm{b}}$ in Table 11.11

19. ${ }^{10} \mathrm{~B}(\mathrm{n}, \gamma){ }^{11} \mathrm{~B} \quad$|  | $Q_{\mathrm{m}}=11.4542$ |
| :--- | :--- |
|  | $Q_{0}=11454.1 \pm 0.2 \mathrm{keV}(86 \mathrm{KO} 19)$ |

The thermal capture cross section is $0.29 \pm 0.04 \mathrm{~b}$ ( 86 KO 19 ). The observed capture $\gamma$-rays are displayed in Table 11.9. See also (88MU05; theor.).
20. ${ }^{10} \mathrm{~B}(\mathrm{n}, \mathrm{n}){ }^{10} \mathrm{~B}$

$$
E_{\mathrm{b}}=11.4542
$$

The scattering amplitude (bound) $a=-0.2 \pm 0.4 \mathrm{fm}$, the total scattering cross section $\sigma($ free $)=2.23 \pm 0.06 \mathrm{~b}$ ( 83 KO 17 ). The total scattering cross section is constant at $2.23 \pm$ 0.06 b for $E_{\mathrm{n}}=0.7$ to 10 keV and then rises to 2.97 b at $E_{\mathrm{n}}=127 \mathrm{keV}$. For a display of cross sections and a listing of measurements see ( 88 MCZT ).

Total cross section measurements in the range $E_{\mathrm{n}}=10$ to 500 keV show a broad maximum near $E_{\mathrm{n}}=0.23 \mathrm{MeV}$, also observed in the ( $\mathrm{n}, \alpha$ ) cross section. At higher energies the total cross section shows broad maxima at $E_{\mathrm{n}}=1.9,2.8$ and 4.3 MeV : see Table 11.10. In the range $E_{\mathrm{n}}=5.5$ to $16 \mathrm{MeV} \sigma_{\text {tot }}$ is constant at 1.5 b .

Polarization measurements ( 0.075 to 2.2 MeV and 2.63 MeV ) and measurements of differential cross sections ( 0.075 to 4.4 MeV ) have been analyzed using $R$-matrix calculations:

Table 11.11: $R$-matrix analysis of resonant state in ${ }^{10} \mathrm{~B}+\mathrm{n}^{\text {a }}$

| $E_{\mathrm{n}}$ | $E_{\text {x }}$ | $J^{\pi}$ | $l_{\text {n }}$ | $\Gamma_{n}$ | $\Gamma_{\alpha_{0}}$ | $\Gamma_{\alpha_{1}}$ | $\begin{aligned} & \Gamma_{\mathrm{c} . \mathrm{m} .} \\ & (\mathrm{keV}) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| (MeV) | (MeV) |  |  | (c.m., MeV) |  |  |  |
|  | 10.60 | 2 | 0 | 0.120 | 0.030 | 0.070 | 220 |
| 0.17 | 11.61 | + | 0 | 0.004 | 0.296 | 0.0 | 300 |
| 0.37 | 11.79 | $\frac{7}{2}+$ | 0 | 0.770 | 0.001 | 0.113 | 884 |
| $0.53{ }^{\text {b }}$ | 11.94 | $\frac{5}{2}$ | 1 | 0.031 | 0.080 | 0.090 | 201 |
| 1.83 | 13.12 | 2 | 1 | 0.100 | 0.275 | 0.050 | 425 |
| 1.88 | 13.16 | $\frac{5}{2}^{+}, \frac{7}{2}^{+}$ | 2 | 0.080 | 0.200 | 0.150 | 430 |
| 2.82 | 14.02 | $\frac{11}{2}{ }^{+}$ | 2 | 0.800 | 0.045 | 0.010 | 855 |
| 4.2 | 15.3 | $\left(\frac{3}{2}, \frac{5}{2}, \frac{7}{2}\right)^{+}$ | 2 | 0.500 | 0.100 | 0.100 | 700 |

${ }^{\text {a }}$ Analysis based on polarization and differential cross-section measurements of the elastic scattering, and on results from ${ }^{10} \mathrm{~B}\left(\mathrm{n}, \alpha_{0}\right)$ and ( $\mathrm{n}, \alpha_{1}$ ). The analysis used a two-level, four-channel $R$-matrix formalism with a non-diagonal background $R$-matrix: see (73HA64). This analysis does not include ${ }^{11} \mathrm{~B}^{*}(14.53)$ because the resonance is weak, narrow and almost entirely in the $\alpha$-channel (73CO05). See also Table 11.10
b (78LA23) report $E_{\text {res }}=495 \pm 5 \mathrm{keV}, \Gamma=140 \pm 15 \mathrm{keV}, \sigma_{\max }[\mathrm{in}(\mathrm{n}$, $\left.\left.\alpha_{1} \gamma\right)\right]=94 \pm 6 \mathrm{mb}$.
the results are shown in Table 11.11. They are consistent with results from ${ }^{10} \mathrm{~B}\left(\mathrm{n}, \mathrm{n}^{\prime} \gamma\right)$ and ${ }^{7} \operatorname{Li}(\alpha, n)$. See (80AJ01) for references.

Elastic and inelastic cross sections have also been reported at $E_{\mathrm{n}}=4$ to 14.1 MeV [see (80AJ01)], at $E_{\mathrm{n}}=3.0$ to 12.0 MeV (86SAZR, 87SAZX; prelim.), at 8.0 to 13.9 MeV (82GL02) and at 10 to 17 MeV (86MU1D; also polarization measurements at 10 and 15 MeV ; prelim.). The yield of $0.7 \mathrm{MeV} \gamma$-rays has been studied from threshold to $E_{\mathrm{n}}=5.2 \mathrm{MeV}$ : observed resonances are displayed in Table 11.10. Inelastic scattering cross sections for formation of various ${ }^{10} \mathrm{~B}$ states have been measured at a number of energies in the range $E_{\mathrm{n}}=1.45$ to 14.8 MeV : see (75AJ02). See also ${ }^{10} \mathrm{~B}$ in (88AJ01), (86BAYL, 86DR1D), (83GO1H, 88MA1H), (88RE09; computer code) and (85CH27, 88HAZT; theor.).
21.
(a) ${ }^{10} \mathrm{~B}(\mathrm{n}, \mathrm{p})^{10} \mathrm{Be}$
$Q_{\mathrm{m}}=0.2262$
$E_{\mathrm{b}}=11.4542$
(b) ${ }^{10} \mathrm{~B}(\mathrm{n}, \mathrm{t})^{4} \mathrm{He}^{4} \mathrm{He}$
$Q_{\mathrm{m}}=0.3226$

The thermal cross section for reaction (a) is $6.4 \pm 0.5 \mathrm{mb}$ (87LA16); that for reaction (b) is $4.47 \pm 0.15 \mathrm{mb}$ (89CL01) [see also for other references], $7 \pm 2 \mathrm{mb}$ (87KA32). The cross section for reaction (b) has also been studied for $E_{\mathrm{n}}=1.4$ to 8.2 MeV [see Table 11.10 and (68AJ02)] and 3 to 8 MeV (86QA1A; prelim.). For various breakup processes see (84TU02). For a display of cross sections and a listing of measurements see (88MCZT). See also (85BO1D, 88MA1H, 88SUZY).
22. ${ }^{10} \mathrm{~B}(\mathrm{n}, \alpha)^{7} \mathrm{Li}$

$$
Q_{\mathrm{m}}=2.7905
$$

$$
E_{\mathrm{b}}=11.4542
$$

The "recommended" value of the thermal isotopic absorption cross section is $3837 \pm 9 \mathrm{~b}$ (81MUZQ). The $\alpha_{0} / \alpha_{1}$ branching for thermal neutrons is ( $6.723 \pm 0.011$ ) \% [mean of values listed in (85AJ01)]. At $E_{\mathrm{n}}=2$ and 24 keV the values are $(7.05 \pm 0.16) \%$ and $(7.13 \pm 0.15) \%$, respectively (79ST1B).

The cross section for this reaction has been measured for $E_{\mathrm{n}}=0.025 \mathrm{eV}$ to 14.8 MeV [see (75A.J02, 80A.J01, 85AJ01)]: for observed and deduced structures see Tables 11.10 and 11.11. For a display of cross sections and a listing of measurements see ( 88 MCZT ). For a review see (86CA28). "Detailed balance" [from ${ }^{7} \mathrm{Li}(\alpha, \mathrm{n})$ measurements] has led to the determination of the ${ }^{10} \mathrm{~B}\left(\mathrm{n}, \alpha_{0}\right)$ cross section from $0<E_{\mathrm{n}} \leq 0.78 \mathrm{MeV}$ : two resonances are inferred at $E_{\mathrm{R}}=241 \pm 18$ and $493 \pm 4 \mathrm{keV}$, with $\sigma_{\mathrm{R}}=17 \pm 3$ and $112 \pm 3 \mathrm{mb}$ and $\Gamma=166 \pm 40$ and $194 \pm 6 \mathrm{keV}$ (84OL05).

A study of the reaction involving polarized thermal neutrons and a polarized ${ }^{10} \mathrm{~B}$ target shows that the transition to ${ }^{7} \mathrm{Li}^{*}(0.48)$ proceeds almost totally through the $J=\frac{7}{2}$ channel (86KO19). The ratio of the ${ }^{10} \mathrm{~B}(\mathrm{n}, \alpha)$ cross section to the ${ }^{6} \mathrm{Li}(\mathrm{n}, \mathrm{t})$ cross section has been measured from $E_{\mathrm{n}} \approx 1$ to 45 eV (86CA29; prelim.).

Parity violation has been studied using polarized thermal neutrons: the $P$-odd asymmetries for the transitions to ${ }^{7} \mathrm{Li}^{*}(0,0.48)$ are $<3.7 \times 10^{-6}$ and $<6.1 \times 10^{-7}$, respectively (86ER05): see also (83VE10), and (85AJ01) for the earlier work. See also ${ }^{7} \mathrm{Li}$ in (88AJ01), (84AL1M, 84XI1A, 86CO1M, 86DR1G, 86GR1F, 86OL1B, 86WI1B; applied) and (86AB1E, 86MI1G, 88MA1H).
23. ${ }^{10} \mathrm{~B}\left(\mathrm{p}, \pi^{+}\right)^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-128.897$

Angular distributions have been obtained at $E_{\mathrm{p}}=168$ to 800 MeV to several states of ${ }^{11} \mathrm{~B}$ [see (80AJ01, 85AJ01)] as have cross sections for $\pi^{+}$production near threshold. At $E_{\overrightarrow{\mathrm{p}}}=200$ to 260 MeV , angular distributions and analyzing powers have been measured for the groups to ${ }^{11} \mathrm{~B}^{*}(0,2.12)$ (85ZI04).
24. ${ }^{10} \mathrm{~B}(\mathrm{~d}, \mathrm{p}){ }^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=9.2296
$$

Table 11.12: Beta decay of ${ }^{11} \mathrm{Be}(82 \mathrm{MI} 08)^{\text {a }}$

| ${ }^{11} \mathrm{~B}(\mathrm{keV})$ | $J^{\text {a b }}$ | Branching ${ }^{\mathrm{c}}$ <br> ratio (\%) | $\log f t$ | $E_{\gamma}(\mathrm{keV})$ | $\begin{aligned} & \hline I_{\gamma}{ }^{c} \\ & (\%) \end{aligned}$ | Transition to ${ }^{11} \mathrm{~B}^{*}(\mathrm{MeV})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| g.s. | $\frac{3}{2}^{-}$ | $54.7 \pm 2.0{ }^{\text {d }}$ | $6.830 \pm 0.016$ |  |  |  |
| $2124.693 \pm 0.027$ | $\frac{1}{2}{ }^{-}$ | $31.4 \pm 1.8$ | $6.648 \pm 0.025$ | $2124.473 \pm 0.027$ | 100 | g.s. |
| $4444.89 \pm 0.50$ | $\frac{5}{2}^{-}$ | $0.054 \pm 0.004$ | $10.93 \pm 0.03{ }^{\text {e }}$ | $4443.90 \pm 0.50$ | 100 | g.s. |
| $5020.31 \pm 0.30$ | $\frac{3}{2}^{-}$ | $0.282 \pm 0.020$ | $7.934 \pm 0.031$ | $5018.98 \pm 0.40$ | $85.6 \pm 0.6$ | g.s. |
|  |  |  |  | $2895.30 \pm 0.40$ | $14.4 \pm 0.6$ | 2.12 |
| $6791.80 \pm 0.30^{\text {f }}$ | $\frac{1}{2}^{+}$ | $6.47 \pm 0.45$ | $5.938 \pm 0.030$ | $6789.81 \pm 0.50$ | $67.5 \pm 1.1$ | g.s. |
|  |  |  |  | $4665.90 \pm 0.40$ | $28.5 \pm 1.1$ | 2.12 |
|  |  |  |  | $1171.31 \pm 0.30$ | $4.0 \pm 0.3$ | 5.02 |
| $7285.51 \pm 0.43$ | $\frac{5}{2}+$ | $<0.03$ | > 8.04 | 7282.92 | $87.0 \pm 2.0$ | g.s. |
| $7977.84 \pm 0.42{ }^{\text {g }}$ | $\frac{3}{2}+$ | $4.00 \pm 0.30$ | $5.576 \pm 0.033$ | 7974.73 | $46.2 \pm 1.1$ | g.s. |
|  |  |  |  | $5851.47 \pm 0.42$ | $53.2 \pm 1.2$ | 2.12 |
|  |  |  |  | $692.31 \pm 0.10$ | $0.85 \pm 0.04$ | 7.29 |
| 9.876 | $\frac{3}{2}+$ | $3.1 \pm 0.4{ }^{\text {h }}$ | $4.04 \pm 0.08$ |  |  |  |

${ }^{\text {a }}$ See also Tables 11.15 in (80AJ01) and 11.13 in (85AJ01).
${ }^{\mathrm{b}}$ From Table 11.3.
${ }^{\text {c }}$ Adopted by (82MI08); based on their work and on the earlier work.
${ }^{\mathrm{d}}$ From the relative intensities of the $\gamma$-rays and $I_{2.13} / I_{\text {total } \beta}=0.355 \pm 0.018$.
${ }^{\mathrm{e}} \log f_{1} t$.
${ }^{\mathrm{f}}$ Transition to ${ }^{11} \mathrm{~B}^{*}(4.44)$ is $<0.04 \%$.
${ }^{\mathrm{g}}$ Transitions to ${ }^{11} \mathrm{~B}^{*}(4.44,5.02,6.79)$ are $<0.06,<0.09$ and $<0.10 \%$.
${ }^{\mathrm{h}}$ From the relative intensities of the $\gamma$-rays and $I_{\alpha} / I_{2.12}$ of (81AL03).

Reported proton groups are displayed in Table 11.14 of (80AJ01). Angular distributions have been studied at many energies in the range $E_{\mathrm{d}}=0.17$ to 28 MeV [see (68AJ02, 75AJ02, 80AJ01)]. The lowest five levels are formed by $l_{\mathrm{n}}=1$ except for ${ }^{11} \mathrm{~B}^{*}(2.12)$ which appears to involve a spin-flip process. They are presumed to comprise the set $\frac{3}{2}^{-}, \frac{1^{-}}{2}, \frac{5^{-}}{2}, \frac{3}{2}^{-}, \frac{7}{2}^{-}$ expected as the lowest $\mathrm{p}^{7}$ levels $(a / K \approx 4.0) .{ }^{11} \mathrm{~B}^{*}(9.19,9.27)\left[J^{\pi}=\frac{7}{2}^{+}, \frac{5}{2}^{+}\right]$show strong $l=0$ stripping and are ascribed to capture of a 2 s neutron by ${ }^{10} \mathrm{~B}$ : see (68AJ02) for a listing of all the relevant references. Studies of p $\gamma$ correlations are discussed in reaction 14 of (68AJ02) and displayed in Table 11.4 of this paper. See also ${ }^{12} \mathrm{C}$.
25. (a) ${ }^{10} \mathrm{~B}(\mathrm{t}, \mathrm{d})^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=5.1969$
(b) ${ }^{10} \mathrm{~B}\left(\alpha,{ }^{3} \mathrm{He}\right)^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-9.1236$

See (68AJ02, 75AJ02).
26. (a) ${ }^{10} \mathrm{~B}\left({ }^{7} \mathrm{Li},{ }^{6} \mathrm{Li}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=4.204$
(b) ${ }^{10} \mathrm{~B}\left({ }^{9} \mathrm{Be},{ }^{8} \mathrm{Be}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=9.7888$
(c) ${ }^{10} \mathrm{~B}\left({ }^{13} \mathrm{C},{ }^{12} \mathrm{C}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=6.5078$

See (80AJ01, 85AJ01).
27. ${ }^{11} \mathrm{~B}\left(\beta^{-}\right){ }^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=11.506
$$

${ }^{11}$ Be decays to many states of ${ }^{11} \mathrm{~B}$ : see Table 11.12 for the observed $\beta$ - and $\gamma$-transitions (82MI08). ${ }^{11} \mathrm{~B}^{*}(9.88)$ decays via $\alpha$-emission for ${ }^{7} \mathrm{Li}^{*}(0,0.48)$ with branching ratios ( $87.4 \pm$ $1.2) \%$ and $(12.6 \pm 1.2) \%$, respectively (81AL03). A study of the $\beta \nu$ angular correlation in the first-forbidden decay of ${ }^{11} \mathrm{Be}$ to the $\frac{1}{2}^{-}$state ${ }^{11} \mathrm{~B}^{*}(2.12)$ has been performed: the $\beta$-transition is dominated by rank- 0 matrix elements and is of interest as a test of meson-exchange effects: see (85AJ01). See also (88WA1E).
28. (a) ${ }^{11} \mathrm{~B}(\gamma, \mathrm{n})^{10} \mathrm{~B}$
$Q_{\mathrm{m}}=-11.4542$
(b) ${ }^{11} \mathrm{~B}(\gamma, \mathrm{p})^{10} \mathrm{Be}$
$Q_{\mathrm{m}}=-11.2279$
(c) ${ }^{11} \mathrm{~B}(\gamma, \mathrm{~d}){ }^{9} \mathrm{Be}$
$Q_{\mathrm{m}}=-15.8153$
(d) ${ }^{11} \mathrm{~B}(\gamma, \mathrm{t}){ }^{8} \mathrm{Be}$
$Q_{\mathrm{m}}=-11.2234$

The giant dipole resonance is shown to consist mainly of $T=\frac{1}{2}$ states in the lower energy region and of $T=\frac{3}{2}$ states in the higher energy region by observing the decay to states in ${ }^{10} \mathrm{~B}$ and ${ }^{10} \mathrm{Be}$ [reactions (a) and (b)]. Absolute measurements of the ${ }^{11} \mathrm{~B}(\gamma$, all n$)$ cross section have been carried out from threshold to 35 MeV : the cross section exhibits a main peak at $E_{\gamma}=25$ to 28 MeV and weak shoulders at 13 and 16 MeV . The integrated cross section to 35 MeV is $69.1 \pm 0.8 \mathrm{MeV} \cdot \mathrm{mb}$ : see (80AJ01) and (88DI02). See also (84AL22). For other structures reported in the ( $\gamma, \mathrm{n}$ ) and ( $\gamma, \mathrm{p}$ ) cross sections see (75AJ02). The ( $\gamma$, $\mathrm{d}_{0}$ ) cross section peaks at $\approx 19 \mathrm{MeV}$, lower than it would if $T=\frac{3}{2}$ states were involved. The yield of $3.37 \mathrm{MeV} \gamma$-rays [from ${ }^{10} \mathrm{Be}^{*}(3.37)$, reaction (b)] has been measured for $E_{\mathrm{bs}}=100$ to 800 MeV . See also (84AL22, 86AL24). For reaction (d) see (86AL24). See (80AJ01, 85AJ01) for references and for other photonuclear processes. See also (85CH27, 85GO1A, 87KI1C, 87LU1B, 88DU04; theor.).
29. ${ }^{11} \mathrm{~B}(\gamma, \gamma){ }^{11} \mathrm{~B}$

Widths of excited states are displayed in Table 11.13. See also (84AL22, 88BEYY).
30. (a) ${ }^{11} \mathrm{~B}(e, e)^{11} \mathrm{~B}$
(b) ${ }^{11} \mathrm{~B}(\mathrm{e}, \text { ep })^{10} \mathrm{Be} \quad Q_{\mathrm{m}}=-11.2279$

$$
\left\langle r^{2}\right\rangle^{1 / 2}=2.43 \pm 0.11 \mathrm{fm} \quad \text { (86DO1E; prelim.). }
$$

[See also unpublished result in (80AJ01).]
Magnetic elastic scattering at $\theta=180^{\circ}$ shows strong M3 effects: the derived ratio of static M3/M1, $2.9 \pm 0.2 \mathrm{fm}^{2}$, suggests a $j$-j coupling scheme for ${ }^{11} \mathrm{~B}$ (g.s.). The quadrupole contribution to the elastic form factor is best accounted for by the undeformed shell model, $Q=3.72( \pm 20 \%) \mathrm{fm}^{2},\left\langle r^{2}\right\rangle^{1 / 2}=2.42 \mathrm{fm}$. See (80AJ01) for references. A recent study of the elastic scattering for $q=2.0$ to $3.9 \mathrm{fm}^{-1}$ is reported by ( 88 HI 02 ): the M3 component is dominant in the elastic form factor for $q>1.5 \mathrm{fm}^{-1}$.

The excitiation of ${ }^{11} \mathrm{~B}^{*}(2.1,4.4,5.0,8.6,8.9)$ has been studied. The giant resonance region, centered at $\approx 18 \mathrm{MeV}$, is characterized by a lack of prominent features except for a pronounced peak at $E_{\mathrm{x}}=13.0 \pm 0.1 \mathrm{MeV}$ (mixed M1-E2) and a broad transverse group at $E_{\mathrm{x}}=15.5 \mathrm{MeV}$. At $E_{\mathrm{e}}=121,186$ and 250 MeV form factors (and $B(\mathrm{E} \lambda) \uparrow$ ) are obtained for ${ }^{11} \mathrm{~B}^{*}(4.4,6.7,8.5,8.9,13.00 \pm 0.15)$ and the excitation of ${ }^{11} \mathrm{~B}^{*}(14.50 \pm 0.15,16.7 \pm 0.2)$ is also reported: see (85AJ01). See also (84DO1A, 87DE1A).

For $\Gamma_{\gamma_{0}}$ see Table 11.13. For reaction (b) see (75AJ02). See also (85KE1E, 86HA1M, 86KE1F, 87AL1M, 87DO12; theor.).

Table 11.13: Gamma widths from ${ }^{11} \mathrm{~B}(\gamma, \gamma){ }^{11} \mathrm{~B}$ and ${ }^{11} \mathrm{~B}(\mathrm{e}, \mathrm{e}){ }^{11} \mathrm{~B}{ }^{\text {a }}$

| $E_{\mathrm{x}}(\mathrm{MeV})$ | $J^{\pi}$ | $\Gamma_{\gamma_{0}}(\mathrm{eV})$ | Reaction |
| :---: | :---: | :---: | :---: |
| 2.12 | $\frac{1}{2}^{-}$ | $0.120 \pm 0.009^{\mathrm{b}}$ | $(\gamma, \gamma)$ |
| 4.44 | $\frac{5}{2}^{-}$ | $0.58 \pm 0.04$ | $(\gamma, \gamma)$ |
|  |  | $0.55 \pm 0.02$ | $(\gamma, \gamma)$ |
|  |  | $0.60 \pm 0.09(\mathrm{M} 1)$ | $(\mathrm{e}, \mathrm{e})$ |
|  |  | $\pm 0.016 \pm 0.002(\mathrm{E} 2)$ |  |
|  |  | $0.56 \pm 0.02^{\mathrm{b}}$ |  |
| 5.02 | $\frac{3}{2}^{-}$ | $1.80 \pm 0.13$ | $(\gamma, \gamma)$ |
|  |  | $1.64 \pm 0.07$ | $(\gamma, \gamma)$ |
|  |  | $1.73 \pm 0.14(\mathrm{M} 1)$ | $(\mathrm{e}, \mathrm{e})$ |
|  |  | $\leq 0.0034(\mathrm{E} 2)$ |  |
|  |  | $1.68 \pm 0.06^{\mathrm{b}}$ |  |
| 6.74 | $\frac{7}{2}^{-}$ | $0.021 \pm 0.005$ | $(\gamma, \gamma)$ |
| 7.79 | $\frac{1}{2}^{+}$ | $0.26 \pm 0.03$ | $(\gamma, \gamma)$ |
| 7.29 | $\frac{5}{2}^{+}$ | $1.00 \pm 0.07^{\mathrm{b}}$ | $(\gamma, \gamma)$ |
| 7.98 | $\frac{3}{2}^{+}$ | $0.53 \pm 0.07$ | $(\gamma, \gamma)$ |
| 8.56 | $\left.\frac{3}{2}^{-}\right)$ | $0.53 \pm 0.05$ | $(\gamma, \gamma)$ |
| 8.92 | $\frac{5}{2}^{-}$ | $4.15 \pm 0.20^{\mathrm{b}}$ | $(\gamma, \gamma) ;(\mathrm{e}, \mathrm{e})$ |

[^2]31. ${ }^{11} \mathrm{~B}\left(\pi^{+}, \pi^{+}\right)^{11} \mathrm{~B}$

The proton matter distribution in ${ }^{11} \mathrm{~B}_{\text {g.s. }}$ has a radius of $2.368 \pm 0.021 \mathrm{fm}$, assuming that for ${ }^{12} \mathrm{C}$ to be 2.44 fm . The result is not sensitive to the details of the optical-model calculations ( $80 \mathrm{BA} 45 ; E_{\pi^{+}}=38.6$ and 47.7 MeV ). See also the "General" section.
32. ${ }^{11} \mathrm{~B}(\mathrm{n}, \mathrm{n}){ }^{10} \mathrm{~B}$

Angular distributions have been reported for $E_{\mathrm{n}}=75 \mathrm{keV}$ to 14.1 MeV [see (80AJ01, 85AJ01)] and at $E_{\mathrm{n}}=8.0$ to $13.9 \mathrm{MeV}\left(82 \mathrm{GL} 02 ; \mathrm{n}_{0} \rightarrow \mathrm{n}_{3}\right)$. Recent work (prelim.) is reported to 17 MeV (86MU1D; $\mathrm{n}_{0}$ ). See also ${ }^{12} \mathrm{~B}$, (85WA1P) and (88HAZT; theor.).
33. (a) ${ }^{11} \mathrm{~B}(\mathrm{p}, \mathrm{p})^{11} \mathrm{~B}$
$\begin{array}{ll}\text { (b) }{ }^{11} \mathrm{~B}(\mathrm{p}, 2 \mathrm{p})^{10} \mathrm{Be} & Q_{\mathrm{m}}=-11.2279 \\ (\mathrm{c}){ }^{11} \mathrm{~B}(\mathrm{p}, \mathrm{pn})^{10} \mathrm{~B} & Q_{\mathrm{m}}=-11.4542\end{array}$
Observed proton groups are displayed in Table 11.4. Angular distributions have been measured for $E_{\mathrm{p}}=6$ to 185 MeV [see (80AJ01)] and at 1 GeV (85AL1F). For reactions (b) and (c) at 1 GeV see (85BE1J, 85DO1B). For pion production see (87AB1E). See also ${ }^{12} \mathrm{C}$, (88BE2B), (85MUZZ) and (85AJ01).
34. ${ }^{11} \mathrm{~B}(\mathrm{~d}, \mathrm{~d}){ }^{11} \mathrm{~B}$

Elastic scattering has been studied at $E_{\mathrm{d}}=5.5$ and 11.8 MeV : see (80AJ01).
35. ${ }^{11} \mathrm{~B}(\mathrm{t}, \mathrm{t})^{11} \mathrm{~B}$

The elastic scattering has been studied at $E_{\mathrm{t}}=1.8$ and 2.1 MeV : see (80AJ01).
36. ${ }^{11} \mathrm{~B}\left({ }^{3} \mathrm{He},{ }^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}$

The elastic scattering has been studied at $E\left({ }^{3} \mathrm{He}\right)=8$ to 74 MeV : see (75AJ02, 80AJ01). At $E\left({ }^{3} \mathrm{He}\right)=17.5$ and 40 MeV angular distributions have also been studied for the ${ }^{3} \mathrm{He}$ ions to ${ }^{11} \mathrm{~B}^{*}(2.12,4.44,5.02,6.74) . T=\frac{3}{2}$ states observed in this reaction are displayed in Table 11.14. See also (85AJ01). There is a weak indication of a state at $E_{\mathrm{x}}=14.51 \mathrm{MeV}$ : see (75AJ02). See also (86JA14) and (87TR01; theor.).

Table 11.14: $T=\frac{3}{2}$ states in ${ }^{11} \mathrm{~B}{ }^{\text {a }}$

| Reaction | $E_{\mathrm{x}}(\mathrm{MeV} \pm \mathrm{keV})$ | $\Gamma_{\text {c.m. }}(\mathrm{keV})$ |
| :---: | :---: | :---: |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ | $12.563 \pm 20$ | $202 \pm 25$ |
| ${ }^{10} \mathrm{Be}(\mathrm{p}, \gamma){ }^{11} \mathrm{~B}$ | $12.56 \pm 30$ | $230 \pm 65$ |
| ${ }^{11} \mathrm{~B}\left({ }^{3} \mathrm{He},{ }^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}^{*}$ | $\underline{12.51 \pm 50}$ | $\underline{260 \pm 50}$ |
|  | $12.557 \pm 16^{\text {b }}$ | $215 \pm 21^{\text {b }}$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ | $12.920 \pm 20$ | $155 \pm 25$ |
| ${ }^{10} \mathrm{Be}(\mathrm{p}, \gamma){ }^{11} \mathrm{~B}$ | $12.91 \pm 20$ | $235 \pm 27$ |
| ${ }^{13} \mathrm{C}\left(\mathrm{p},{ }^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}$ | $12.94 \pm 50$ | $350 \pm 50$ |
| ${ }^{13} \mathrm{C}\left(\mathrm{p},{ }^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}$ | $12.91 \pm 30$ | $260 \pm 50$ |
| ${ }^{14} \mathrm{C}(\mathrm{p}, \alpha){ }^{11} \mathrm{~B}$ | $\underline{12.92 \pm 20}{ }^{\text {c }}$ | $\underline{238 \pm 15}$ |
|  | $12.916 \pm 12^{\text {d }}$ | $155 \pm 25^{\text {d }}$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ | $14.40{ }^{\text {e }}$ | $261 \pm 25$ |
| ${ }^{10} \mathrm{Be}(\mathrm{p}, \gamma){ }^{11} \mathrm{~B}$ | $14.33 \pm 20$ | $255 \pm 30$ |
| ${ }^{11} \mathrm{~B}\left({ }^{3} \mathrm{He},{ }^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}^{*}$ | $\underline{14.40 \pm 50}$ | $\underline{220 \pm 50}$ |
|  | $14.34 \pm 20^{\text {b }}$ | $254 \pm 18^{\text {b }}$ |
| ${ }^{10} \mathrm{Be}(\mathrm{p}, \gamma){ }^{11} \mathrm{~B}$ | $15.32 \pm 100^{\text {c }}$ | $635 \pm 180$ |
| ${ }^{14} \mathrm{C}(\mathrm{p}, \alpha)^{11} \mathrm{~B}$ | $15.29 \pm 25^{\text {c }}$ | $282 \pm 15$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ | $16.437 \pm 2{ }^{\text {f }}$ | $\leq 30$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ | 17.69 | $91 \pm 25$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ | $18.0 \pm 100$ | $870 \pm 100$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ | $19.146 \pm 30^{\text {f }}$ | $115 \pm 25$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{p}\right){ }^{11} \mathrm{~B}$ | $21.27 \pm 50$ | $300 \pm 30$ |

${ }^{\text {a }}$ See also Table 11.18 in (80AJ01). See Table 11.16 in (85AJ01) for references.
${ }^{\mathrm{b}}$ Mean value.
c See Table 11.3
d "Best" value.
${ }^{e}$ May have mixed isospin $\left(T=\frac{1}{2}+T=\frac{3}{2}\right)$.
${ }^{\mathrm{f}}$ See also reaction 60 (85AR03).
$37 .{ }^{11} \mathrm{~B}(\alpha, \alpha){ }^{11} \mathrm{~B}$

Angular distributions have been reported at $E_{\alpha}=24$ to 31.2 MeV : [see (75AJ02, 80AJ01, 85AJ01)] and at 48.7 and 54.1 MeV ( $87 \mathrm{AB} 03 ; \alpha_{0}$ ). See also (83SA07) and (85SH1D; theor.).
38. (a) ${ }^{11} \mathrm{~B}\left({ }^{6} \mathrm{Li},{ }^{6} \mathrm{Li}\right)^{11} \mathrm{~B}$
(b) ${ }^{11} \mathrm{~B}\left({ }^{7} \mathrm{Li},{ }^{7} \mathrm{Li}\right){ }^{11} \mathrm{~B}$

The elastic scattering has been studied at $E\left({ }^{6} \mathrm{Li}\right)=28 \mathrm{MeV}$ : see (75AJ02). At $E\left({ }^{7} \mathrm{Li}\right)=$ 34 MeV angular distributions have been reported to ${ }^{11} \mathrm{~B}^{*}(0,2.12,4.44,5.02,6.74,7.29,8.92)$ (87CO02, 87CO16). See also (88HN01; theor.).
39. (a) ${ }^{11} \mathrm{~B}\left({ }^{9} \mathrm{Be},{ }^{9} \mathrm{Be}\right){ }^{11} \mathrm{~B}$
(b) ${ }^{11} \mathrm{~B}\left({ }^{10} \mathrm{~B},{ }^{10} \mathrm{~B}\right){ }^{11} \mathrm{~B}$
(c) ${ }^{11} \mathrm{~B}\left({ }^{11} \mathrm{~B},{ }^{11} \mathrm{~B}\right){ }^{11} \mathrm{~B}$

For reaction (a) see (84DA17, 86CU02). For fusion cross sections (reactions (b) and (c)) see (89SZ01). See also (75AJ02, 80AJ01), (85BE1A, 85CU1A) and (84HA43, 86RO12; theor.).
40. (a) ${ }^{11} \mathrm{~B}\left({ }^{12} \mathrm{C},{ }^{12} \mathrm{C}\right){ }^{11} \mathrm{~B}$
(b) ${ }^{11} \mathrm{~B}\left({ }^{13} \mathrm{C},{ }^{13} \mathrm{C}\right){ }^{11} \mathrm{~B}$

The elastic scattering has been studied at $E\left({ }^{11} \mathrm{~B}\right)=18.8$ to 50 MeV and at $E\left({ }^{12} \mathrm{C}\right)=15$ to 24 MeV and 87 MeV [see (80AJ01, 85AJ01)] as well as at $E\left({ }^{11} \mathrm{~B}\right)=10.4,12.4$ and 14.6 MeV (JA85), at $E_{\text {c.m. }}=25 \mathrm{MeV}(86 \mathrm{MA13})$, at $E\left({ }^{11} \mathrm{~B}\right)=42.5$ to 100 MeV (85MA10) and at $E\left({ }^{12} \mathrm{C}\right)=65 \mathrm{MeV}\left(85 \mathrm{GO} 1 \mathrm{H}\right.$; prelim.; involving various states of $\left.{ }^{12} \mathrm{C}\right)$ [see $\left.{ }^{12} \mathrm{C}\right]$. The population of ${ }^{11} \mathrm{~B}^{*}(2.12,4.44,6.79)$ is also reported. For yields, fusion and breakup studies see (85AJ01) and (85MA10, 86MA13). For reaction (b) see (84DE1J, 84HAZK; prelim.). See also (87PO15 ), (84FR1A, 84HA53, 85BE1A, 85CU1A, 88MA07), (82BA1D, 85BA1T; astrophys.) and (84HA43, 84IN03, 85KO1J, 86BA69, 86HA13; theor.).
41. ${ }^{11} \mathrm{~B}\left({ }^{14} \mathrm{~N},{ }^{14} \mathrm{~N}\right){ }^{11} \mathrm{~B}$

The elastic scattering has been investigated at $E\left({ }^{14} \mathrm{~N}\right)=41,77$ and 133 MeV : see (75AJ02, 85AJ01). See also (85BE1A, 85CU1A) and (84HA43; thoer.).
42. (a) ${ }^{11} \mathrm{~B}\left({ }^{16} \mathrm{O},{ }^{16} \mathrm{O}\right)^{11} \mathrm{~B}$
(b) ${ }^{11} \mathrm{~B}\left({ }^{18} \mathrm{O},{ }^{18} \mathrm{O}\right){ }^{11} \mathrm{~B}$

The elastic scattering in reaction (a) has been studied at $E\left({ }^{16} \mathrm{O}\right)=14.5$ to 60 MeV and at $E\left({ }^{11} \mathrm{~B}\right)=41.6,49.5$ and 115 MeV . The elastic scattering in reaction $(\mathrm{b})$ is reported at $E\left({ }^{11} \mathrm{~B}\right)=115 \mathrm{MeV}$. For references see (75AJ02, 80AJ01, 85AJ01).
43. ${ }^{11} \mathrm{~B}\left({ }^{20} \mathrm{Ne},{ }^{20} \mathrm{Ne}\right){ }^{11} \mathrm{~B}$

The elastic angular distribution has been studied at $E\left({ }^{11} \mathrm{~B}\right)=115 \mathrm{MeV}$ : see (85AJ01).
44. (a) ${ }^{11} \mathrm{~B}\left({ }^{24} \mathrm{Mg},{ }^{24} \mathrm{Mg}\right)^{11} \mathrm{~B}$
(b) ${ }^{11} \mathrm{~B}\left({ }^{25} \mathrm{Mg},{ }^{25} \mathrm{Mg}\right){ }^{11} \mathrm{~B}$
(c) ${ }^{11} \mathrm{~B}\left({ }^{26} \mathrm{Mg},{ }^{26} \mathrm{Mg}\right){ }^{11} \mathrm{~B}$
(d) ${ }^{11} \mathrm{~B}\left({ }^{27} \mathrm{Al},{ }^{27} \mathrm{Al}\right){ }^{11} \mathrm{~B}$
(e) ${ }^{11} \mathrm{~B}\left({ }^{28} \mathrm{Si},{ }^{28} \mathrm{Si}\right){ }^{11} \mathrm{~B}$

The elastic angular distributions for reactions (a) to (d) have been studied at $E\left({ }^{11} \mathrm{~B}\right)=$ 79.6 MeV: see (85AJ01). See also (87PO15 ). For reaction (e) see (84TE1A).
45. (a) ${ }^{11} \mathrm{~B}\left({ }^{40} \mathrm{Ar},{ }^{40} \mathrm{Ar}\right){ }^{11} \mathrm{~B}$
(b) ${ }^{11} \mathrm{~B}\left({ }^{40} \mathrm{Ca},{ }^{40} \mathrm{Ca}\right){ }^{11} \mathrm{~B}$

For reaction (a) see ( 85 MO 1 K ; prelim.). Angular distributions have been reported in reaction $(\mathrm{b})$ at $E\left({ }^{11} \mathrm{~B}\right)=51.5 \mathrm{MeV}$ to ${ }^{11} \mathrm{~B}^{*}(0,2.12)$ : see (85AJ01).
46. ${ }^{11} \mathrm{C}\left(\beta^{+}\right)^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=1.982$

See ${ }^{11} \mathrm{C}$.
47. (a) ${ }^{12} \mathrm{C}(\gamma, \mathrm{p})^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-15.9572$
(b) ${ }^{12} \mathrm{C}(e, p){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-15.9572$

The fraction of transitions to the ground and to excited states of ${ }^{11} \mathrm{~B}$ (and to ${ }^{11} \mathrm{C}$ states reached in the $(\gamma, \mathrm{n})$ reaction) has been measured at $E_{\mathrm{b} . \mathrm{s} .}=21.7$ to 42 MeV : the ground state is predominantly populated: see (80AJ01). The predominant population of ${ }^{11} \mathrm{~B}_{\mathrm{g} . \mathrm{s} .}$ has also recently been observed at $E_{\gamma}=28 \mathrm{MeV}$ (89FE01). Analog states are populated similarly in the $(\gamma, \mathrm{n})$ and $(\gamma, \mathrm{p})$ reactions. Angular distributions for the protons to several states of ${ }^{11} \mathrm{~B}$ have been measured at $E_{\gamma}=21.7 \rightarrow 31 \mathrm{MeV}$ and at 60,80 and 100 MeV [see (80AJ01, 85AJ01)] as well as in the giant resonance region [see ${ }^{12} \mathrm{C}$ ] (86KE06; $\mathrm{p}_{0}$ ) and at 60 MeV (88SH08; p to ${ }^{11} \mathrm{~B}^{*}\left(0,2.12,5.0,6.8\right.$ (unres.)). The relative population of ${ }^{11} \mathrm{~B}^{*}(6.8)$ is much greater than that reported in (e, ep) (88SH08). Spectra have also been studied by (86AN25, 86MC15). For reaction (b) see (85AJ01). See also ${ }^{12} \mathrm{C}$, (87VO08) and (84BO18, 87GO37, 88OR02, 89PIZZ; theor.).
48. ${ }^{12} \mathrm{C}(e, e p){ }^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=-15.9572
$$

(88VA09) have studied the $l=1$ knockout to ${ }^{11} \mathrm{~B}^{*}(0,2.12,5.02)$ at $E_{\mathrm{e}}=284.5$ to 481.1 MeV. One-third to one-half of the sum-rule strength predicted by the independentparticle shell model is observed. See (88VA09) also for a review of spectroscopic factors. ${ }^{11} B^{*}(4.44)$ is not observed: the two-step processes which are necessary to excite it in this reaction appear to be weak (85VA16, 88VA21). Weak transitions have been studied to states at $E_{\mathrm{x}}=6.751$ (unresolved), $7.278,7.954,8.61,9.820( \pm 25 \mathrm{keV}$, except $\pm 50 \mathrm{keV}$ for 8.61) and to a broad structure at $11.5 \mathrm{MeV} . l=0$ and 1 are suggested for the structures at 9.8 and 11.5 MeV (88VA21; also $S_{\alpha}$ ). See also the earlier work in (85VA05). The effects of the nuclear medium have been studied by (86VA17, 87UL03, 88VA09): see ${ }^{12} \mathrm{C}$. See ${ }^{12} \mathrm{C}$ and (84CA34, 87CAZY) for the decay of ${ }^{12} \mathrm{C}$ states to ${ }^{11} \mathrm{~B}^{*}(0,2.12)$. See also (85DE56, 86DE1U, 86LA1T, 88HA12, 89BOZZ) and (84LA16, 85CA32, 85LA1F, 86DE05, 87BL10, 87GOZ0, 87VA15, 88HO10, 88SU02, 89RY03; theor.).
49. ${ }^{12} \mathrm{C}\left(\pi^{+}, \pi^{+} \mathrm{p}\right)^{11} \mathrm{~B} \quad Q_{\mathrm{m}}=-15.9572$

At $E_{\pi^{+}}=100$ to 200 MeV the reaction proceeds primarily to ${ }^{11} \mathrm{~B}_{\text {g.s. }}$. At $E_{\pi}=200 \mathrm{MeV}$ the ratios for $\sigma_{n} / \sigma_{p}$ for the first excited states in ${ }^{11} \mathrm{C} /{ }^{11} \mathrm{~B}$ are $1.4 \pm 0.2$ for $\pi^{-}$and $1 / 1.8 \pm 0.2$ for $\pi^{+}$. At $E_{\pi^{+}}=60$ to $300 \mathrm{MeV}^{11} \mathrm{~B}^{*}(4.44)\left[J^{\pi}=\frac{5}{2}^{-}\right]$is strongly populated as is the analog state in the mirror reaction: see (80AJ01, 85AJ01) for references. At $E_{\pi^{ \pm}}=220 \mathrm{MeV}$ the quasi-elastic nature of the scattering has been studied by (84FA11). See also the studies by (84ZI1B, 87HU02), ${ }^{12} \mathrm{C}$, (84GO1F), (86CH1J) and (85CO03; theor.).
50. ${ }^{12} \mathrm{C}(\mathrm{n}, \mathrm{d}){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-13.7326$

See (85FR07, 87FR16, 89ROZW) and in ${ }^{13} \mathrm{C}$ in (86AJ01, 90AJ01). See also (86DO12, 88YOZX).
51. ${ }^{12} \mathrm{C}(\mathrm{p}, 2 \mathrm{p})^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-15.9572$

At $E_{\mathrm{p}}=98.7 \mathrm{MeV}$ groups are observed to ${ }^{11} \mathrm{~B}^{*}(0,2.12,4.44,5.02,6.79)$. DWIA lead to relative spectroscopic factors of $2.0,0.37,0.15,1.08,0.25$ for these states. No evidence is seen for multistep reaction processes which would be necessary to populate ${ }^{11} \mathrm{~B}^{*}(4.44,6.74)$ : see (85AJ01). At $E_{\mathrm{p}}=1 \mathrm{GeV}$ the separation energy between 6 and 14 MeV broad $1 \mathrm{p}_{3 / 2}$ and $1 \mathrm{~s}_{1 / 2}$ groups is 18 MeV (85BE1J, 85DO1B). See also (84VD1B, 86VD1C; $E_{\mathrm{p}}=50 \mathrm{MeV}$ ), (89TEZZ) and (85DE56, 87VD1A).
52. ${ }^{12} \mathrm{C}\left(\mathrm{d},{ }^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-10.4637$

Angular distributions of ${ }^{3} \mathrm{He}$ ions have been measured for $E_{\mathrm{d}}=20$ to 80 MeV and spectroscopic factors have been derived for ${ }^{11} \mathrm{~B}^{*}(0,2.12,5.02)$ : see (75AJ02, 80AJ01, 85AJ01).
53. ${ }^{12} \mathrm{C}(\mathrm{t}, \alpha)^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=3.8568
$$

Angular distributions have been measured at $E_{\mathrm{t}}=33$ and 38 MeV to ${ }^{11} \mathrm{~B}^{*}(0,2.12,4.44$, $5.02,6.74,7.29,7.98,8.56)$. As expected, the $\frac{5^{-}}{2}$ and $\frac{7}{2}^{-}$states ${ }^{11} \mathrm{~B}^{*}(4.44,6.74)$ are populated by two-step processes. The best $J^{\pi}$ value for ${ }^{11} \mathrm{~B}^{*}(8.56)$ is $\frac{3^{-}}{2}$ but this assumes some direct population which may not be the case (87FO21, 88SI08) [see for spectroscopic factors]. For the earlier work see (75AJ02).
54. ${ }^{12} \mathrm{C}\left(\alpha,{ }^{5} \mathrm{Li}\right){ }^{11} \mathrm{~B} \quad Q_{\mathrm{m}}=-17.92$

See (87GA20) and (85AJ01).
55. ${ }^{12} \mathrm{C}\left({ }^{6} \mathrm{Li},{ }^{7} \mathrm{Be}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-10.351$

At $E\left({ }^{6} \mathrm{Li}\right)=93 \mathrm{MeV},{ }^{11} \mathrm{~B}^{*}(0,2.12,5.0,6.8,8.9)$ are populated (88BUZI; prelim.). See also (86GL1E; prelim.).
56. ${ }^{12} \mathrm{C}\left({ }^{12} \mathrm{C},{ }^{13} \mathrm{~N}\right){ }^{11} \mathrm{~B} \quad Q_{\mathrm{m}}=-14.0134$

Angular distributions involving ${ }^{11} \mathrm{Bg}$.s. have been measured at $E\left({ }^{12} \mathrm{C}\right)=93.8$ and 114 MeV : see (85AJ01). See also (87WIZW).
57. ${ }^{12} \mathrm{C}\left({ }^{13} \mathrm{C},{ }^{14} \mathrm{~N}\right){ }^{11} \mathrm{~B} \quad Q_{\mathrm{m}}=-8.4066$

See (87AD07, 88VO08) and ${ }^{14} \mathrm{~N}$ in (90AJ01). See also (89VO1D).
58. ${ }^{12} \mathrm{C}\left({ }^{19} \mathrm{~F},{ }^{20} \mathrm{Ne}\right){ }^{11} \mathrm{~B} \quad Q_{\mathrm{m}}=-3.108$

At $E\left({ }^{19} \mathrm{~F}\right)=40,60$ and 68.8 MeV angular distributions involving ${ }^{11} \mathrm{~B}^{*}(0,2.12)$ and ${ }^{20} \mathrm{Ne}^{*}(0,1.63)$ have been measured: see (80AJ01). See also (86HE1A, 88DI08; theor.).
59. ${ }^{13} \mathrm{C}\left(\mathrm{p},{ }^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-13.1855$

At $E_{\mathrm{p}}=50.5 \mathrm{MeV}$, in addition to ${ }^{11} \mathrm{~B}^{*}(0,2.12,4.44,5.02,6.74,8.92)$, a state is observed at $E_{\mathrm{x}}=12.94 \pm 0.05 \mathrm{MeV}, \Gamma=350 \pm 50 \mathrm{keV}$. Comparison of the angular distributions of the ${ }^{3} \mathrm{He}$ and of the tritons [to the analog state] at $E_{\mathrm{p}}=43.7$ and 50.5 MeV lead to the assignments $J^{\pi}=\frac{1}{2}^{-}, T=\frac{3}{2}$ for this state and for ${ }^{11} \mathrm{C}^{*}(12.50)$ : the strong proton and the weak $\alpha$ decay are consistent with this assignment: see Table 11.14. Angular distributions have been measured at $E_{\mathrm{p}}=26.9$ to 49.6 MeV involving the above states except for ${ }^{11} \mathrm{~B}^{*}(8.92)$ and at $E_{\overrightarrow{\mathrm{p}}}=65 \mathrm{MeV}$ (to ${ }^{11} \mathrm{~B}^{*}(0,2.12)$ ): see (75AJ02, 80AJ01, 85AJ01). See also ${ }^{14} \mathrm{~N}$ in (86AJ01) and (85HA1J).
60. ${ }^{13} \mathrm{C}(\mathrm{d}, \alpha){ }^{11} \mathrm{~B} \quad Q_{\mathrm{m}}=5.1677$

Observed proton groups are displayed in Table 11.15. Angular distributions are reported at $E_{\mathrm{d}}=0.41$ to 14.1 MeV : See (75AJ02). See also (85HA1J).
61. ${ }^{14} \mathrm{C}(\mathrm{p}, \alpha){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-0.7842$

Table 11.15: States of ${ }^{11} \mathrm{~B}$ from ${ }^{11} \mathrm{~B}\left(\mathrm{p}, \mathrm{p}^{\prime}\right){ }^{11} \mathrm{~B}^{*},{ }^{13} \mathrm{C}(\mathrm{d}, \alpha){ }^{11} \mathrm{~B}$ and ${ }^{14} \mathrm{C}(\mathrm{p}, \alpha){ }^{11} \mathrm{~B}^{\mathrm{a}}$

| $E_{\mathrm{x}}(\mathrm{keV})^{\mathrm{b}}$ | $E_{\mathrm{x}}(\mathrm{keV})^{\mathrm{c}}$ | $E_{\mathrm{x}}(\mathrm{keV})^{\mathrm{d}}$ | $\Gamma_{\text {c.m. }}(\mathrm{keV})^{\mathrm{d}}$ |
| :---: | :---: | :---: | :---: |
| 0 | 0 | 0 |  |
| $2124.7 \pm 0.5$ | $2125.4 \pm 1.4$ | $2120 \pm 10$ |  |
| $4445.2 \pm 0.5$ | $4444.5 \pm 1.6$ | $4450 \pm 10$ |  |
| $5021.1 \pm 0.6$ | $5020.2 \pm 1.9$ | $5025 \pm 8$ |  |
| $6743.0 \pm 0 . \mathrm{e}^{\mathrm{e}}$ | $6745.8 \pm 3.4$ | $6746 \pm 5^{\mathrm{f}}$ |  |
| $6792.6 \pm 1.6$ | $6795 \pm 3.0$ |  |  |
| $7285.6 \pm 1.5$ |  |  |  |
| $7978.0 \pm 1.7$ |  | $8560 \pm 10^{\mathrm{g}}$ |  |
| $8559.4 \pm 1.9$ | $8520 \pm 70$ | $8920 \pm 10^{\mathrm{h}}$ |  |
| $8920.2 \pm 2.0$ | $8910 \pm 60$ |  |  |
| $9185.0 \pm 2.0$ |  | $10300 \pm 60^{\mathrm{i}}$ | $133 \pm 10$ |
| $9274.4 \pm 2.0$ |  | $11620 \pm 30$ | $186 \pm 25$ |
| $10450 \pm 150$ |  | $12920 \pm 20$ | $238 \pm 15$ |
| $11650 \pm 150$ |  | $14560 \pm 15$ | $42 \pm 27$ |
| $12850 \pm 100$ |  | $15290 \pm 25$ | $282 \pm 15$ |
|  |  | $16500 \pm 50$ | $201 \pm 10$ |
| $15200 \pm 150$ |  | $19070 \pm 50$ | $294 \pm 10$ |
| $16400 \pm 150$ |  |  |  |
|  |  |  |  |

${ }^{\text {a }}$ For references see Table 11.17 in (80AJ01).
${ }^{b}{ }^{11} \mathrm{~B}\left(\mathrm{p}, \mathrm{p}^{\prime}\right)^{11} \mathrm{~B}$.
c ${ }^{13} \mathrm{C}(\mathrm{d}, \alpha)^{11} \mathrm{~B}$.

${ }^{e}$ Values below are normalized to $E_{\mathrm{x}}=4445.3,5020.0$ and 6743.4 keV .
${ }^{f}$ Very strongly excited.
${ }^{8}$ Very weakly excited.
${ }^{\mathrm{h}}$ On the basis of the similarity with the angular distribution to ${ }^{11} \mathrm{~B}^{*}(4.44), J^{\pi}=\frac{5}{2}^{-}$is assigned.
${ }^{\mathrm{i}}$ This state and the ones below may be unresolved.

Observed states are displayed in Table 11.14 (85AR03). It is suggested ${ }^{11} \mathrm{~B}^{*}(12.92,15.29$, $16.50,19.07$ ) are $T=\frac{3}{2}$, negative-parity states. Spectroscopic factors have also been derived (85AR03).
62. (a) ${ }^{14} \mathrm{~N}(\mathrm{n}, \alpha)^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-0.1583$
(b) ${ }^{14} \mathrm{~N}(\mathrm{n}, 2 \alpha)^{7} \mathrm{Li}$
$Q_{\mathrm{m}}=-8.8220$

Angular distributions have been measured for $E_{\mathrm{n}}=4.9$ to 18.8 MeV [see (75AJ02, 80AJ01, 85AJ01)] and at $12.2,14.1$ and 18.0 MeV (86RU1B; $\alpha_{0}, \alpha_{1}$ ). At $E_{\mathrm{n}}=14.1$ and 15.7 MeV various states of ${ }^{11} \mathrm{~B}$ with $8.9<E_{\mathrm{x}}<14.5 \mathrm{MeV}$ appear to be involved in the sequential decay to ${ }^{7} \mathrm{Li}$. Angular correlation results are consistent with $J=\frac{7}{2}$ and $\frac{5}{2}$ for ${ }^{11} \mathrm{~B}^{*}(9.19,9.27)$ respectively: see (75AJ02). See also (85HA1J).
63. ${ }^{14} \mathrm{~N}\left(\mathrm{p}, \mathrm{p}^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=-20.7361$

See (86VD1C; prelim.; 50 MeV ).
64. ${ }^{16} \mathrm{O}\left(\mathrm{d},{ }^{7} \mathrm{Be}\right){ }^{11} \mathrm{~B}$

$$
Q_{\mathrm{m}}=-16.038
$$

At $E_{\mathrm{d}}=80 \mathrm{MeV}$ angular distributions have been measured to ${ }^{11} \mathrm{~B}^{*}(0,2.12,4.44+5.02$, $6.74+6.79+7.29):$ see (80AJ01).

## ${ }^{11} \mathrm{C}$

(Figs. 3 and 4)
GENERAL: (See also (85AJ01).)
Model calculations: (88WO04)
Special states: (85SH24, 86AN07, 88KW1A)
Astrophysical Questions: (87RA1D)
Complex reactions involving ${ }^{11} \mathrm{C}$ : (81AS04, 85AR09, 85HI1C, 85MO08, 86AV1B, 86AV07, 86BA3G, 86HA1B, 86HI1D, 86UT01, 87AR19, 87BA38, 87DE37, 87NA01, 87RI03, 87SN01, 87ST01, 87YA16, 88CA06, 88KI05, 88KI06, 88SA19, 88SM07, 88VUZZ, 89AR1G, 89HA1L, 89SA10, 89SE03, 89YO02) Applications: (85TA1D, 86WE1E, 87BO16, 87HI1B, 88FA1C, 88HI1F, 88VO1D, 89TR1B, 89WO1B)

Pion and kaon capture and reactions (see also reactions 19, 20 and 27): (84OH04, 88AB05, 88GIZU)

Hypernuclei: (AS84D, ZH84B, GA85A, DA86, DA86A)
Other topics: (85AN28, 85SH24, 85TA26, 86HE01, 88KW1A)
Ground-state properties of ${ }^{11} \mathrm{C}$ : (84ZI04, 85AN28, 85HA18, 85FA01, 85ZI05, 86GL1A, 87FU06, 87SA15, 88VA03, 88WA08, 88WO04, 89SA10)

$$
\begin{gathered}
\mu=-0.964 \pm 0.001 \mathrm{~nm} \text { (69WO03) } \\
Q=34.26 \mathrm{mb}(78 \mathrm{LEZA})
\end{gathered}
$$

1. ${ }^{11} \mathrm{C}\left(\beta^{+}\right){ }^{11} \mathrm{~B}$
$Q_{\mathrm{m}}=1.982$

The half life of ${ }^{11} \mathrm{C}$ is $1223.1 \pm 1.2 \mathrm{~s}$. Log $f t=3.599 \pm 0.002$. The ratio of $K$-capture to positron emission is $\left(0.230_{-0.011}^{+0.014}\right) \%$. See (80AJ01) for references. See also (85AJ01) and (87BO1Y).

$$
\text { 2. }{ }^{6} \mathrm{Li}\left({ }^{6} \mathrm{Li}, \mathrm{n}\right){ }^{11} \mathrm{C} \quad Q_{\mathrm{m}}=9.450
$$

At $\mathrm{E}\left({ }^{6} \mathrm{Li}\right)=4.1 \mathrm{MeV}$ angular distributions have been obtained for the neutrons to ${ }^{11} \mathrm{C}^{*}(2.00,4.32,4.80,6.34+6.48,6.90,7.50)$. In addition, $\mathrm{n} \gamma$ - coincidences via ${ }^{11} \mathrm{C}^{*}(8.42)$ [and an $8.42 \mathrm{MeV} \gamma$-ray] are reported. ${ }^{11} \mathrm{C}^{*}(8.10)$ was not observed. The lifetimes, $\tau_{\mathrm{m}}$, for ${ }^{11} \mathrm{C}^{*}(4.32,6.90,7.50)$ are $<140,<69$ and $<91 \mathrm{fs}$, respectively. See (80AJ01) for references. For yields see ${ }^{12} \mathrm{C}$ and (87DO05).

Table 11.16: Energy levels of ${ }^{11} \mathrm{C}^{\mathrm{a}}$

| $\begin{gathered} E_{\mathrm{x}} \text { in }{ }^{11} \mathrm{C} \\ (\mathrm{MeV} \pm \mathrm{keV}) \end{gathered}$ | $J^{\pi} ; T$ | $\tau$ or $\Gamma_{\text {c.m }}$. | Decay | Reactions |
| :---: | :---: | :---: | :---: | :---: |
| 0 | $\frac{3}{2}^{-} ; \frac{1}{2}$ | $\tau_{1 / 2}=20.39 \pm 0.02 \mathrm{~min}$ | $\beta^{+}$ | $1,2,4,6,12,13$, $14,15,16,17,18$, 19, 20, 21, 22, 23, $24,25,26,27,28$, 29, 31 |
| $2.0000 \pm 0.5$ | $\frac{1}{2}^{-}$ | $\tau_{\mathrm{m}}=10.3 \pm 0.7 \mathrm{fs}$ | $\gamma$ | $2,5,6,12,13,14$, $15,16,20,21,22$, $23,24,25,28,29$ |
| $4.3188 \pm 1.2$ | $\frac{5}{2}^{-}$ | $<12$ fs | $\gamma$ | $\begin{aligned} & 2,5,6,12,13,15, \\ & 16,17,19,20,21, \\ & 22,23,28 \end{aligned}$ |
| $4.8042 \pm 1.2$ | $\frac{3}{2}^{-}$ | $<11$ fs | $\gamma$ | $\begin{aligned} & 2,5,12,15,16 \\ & 17,20,21,23,28 \end{aligned}$ |
| $6.3392 \pm 1.4$ | $\frac{1}{2}^{+}$ | $<110 \mathrm{fs}$ | $\gamma$ | 2, 5, 13, 23 |
| $6.4782 \pm 1.3$ | $\frac{7}{2}^{-}$ | $<8 \mathrm{fs}$ | $\gamma$ | $\begin{aligned} & 2,5,6,12,13,15 \\ & 16,20,21,23,27, \\ & 28 \end{aligned}$ |
| $6.9048 \pm 1.4$ | $\frac{5}{2}+$ | $<69$ fs | $\gamma$ | $\begin{aligned} & 2,5,12,13,16 \\ & 21,23 \end{aligned}$ |
| $7.4997 \pm 1.5$ | $\frac{3}{2}^{+}$ | $<91$ fs | $\gamma$ | $\begin{aligned} & 2,5,13,16,21 \\ & 23,28 \end{aligned}$ |
| $8.1045 \pm 1.7$ | $\frac{3}{2}-$ | $0.06 \pm 0.04 \mathrm{fs}^{\text {b }}$ | $\gamma, \alpha$ | 4, 13, 17, 21, 23 |
| $8.420 \pm 2$ | $\frac{5}{2}-$ | $0.43 \pm 0.011 \mathrm{fs}^{\text {b }}$ | $\gamma, \alpha$ | $\begin{aligned} & 2,4,5,12,13,15, \\ & 21,23 \end{aligned}$ |
| $8.655 \pm 8$ | $\frac{7}{2}+$ | $\Gamma \leq 5 \mathrm{keV}$ | ( $\gamma$ ) | 12, 13, 15, 21 |
| $8.699 \pm 10$ | $\frac{5}{2}+$ | $15 \pm 1$ | $\gamma, \mathrm{p}$ | $6,12,13,15$ |
| $9.20 \pm 50$ | $\frac{5}{2}^{+}$ | $500 \pm 100$ | $\gamma, \mathrm{p}$ | 6 |
| $9.65 \pm 50$ | $\left(\frac{3}{2}^{-}\right)$ | $210 \pm 50$ | $\gamma, \mathrm{p}, \alpha$ | 6, 8, 11, 21 |
| $9.78 \pm 50$ | $\left(\frac{5}{2}^{-}\right)$ | $240 \pm 60$ | $\gamma, \mathrm{p}$ | 6, 8, 11, 21 |
| $9.97 \pm 50$ | $\left(\frac{7}{2}^{-}\right)$ | $120 \pm 20$ | $\gamma, \mathrm{p}$ | 6, 21 |
| $10.083 \pm 5$ | $\frac{7}{2}^{+}$ | $\approx 230$ | $\gamma, \mathrm{p}, \alpha$ | $6,8,11,13,21$ |
| $10.679 \pm 5$ | $\frac{9}{2}+$ | $200 \pm 30$ | $\gamma, \mathrm{p}, \alpha$ | $6,8,11,12,21$ |
| $11.03 \pm 30$ | $T=\frac{1}{2}$ | $300 \pm 60$ |  | 21, 23, 28 |

Table 11.16: Energy levels of ${ }^{11} \mathrm{C}^{\text {a }}$ (continued)

| $E_{\mathrm{x}}$ in ${ }^{11} \mathrm{C}$ | $J^{\pi} ; T$ | $\tau$ or $\Gamma_{\text {c.m. }}$ | Decay | Reactions |
| :---: | :---: | :---: | :---: | :--- |
| $(\mathrm{MeV} \pm \mathrm{keV})$ |  |  | $\mathrm{p}, \alpha$ | 11,21 |
| $11.44 \pm 10$ |  | 360 | p | $5,9,17$ |
| $12.16 \pm 40$ | $T=\frac{3}{2}^{2}$ | $270 \pm 50$ | $\gamma, \mathrm{p}$ | 6,23 |
| 12.4 | $\pi=-$ | $1-2 \mathrm{MeV}$ | p | $5,9,17,20,28$ |
| $12.51 \pm 30$ | $\frac{1}{2}^{-} ; \frac{3}{2}$ | $490 \pm 40 \mathrm{keV}$ | $\mathrm{p},{ }^{3} \mathrm{He}, \alpha$ | $6,10,11$ |
| $12.65 \pm 20$ | $\left(\frac{7}{2}^{+}\right)$ | 360 | $\gamma, \mathrm{p}$ | 6 |
| $(13.01)$ |  |  |  | 20,28 |
| $13.33 \pm 60$ |  | $270 \pm 80$ | $\mathrm{p}, \alpha$ | 11,21 |
| 13.4 |  | $1100 \pm 100$ | p | $6,9,17,28$ |
| $13.90 \pm 20$ | $\left(T=\frac{3}{2}\right)$ | $200 \pm 100$ | $\mathrm{n}, \mathrm{p}$ | 7,28 |
| $14.07 \pm 20$ |  | $135 \pm 50$ | $\mathrm{n}, \mathrm{p},{ }^{3} \mathrm{He}$ | $5,7,9,10$ |
| $14.76 \pm 20$ |  | $\approx 450$ | $\gamma, \mathrm{n}, \mathrm{p}$ | $6,7,9,23$ |
| $15.35 \pm 50$ | $\pi=-$ | broad | $\mathrm{n}, \mathrm{p}$ | 7,9 |
| $15.59 \pm 50$ |  | $\approx 450$ | $\gamma, \mathrm{p}$ | 6 |
| 16.7 | $\pi=-$ | $800 \pm 100$ | $\gamma, \mathrm{p}$ | 6 |
| $(18.2)$ |  |  |  | 23 |
| $(23.0)$ |  |  |  |  |
| $(28.0)$ |  |  |  |  |

${ }^{\text {a }}$ See also Table 11.17
${ }^{\mathrm{b}} \Gamma_{\mathrm{c} . \mathrm{m} .}=\Gamma_{\alpha}+\Gamma_{\gamma}=11 \pm 7 \mathrm{eV}$ and $15.2 \pm 3.8 \mathrm{eV}$ for ${ }^{11} \mathrm{C}^{*}(8.10,8.42)$ : see reaction 4.
${ }^{\text {c }}$ I am grateful to Professor F. C. Barker for his comments.
3. ${ }^{7} \mathrm{Li}\left({ }^{7} \mathrm{Li}, 3 \mathrm{n}\right){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-5.050$

At $E\left({ }^{7} \mathrm{Li}\right)=82 \mathrm{MeV}$ no states of ${ }^{11} \mathrm{C}$ are populated (87AL10).
4. ${ }^{7} \operatorname{Be}(\alpha, \gamma){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=7.543$

Table 11.17: Gamma decay of ${ }^{11} \mathrm{C}$ levels ${ }^{\text {a }}$

| $E_{\mathrm{i}}(\mathrm{MeV})$ | $J^{\pi}$ | $\tau_{\mathrm{m}}(\mathrm{fs})$ | $E_{\mathrm{f}}(\mathrm{MeV})$ | Branch |
| :---: | :---: | :---: | :---: | :---: |
| 2.00 | $\begin{aligned} & \frac{1}{2}^{-} \\ & \frac{5}{2}^{-} \\ & \frac{3}{2}^{-} \end{aligned}$ | $10.3 \pm 0.7 \mathrm{fs}$ | 0 | 100 |
| $4.32{ }^{\text {b }}$ |  | $<12^{\text {h }}$ | 0 | 100 |
| 4.80 |  | $<11^{\text {h }}$ | 0 | $85.2 \pm 1.4$ |
|  |  |  | 2.00 | $14.8 \pm 1.4$ |
| $6.34{ }^{\text {c }}$ | $\frac{1}{2}^{+}$ | $<110$ | 0 | $66.5 \pm 2.1$ |
|  |  |  | 2.00 | $33.5 \pm 2.1$ |
| $6.48{ }^{\text {d }}$ | $\frac{7}{2}^{-}$ | $<8^{\text {h }}$ | 0 | $88.5 \pm 1.4$ |
|  |  |  | 4.32 | $11.5 \pm 1.4$ |
| $6.90{ }^{\text {e }}$ | $\frac{5}{2}^{+}$ | < 69 | 0 | $91 \pm 2$ |
|  |  |  | 4.32 | $4.5 \pm 1$ |
|  |  |  | 4.80 | $4.5 \pm 1$ |
| $7.50{ }^{\text {f }}$ | $\frac{3}{2}^{+}$ | $<91$ | 0 | $36 \pm 2$ |
|  |  |  | 2.000 | $64 \pm 2$ |
| $8.10^{\text {i }}$ | $\frac{3}{2}^{-}$ | $0.06 \pm 0.04$ |  | $74 \pm 12$ |
|  |  |  | 0 200 | $26 \pm 5$ |
| $\begin{aligned} & 8.42^{\mathrm{i}, \mathrm{l}} \\ & 8.70^{\mathrm{k}, \mathrm{l}} \end{aligned}$ | $\begin{aligned} & \frac{5}{2}^{-} \\ & \frac{5}{2}^{+} \end{aligned}$ | $0.043 \pm 0.011$ | 0 | $100^{\text {j }}$ |
|  |  |  | 0 | $42 \pm 10$ |
|  |  |  | 4.32 | $42 \pm 10$ |
|  |  |  | 4.80 | $2.4 \pm 1.5$ |
|  |  |  | 6.48 | $13.6 \pm 4.6$ |
| $9.20^{\mathrm{k}}$ | $\frac{5}{2}^{+}$ |  | 0 | $74 \pm 18$ |
|  |  |  | 4.32 | $6 \pm 5$ |
|  |  |  | 6.48 | $20 \pm 10$ |
| $9.65^{\text {g,k }}$ | $\left(\frac{3}{2}^{-}\right)$ |  | 0 | $60 \pm 5$ |
|  |  |  | 4.32 | $32 \pm 10$ |
|  |  |  | 4.80 | $8 \pm 4$ |
| $9.78{ }^{\text {g,k }}$ | $\left(\frac{5}{2}^{-}\right)$ |  | 0 | $76 \pm 16$ |
|  |  |  | 4.32 | $8 \pm 2$ |
|  |  |  | 4.80 | $4 \pm 2$ |
|  |  |  | 6.48 | $12 \pm 4$ |
| $9.97^{\text {k }}$ | $\left(\frac{7}{2}^{-}\right)$ |  | 4.32 | $90 \pm 10$ |

Table 11.17: Gamma decay of ${ }^{11} \mathrm{C}$ levels ${ }^{\text {a }}$ (continued)

| $E_{\mathrm{i}}(\mathrm{MeV})$ | $J^{\pi}$ | $\tau_{\mathrm{m}}(\mathrm{fs})$ | $E_{\mathrm{f}}(\mathrm{MeV})$ | Branch |
| :---: | :---: | :---: | :---: | :---: |
|  |  |  | 6.48 | $10 \pm 7$ |
| $10.08^{\mathrm{k}}$ | $\frac{7}{2}^{+}$ |  | 4.32 | $67 \pm 8$ |
|  |  |  | 6.48 | $13 \pm 6$ |
| $10.68^{\mathrm{k}}$ | $\frac{9}{2}^{+}$ |  | 6.48 | 100 |

${ }^{\text {a }}$ Mostly from (65OL03) and (68EA03): see Table 11.20 in (80AJ01) for other references and additional information.
${ }^{\mathrm{b}}$ Cascade via ${ }^{11} \mathrm{C}^{*}(2.0)$ is $<2 \%$.
${ }^{\text {c }}$ Cascade via ${ }^{11} \mathrm{C}^{*}(4.32)$ is $<7 \%$; that through ${ }^{11} \mathrm{C}^{*}(4.80)$ is $<3 \%$.
${ }^{\mathrm{d}}$ Cascades via ${ }^{11} \mathrm{C}^{*}(2.00,4.80)$ are $<2 \%$.
${ }^{e}$ Cascade via ${ }^{11} \mathrm{C}^{*}(2.00,6.34,6.48)$ are $<1,<5,<5 \%$, respectively. The cascade via ${ }^{11} \mathrm{C}^{*}(4.80)$ is not reported by (65OL03) [they suggest $<3 \%$ ].
${ }^{\mathrm{f}}$ Cascades via ${ }^{11} \mathrm{C}^{*}(4.32,4.80,6.34,6.48,6.90)$ are $<1,<1,<3,<3$ and $<4 \%$.
${ }^{\mathrm{g}}$ See also (79AN16).
${ }^{\mathrm{h}}$ (79AN16). See also (81CA06) for $\tau_{\mathrm{m}}$ of ${ }^{11} \mathrm{C}^{*}(4.32,4.80,6.48)$.
${ }^{\mathrm{i}}$ (84HA13).
${ }^{\mathrm{j}}$ Branching ratio to ${ }^{11} \mathrm{C}^{*}(4.32)$ is $<7 \%$ ( 84 HA 13 ).
k (83WI09).
${ }^{1} \Gamma_{\gamma} / \Gamma=0.20 \pm 0.05,<0.06$ and $\leq 0.1$ for ${ }^{11} \mathrm{C}^{*}(8.42,8.66,8.70)$, respectively: $\Gamma_{\text {total }}($ c.m. $) \leq 4.5$, $\leq 4.5$ and $15 \pm 1 \mathrm{keV}$ (83WI09).

At the resonances at $E_{\alpha}=0.884 \pm 0.008$ and $1.376 \pm 0.003 \mathrm{MeV}\left[{ }^{11} \mathrm{C}^{*}(8.106,8.419)\right], \omega \gamma=$ $0.331 \pm 0.041$ and $3.80 \pm 0.57 \mathrm{eV}, \Gamma_{\gamma}=0.350 \pm 0.056$ and $3.1 \pm 1.3 \mathrm{eV}$ for these two states and $\Gamma_{\alpha}=6_{-2}^{+12}$ and $12.6 \pm 3.8 \mathrm{eV}$, respectively (84HA13). See also (83HA1B, 84YA1A, 85CA41, 88BU01, 88CA26; astrophysics).
5. ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{n}\right){ }^{11} \mathrm{C}$

$$
Q_{\mathrm{m}}=7.5572
$$

Reported neutron groups are listed in Table 11.16 of (68AJ02). Angular distributions have been studied in the range $E\left({ }^{3} \mathrm{He}\right)=1.3$ to 13 MeV : see (80AJ01). The dominant $L$-values are 0 for ${ }^{11} \mathrm{C}^{*}(0,8.10), 1$ for ${ }^{11} \mathrm{C}^{*}(6.34,7.50), 2$ for ${ }^{11} \mathrm{C}^{*}(2.00,4.32,4.80,6.48,8.42)$ and 3 for ${ }^{11} \mathrm{C}^{*}(6.90)$. Neutron groups to $T=\frac{3}{2}$ states have been reported at $E_{\mathrm{x}}=12.17 \pm 0.05$ [see, however, reaction 28], $12.55 \pm 0.05 \mathrm{MeV}$ and $14.7 \pm 0.1 \mathrm{MeV}$ : see Table 11.18.

Gamma branching ratios and multipolarities for ${ }^{11} \mathrm{C}$ levels up to $E_{\mathrm{x}}=7.5 \mathrm{MeV}$ have been studied by (65OL03): see Table 11.17. Together with evidence from reactions 12 and 21 they lead to assignments of $J^{\pi}=\frac{1}{2}^{-}, \frac{5}{2}^{-}, \frac{3}{2}^{-}, \frac{1}{2}^{+}, \frac{7}{2}^{-}, \frac{5}{2}^{+}, \frac{3}{2}^{+}$for ${ }^{11} \mathrm{C}^{*}(2.00,4.32,4.80,6.34$,

Table 11.18: $T=\frac{3}{2}$ states in ${ }^{11} \mathrm{C}^{\text {a }}$

| Reaction | $E_{\mathrm{x}}(\mathrm{MeV})$ | $\Gamma_{\text {c.m. }}(\mathrm{keV})$ |
| :---: | :---: | :---: |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{n}\right)^{11} \mathrm{C}$ | $12.17 \pm 0.05$ | $200 \pm 100$ |
| ${ }^{10} \mathrm{~B}(\mathrm{p}, \mathrm{p} /)^{10} \mathrm{~B}^{*}$ | $12.20 \pm 0.10$ |  |
| ${ }^{11} \mathrm{~B}\left({ }^{3} \mathrm{He}, \mathrm{t}\right){ }^{11} \mathrm{C}$ | $\underline{12.15 \pm 0.05}$ | $\underline{290 \pm 50}$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{n}\right){ }^{11} \mathrm{C}$ | $12.16 \pm 0.04$ | ${ }^{\mathrm{b}}$ |
| ${ }^{10} \mathrm{~B}\left(\mathrm{p}, \mathrm{p}_{2}\right)^{10} \mathrm{~B}^{*}$ | $12.45 \pm 0.05$ | $350 \pm 50^{\mathrm{b}}$ |
| ${ }^{11} \mathrm{~B}\left({ }^{3} \mathrm{He}, \mathrm{t}\right){ }^{11} \mathrm{C}$ | $12.57 \pm 0.07$ | $400 \pm 100$ |
| ${ }^{13} \mathrm{C}(\mathrm{p}, \mathrm{t}){ }^{11} \mathrm{C}$ | $12.47 \pm 0.06$ | $370 \pm 90$ |
| ${ }^{13} \mathrm{C}(\mathrm{p}, \mathrm{t}){ }^{11} \mathrm{C}$ | $\underline{12.48 \pm 0.04}$ | $550 \pm 50$ |
| ${ }^{9} \mathrm{Be}\left({ }^{3} \mathrm{He}, \mathrm{n}\right)^{11} \mathrm{C}$ | $12.51 \pm 0.03^{\mathrm{b}}$ | $490 \pm 60$ |
| ${ }^{11} \mathrm{~B}\left({ }^{3} \mathrm{He}, \mathrm{n}\right)^{11} \mathrm{C}$ | $13.7 \pm 0.1$ |  |

${ }^{\text {a }}$ See also Table 11.14 for $T=\frac{3}{2}$ states in ${ }^{11} \mathrm{~B}$, and Table 11.21 in (80AJ01). For references see Table 11.19 in (85AJ01).
${ }^{\mathrm{b}}$ Mean.
$6.48,6.90,7.50$ ): see ( 65 OL 03 ) and reaction 3 in ( 68 AJ 02 ) for a summary of the evidence concerning these assignments. See (80AJ01) for references. See also ${ }^{12} \mathrm{C}$ and (84SU1E).
6. ${ }^{10} \mathrm{~B}(\mathrm{p}, \gamma){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=8.6896$

This reaction has been investigated for $E_{\mathrm{p}}=0.07$ to 17.0 MeV . Reported resonances are displayed in Table 11.19. Observed capture $\gamma$-rays are displayed in Table 11.17 [see also for $\tau_{\mathrm{m}}$ measurements]. Capture measurements for $E_{\mathrm{p}}=0.07$ to 2.20 MeV are consistent with five new resonances (see Tables 11.19 and 11.17), the lowest two (at $E_{\mathrm{p}}=10$ and 560 keV ) of which are s-wave resonances. Thermonuclear reaction rates for $T=(0.01 \rightarrow 5) \times 10^{9} \mathrm{~K}$ are deduced from the results (83WI09; see also for spectroscopic factors).

The $90^{\circ}$ yield of $\gamma_{0}$ has been measured for $E_{\mathrm{p}}=2.6$ to 17 MeV and angular distributions have been obtained for $E_{\mathrm{p}}=2.8$ to 14 MeV . The excitation function is consistent with the giant resonance centered at $E_{\mathrm{x}} \approx 16 \mathrm{MeV}$. In addition to weak structures at $E_{\mathrm{p}}=4.75$ MeV and 10.5 MeV , there are three major peaks at $E_{\mathrm{p}}=4.1,7.0$ and $8.8 \mathrm{MeV}(\Gamma=1-2$ $\mathrm{MeV})\left[E_{\mathrm{x}}=12.4,15.0,16.7 \mathrm{MeV}\right]$. At ${ }^{11} \mathrm{C}^{*}(12.4)$, the $\gamma_{0}$ angular distribution is essentially isotropic: $\Gamma_{\mathrm{p}} \Gamma_{\gamma} / \Gamma \approx 200 \mathrm{eV}, \Gamma_{\gamma} \approx 5 \mathrm{keV}$ (assuming $\Gamma_{\mathrm{p}} \approx 10 \mathrm{keV}$ ). The $E_{\mathrm{p}}=4.1 \mathrm{MeV}$ resonance is probably part of the E 1 giant resonance and is formed by s-wave capture. At

Table 11.19: Resonances ${ }^{\text {a }}$ in ${ }^{10} \mathrm{~B}+\mathrm{p}$

| $E_{\text {res }}$ <br> $(\mathrm{MeV} \pm \mathrm{keV})$ | $E_{\mathrm{x}}$ <br> $(\mathrm{MeV})$ | $J^{\pi}$ | $\Gamma_{\text {lab }}$ <br> $(\mathrm{keV})$ | Decay |
| :---: | :---: | :---: | :---: | :---: |
| $0.010 \pm 2^{\mathrm{b}}$ | $8.699 \pm 10$ | $\frac{5}{2}^{+}$ | $16 \pm 1^{\mathrm{c}}$ | $\gamma$ |
| $0.56 \pm 60^{\mathrm{b}}$ | $9.20 \pm 50$ | $\frac{5}{2}^{+}$ | $550 \pm 100$ | $\gamma$ |
| $1.05 \pm 60^{\mathrm{b}}$ | $9.64 \pm 50$ | ${\left(\frac{3}{2}^{-}\right)}^{5^{-}}$ | $230 \pm 50$ | $\gamma,\left(\mathrm{p}_{0}, \alpha_{0}\right)$ |
| $1.20 \pm 50^{\mathrm{b}}$ | $9.78 \pm 50$ | $\left.\left(\frac{5}{2}\right)^{-}\right)$ | $260 \pm 60$ | $\gamma,\left(\mathrm{p}_{0}, \alpha_{0}\right)$ |
| $1.41 \pm 50^{\mathrm{b}}$ | $9.97 \pm 50$ | $\left(\frac{7}{2}^{-}\right)$ | $130 \pm 20$ | $\gamma$ |
| $1.533 \pm 5$ | 10.083 | $\frac{7}{2}^{+}$ | $\approx 250$ | $\mathrm{p}_{0}, \alpha_{0}, \alpha_{1}$ |
| $2.189 \pm 5$ | 10.679 | $\frac{9}{2}^{+}$ | $220 \pm 30$ | $\mathrm{p}_{0}, \alpha_{0}, \alpha_{1}$ |
| $3.03 \pm 10$ | 11.44 |  | 400 | $\alpha_{0}, \alpha_{1}$ |
| $3.9 \pm 10$ | 12.20 | $T=\frac{3}{2}$ |  | $\mathrm{p}_{2}$ |
| $4.1 \pm 100$ | 12.45 | $T=\frac{3}{2}$ | $440 \pm 100$ | $\mathrm{p}_{2}$ |
| $4.1^{\mathrm{d}, \mathrm{e}}$ | 12.4 | $\pi=-$ | $1-2 \mathrm{MeV}$ | $\gamma_{0}$ |
| $4.36 \pm 20$ | 12.65 | $\left(\frac{7^{+}}{2}\right)$ | 400 | $\gamma_{1}, \alpha_{0}, \alpha_{1},{ }^{3} \mathrm{He}$ |
| $(4.75)$ | $(13.01)$ |  | $\gamma_{0}$ |  |
| 5.2 | 13.4 |  | $1200 \pm 100$ | $\alpha_{0}, \alpha_{1}$ |
| $5.73 \pm 20$ | 13.90 |  | $\approx 500$ | $\gamma_{1}, \mathrm{p}$ |
| $5.92 \pm 20$ | 14.07 |  | broad | n |
| $6.68 \pm 40$ | 14.76 |  | $\approx 500$ | $\mathrm{n}, \mathrm{p},{ }^{3} \mathrm{He}$ |
| $7.33 \pm 50^{\mathrm{e}}$ | 15.35 | $\pi=-$ | broad | $\gamma_{0}, \mathrm{n}, \mathrm{p}$ |
| $7.60 \pm 50$ | 15.59 |  | $\approx 500$ | $\mathrm{n}, \mathrm{p}$ |
| $8.8^{\mathrm{e}}$ | 16.7 | $\pi=-$ | $900 \pm 100$ | $\gamma_{0}$ |
| $(10.5)$ | $(18.2)$ |  |  | $\gamma_{0}$ |

[^3]the two higher resonances the angular distributions are characteristic of E1 giant resonances in light nuclei. The ${ }^{10} \mathrm{~B}\left(\mathrm{p}, \gamma_{1}\right)$ cross section is small for $E_{\mathrm{p}}=2.6$ to 17 MeV : see (80AJ01). See also (84YA1A, 85CA41, 88CA26; astrophysics).
7. ${ }^{10} \mathrm{~B}(\mathrm{p}, \mathrm{n}){ }^{10} \mathrm{C}$
$Q_{\mathrm{m}}=-4.4305$
$$
E_{\mathrm{b}}=8.6896
$$

The total (p, n) cross section has been measured to $E_{\mathrm{p}}=10.6 \mathrm{MeV}$ : broad maxima are observed at $E_{\mathrm{p}}=5.92 \pm 0.02,6.68 \pm 0.04,7.33 \pm 0.05$ and $7.60 \pm 0.05 \mathrm{MeV}$ (see Table 11.19). The cross section for formation of ${ }^{10} \mathrm{C}$ (g.s.) measured up to 12 MeV shows similar behavior to 8 MeV . At $E_{\mathrm{p}} \approx 8 \mathrm{MeV}$, a sharp maximum is observed. The cross section for production of $3.35 \mathrm{MeV} \gamma$-rays $\left(\right.$ from $\left.{ }^{10} \mathrm{C}^{*}\right)$ does not appear to show structure for $E_{\mathrm{p}}=8.5$ to 12 MeV . For references see (80AJ01). For $\mathrm{n}_{0}$ and $\mathrm{n}_{1}$ excitiation curves from $E_{\mathrm{p}}=13.7$ to 14.7 MeV see (85SC08). See also (84BA1R, 84BA1U).
8. ${ }^{10} \mathrm{~B}(\mathrm{p}, \mathrm{p})^{10} \mathrm{~B}$

$$
E_{\mathrm{b}}=8.6896
$$

Below $E_{\mathrm{p}}=0.7 \mathrm{MeV}$ the scattering can be explained in terms of pure s-wave potential scattering but the possibility of a state near $E_{\mathrm{p}}=0.27 \mathrm{MeV}\left(E_{\mathrm{x}}=8.95 \mathrm{MeV}\right)$ cannot be excluded. The elastic scattering then shows two conspicuous anomalies at $E_{\mathrm{p}}=1.50 \pm$ 0.02 MeV and at $2.18 \mathrm{MeV}\left[E_{\mathrm{x}}=10.05\right.$ and 10.67 MeV$]$ with $J^{\pi}=\frac{7}{2}^{+}$and $\frac{9}{2}^{+}$: see Table 11.19. At higher energies (to $E_{\mathrm{p}}=10.5 \mathrm{MeV}$ ) a single broad resonance is reported at $E_{\mathrm{p}} \approx 5 \mathrm{MeV}$. Polarization measurements are reported at 30.3 MeV : optical model parameters have been derived. The depolarization parameter $D$ has been measured for polarized protons at 26 and 50 MeV . For references see (80AJ01, 85AJ01). See also (84BA1U) and (86MU1D).
9. ${ }^{10} \mathrm{~B}\left(\mathrm{p}, \mathrm{p}^{\prime}\right)^{10} \mathrm{~B}$

$$
E_{\mathrm{b}}=8.6896
$$

The yield of $\gamma_{1}\left[\right.$ from $\left.{ }^{10} \mathrm{~B}^{*}(0.72)\right]$ rises monotonically from $E_{\mathrm{p}}=1.5$ to 4.1 MeV and then shows resonance behavior at $E_{\mathrm{p}}=4.36$ and 5.73 MeV : see Table 11.19. For $E_{\mathrm{p}}=6$ to 12 MeV , the cross section for $\gamma_{1}$ shows several sharp maxima superposed on a broad maximum $(\Gamma \approx 2.5 \mathrm{MeV})$ at $E_{\mathrm{p}} \approx 7.2 \mathrm{MeV}$. See however (75AJ02). Yields of five other $\gamma$-rays involved in the decay of ${ }^{10} \mathrm{~B}^{*}(1.74,2.16,3.59,5.18)$ have also been measured in the range $E_{\mathrm{p}}=4$ to 12 MeV [see (75AJ02)].

Excitation curves for the $\mathrm{p}_{1}, \mathrm{p}_{2}$ and $\mathrm{p}_{3}$ groups have been measured for $E_{\mathrm{p}}=3.5$ to 5.0 MeV . Possible resonances are observed in the $\mathrm{p}_{2}$ yield [to the $T=1$ state ${ }^{10} \mathrm{~B}^{*}(1.74)$ ] corresponding to the first $T=\frac{3}{2}$ states at $E_{\mathrm{x}}=12.16$ [see however reaction 28] and 12.50 MeV [see Table 11.18: these do not occur in the yield of $\mathrm{p}_{1}$ and $\mathrm{p}_{3}$. Yield curves for inelastically scattered protons have also been measured at $E_{\mathrm{p}}=5.0$ to $16.4 \mathrm{MeV}\left(\mathrm{p}_{1}, \mathrm{p}_{2}\right.$,
$\mathrm{p}_{3}$ ), 6.6 to $16.4 \mathrm{MeV}\left(\mathrm{p}_{4}\right), 8.9$ to $16.4 \mathrm{MeV}\left(\mathrm{p}_{5}\right)$ and 10.9 to $16.4 \mathrm{MeV}\left(\mathrm{p}\right.$ to ${ }^{10} \mathrm{~B}^{*}(6.03)$ ): the principal feature for all groups, except that to ${ }^{10} \mathrm{~B}^{*}(6.03)$, is a structure at $E_{\mathrm{p}} \approx 7.5 \mathrm{MeV}$, $\Gamma \approx 4 \mathrm{MeV}$. In addition narrower structures are observed, including three at $E_{\mathrm{p}}=5.75,6.90$ and $7.80 \mathrm{MeV}( \pm 0.2 \mathrm{MeV})$ and widths of $\approx 500 \mathrm{keV}$. For references see (80AJ01, 85AJ01).
10. (a) ${ }^{10} \mathrm{~B}(\mathrm{p}, \mathrm{d})^{9} \mathrm{~B}$
$Q_{\mathrm{m}}=-6.212$
$E_{\mathrm{b}}=8.6896$
(b) ${ }^{10} \mathrm{~B}\left(\mathrm{p},{ }^{3} \mathrm{He}\right)^{8} \mathrm{Be}$
$Q_{\mathrm{m}}=-0.5330$

Polarization measurements (reaction (a)) have been carried out at $E_{\mathrm{p}}=49.6 \mathrm{MeV}$ for the deuterons to ${ }^{9} \mathrm{~B}^{*}(0,2.36)$ : see ( 75 AJ 02 ). In reaction (b) two strong maxima are observed at $E_{\mathrm{p}} \approx 4.5$ and 6.5 MeV : see Table 11.19. See also (75AJ02).
11. ${ }^{10} \mathrm{~B}(\mathrm{p}, \alpha)^{7} \mathrm{Be} \quad Q_{\mathrm{m}}=1.1462 \quad E_{\mathrm{b}}=8.6896$

The total cross section for this reaction has been measured for $E_{\mathrm{p}}=60$ to 180 keV : the extrapolated cross section at the Gamow energy, taken to be 19.1 keV , is $\approx 10^{-12} \mathrm{~b}$. The thick target yield for $E_{\mathrm{p}}=75 \mathrm{keV}$ to 3 MeV shows that the ${ }^{7} \mathrm{Be}$ yield constitutes a potential problem if natural boron is used as fuel in CTR devices.

The parameters of observed resonances are displayed in Table 11.19. The ground state $\left(\alpha_{0}\right) \alpha$-particles exhibit broad resonances at $E_{\mathrm{p}}=1.17,1.53,2.18,3.0,4.4,5.1$ and 6.3 MeV . Alpha particles to ${ }^{7} \mathrm{Be}^{*}(0.43)\left[\alpha_{1}\right]$ and $0.43-\mathrm{MeV} \gamma$-rays exhibit all but the 1.2 MeV resonance: see (75AJ02). A broad maximum dominates the region from $E_{\mathrm{p}}=4 \mathrm{MeV}$ to about 7.5 MeV . A study of the yield of $0.43 \mathrm{MeV} \gamma$-rays for $E_{\mathrm{p}}=2.0$ to 4.1 MeV suggests that the 3.0 MeV resonance, which is asymmetric, is due to two broad states. A weak structure at $E_{\mathrm{p}}=2.5 \mathrm{MeV}$ is also reported. For references see (80AJ01, 85AJ01). See also ${ }^{7} \mathrm{Be}$ in (88AJ01) and (84YA1A, 85CA41; astrophysics).
12. ${ }^{10} \mathrm{~B}(\mathrm{~d}, \mathrm{n}){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=6.4650$

Table 11.20 presents the results obtained in this reaction and in the $\left({ }^{3} \mathrm{He}, \mathrm{d}\right)$ reaction. Information on $\tau_{\mathrm{m}}$ and on the $\gamma$-decay of ${ }^{11} \mathrm{C}$ states is displayed in Table 11.16: see (68AJ02, 75AJ02) for references. See also (86WE1E; applied) and ${ }^{12} \mathrm{C}$.
13. ${ }^{10} \mathrm{~B}\left({ }^{3} \mathrm{He}, \mathrm{d}\right){ }^{11} \mathrm{C} \quad Q_{\mathrm{m}}=3.1961$

Table 11.20: Energy levels of ${ }^{11} \mathrm{C}$ from ${ }^{10} \mathrm{~B}(\mathrm{~d}, \mathrm{n}){ }^{11} \mathrm{C}$ and ${ }^{10} \mathrm{~B}\left({ }^{3} \mathrm{He}, \mathrm{d}\right){ }^{11} \mathrm{C}^{\mathrm{a}}$

| $\begin{gathered} E_{\mathrm{x}} \\ (\mathrm{MeV} \pm \mathrm{keV}) \end{gathered}$ | $J^{\pi}$ | $l{ }^{\text {b }}$ | $l^{\text {c }}$ | $S_{\text {d, }}{ }^{\text {c }}$ | $S^{3}{ }^{\text {He, }}{ }^{\text {c }}$ | $l^{\text {d }}$ | $S_{3}{ }^{\text {He, }}$ d ${ }^{\text {d }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | $\frac{3}{2}^{-}$ | 1 | 1 | 1.12 | 0.88 | 1 | 1.09 |
| $2.0006 \pm 0.9$ | $\frac{1}{2}^{-}$ | (1) | (1) | (0.18) | (0.036) |  |  |
|  |  |  | (3) |  | $\leq 0.09$ | (3) | $<0.40$ |
| $4.322 \pm 10$ | $\frac{5}{2}^{-}$ | 1 | 1 | 0.27 | 0.20 | 1 | 0.17, 0.19 |
| $4.808 \pm 10$ | $\frac{3}{2}^{-}$ | 1 | 1 | < 0.02 |  | (1) | $<0.08$ |
|  |  |  |  |  |  | (3) | $<0.35$ |
| $6.345 \pm 10$ | $\begin{aligned} & \frac{1}{2}^{+} \\ & \frac{7}{2}^{-} \\ & \frac{5}{2}^{+} \end{aligned}$ | $\begin{gathered} 1 \\ (1) \end{gathered}$ | 2 | 0.86 | 0.07 | 2 | 0.08 |
| $6.476 \pm 10$ |  |  | 1 |  | 0.56 | 1 | 0.73, 0.79 |
| $6.903 \pm 10$ |  |  |  |  |  | 2 | 0.06 |
|  |  |  |  |  |  | 0 | $<0.04$ |
| $7.498 \pm 10$ | $\frac{3}{2}^{+}$ |  |  |  |  | 2 | 0.08 |
| $8.107 \pm 10$ | $\frac{3}{2}^{-}$ |  |  |  |  | 1 | 0.07 |
| $8.424 \pm 8$ | $\frac{5}{2}^{-}$ | 1 | 1 | 0.65 | 0.46 | 1 | 0.73, 0.79 |
| $8.655 \pm 8$ | $\frac{5}{2}^{+}$ | 0 | 0 | $\underline{0.84}$ | 0.45 |  |  |
|  |  |  | 2 | 0.8 | 0.32 |  |  |
|  | $\frac{7}{2}+$ |  | 0 | $\underline{0.63}$ | 0.33 | 2 | 0.41 |
|  |  |  | 2 | 0.6 | $\underline{0.24}$ | 0 | $<0.34$ |
| $8.701 \pm 20$ | $\begin{aligned} & \frac{5}{2}^{+} \\ & \frac{7}{2}^{+} \end{aligned}$ | (0) | 0 | $\underline{0.40}$ | 0.14 | 0 | <0.8 |
|  |  |  | 2 | $\leq 0.2$ | 0.13 |  |  |
|  |  |  | 0 | $\underline{0.30}$ | 0.11 |  |  |
|  |  |  | 2 | $\leq 0.15$ | 0.10 |  |  |
| 10.08 |  |  |  |  |  |  |  |
| $10.68{ }^{\text {e }}$ |  |  | $(0,2)$ |  |  |  |  |

a See Table 11.23 in (80AJ01) for references.
${ }^{\mathrm{b}}$ From (d,n) work summarized in Table 11.20 of (68AJ02).
${ }^{\text {c }} S_{\mathrm{d}, \mathrm{n}}$ obtained at $E_{\mathrm{d}}=5.8 \mathrm{MeV}, S_{3} \mathrm{He}, \mathrm{d}$ obtained at $E\left({ }^{3} \mathrm{He}\right)=11.0 \mathrm{MeV}$ [both $\pm 30 \%$ ]. When $S_{\mathrm{d}, \mathrm{n}}$ and $S_{3^{\mathrm{He}, \mathrm{d}}}$ differ appreciably, the more reliable value is underlined.
${ }^{\text {d }} E\left({ }^{3} \mathrm{He}\right)=21 \mathrm{MeV}$; when two values are shown for $S_{3^{3} \mathrm{He}, \mathrm{d}}$, they are in order of descending j .
${ }^{e} \Gamma \approx 200 \mathrm{keV}$.

Table 11.20 displays the information derived from this reaction and from the (d, n) reaction. The study of the angular distributions of the deuterons to ${ }^{11} \mathrm{C}^{*}(8.66,8.70)$ shows that these levels are the analogs, respectively, of ${ }^{11} \mathrm{~B}^{*}(9.19,9.27)$ whose $J^{\pi}$ are $\frac{7}{2}^{+}$and $\frac{5}{2}^{+}$ [the ${ }^{11} \mathrm{~B}$ states were studied in the ( $\mathrm{d}, \mathrm{p}$ ) reaction]: $\Gamma_{\text {c.m. }}$ are $\ll 9 \mathrm{keV}$ and $15 \pm 1 \mathrm{keV}$, respectively, for ${ }^{11} \mathrm{C}^{*}(8.66,8.70)$ : see (75AJ02) for references.
14. ${ }^{10} \mathrm{~B}(\alpha, \mathrm{t}){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-11.1244$

Angular distributions have been measured at $E_{\alpha}=25.1$ and 56 MeV [see (80AJ01)] and at 24.8 and 30.1 MeV (83VA28; $\mathrm{t}_{0}, \mathrm{t}_{1}$ ). See also (84BE23; theor.)
15. ${ }^{10} \mathrm{~B}\left({ }^{7} \mathrm{Li},{ }^{6} \mathrm{He}\right){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-1.285$

Angular distributions of ${ }^{6} \mathrm{He}$ ions have been measured at $E\left({ }^{7} \mathrm{Li}\right)=3.0$ to 3.8 MeV and at 24 MeV [to $\left.{ }^{11} \mathrm{C}^{*}(0,4.32,6.48)\right] .{ }^{11} \mathrm{C}^{*}(2.0,4.80,8.42,8.66+8.70)$ are also populated: see (80AJ01) for references.
16. ${ }^{11} \mathrm{~B}(\mathrm{p}, \mathrm{n}){ }^{11} \mathrm{C}$

$$
Q_{\mathrm{m}}=-2.7646
$$

Angular distributions have been measured at many energies up to 49.5 MeV [see (80AJ01, 85AJ01)] and at $E_{\mathrm{p}}=14.0,14.3$ and $14.6 \mathrm{MeV}\left(85 \mathrm{SC} 08 ; \mathrm{n}_{0}, \mathrm{n}_{1}, \mathrm{n}_{2}, \mathrm{n}_{3},\left(\mathrm{n}_{4+5}\right), \mathrm{n}_{6}, \mathrm{n}_{7}\right), 15.8$ and 18.6 MeV ( $88 \mathrm{KA} 30 ; \mathrm{n}_{0}, \mathrm{n}_{1}$ ) and 16 to 26 MeV (85GR09; $\mathrm{n}_{0}, \mathrm{n}_{1}, \mathrm{n}_{2}, \mathrm{n}_{3}$ ) [see also for a study of the GT matrix elements]. For $0^{\circ}$ cross sections at $E_{\mathrm{p}}=492$ and 590 MeV see (89RA09). See also ${ }^{12} \mathrm{C}$, (84BA1R, 85GU1C), (88CA26; astrophysics), (84TA1F, 86MU1D, 87RA32) and (86HU06; theor.).
17. ${ }^{11} \mathrm{~B}\left({ }^{3} \mathrm{He}, \mathrm{t}\right){ }^{11} \mathrm{C}$

$$
Q_{\mathrm{m}}=-2.0008
$$

Angular distributions of $\mathrm{t}_{0}$ and $\mathrm{t}_{1}$ have been measured at $E\left({ }^{3} \mathrm{He}\right)=10,14$, and 217 MeV [the latter also for the triton groups to ${ }^{11} \mathrm{C}^{*}(4.3,4.8,6.48,8.10]$ and at $E\left({ }^{3} \overrightarrow{\mathrm{He}}\right)=33 \mathrm{MeV}$. At $E\left({ }^{3} \mathrm{He}\right)=26 \mathrm{MeV}$ the known states of ${ }^{11} \mathrm{C}$ below $E_{\mathrm{x}}=11 \mathrm{MeV}$ are populated and triton groups are also observed to the possibly $T=\frac{3}{2}$ states displayed in Table 11.18 as well as a state at 14.15 MeV . For references see (80AJ01, 85AJ01).
18.
(a) ${ }^{12} \mathrm{C}(\gamma, \mathrm{n})^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-18.7215$
(b) ${ }^{12} \mathrm{C}(\mathrm{e}, \text { en })^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-18.7215$

The fraction of transitions to the ground and to excited states of ${ }^{11} \mathrm{C}$ [and to ${ }^{11} \mathrm{~B}$ states reached in the $(\gamma, \mathrm{p})$ reaction] has been measured at $E_{\mathrm{bs}}=24.5,27,33$ and 42 MeV : the ground state is predominantly populated. The population of analog states in the ( $\gamma, \mathrm{n}$ ) and $(\gamma, \mathrm{p})$ reactions are similar. And a significant decay strength is found to the positiveparity states with $6<E_{\mathrm{x}}<8 \mathrm{MeV}$. In general the main contribution to the strength of the transitions to the various excited states of ${ }^{11} \mathrm{~B},{ }^{11} \mathrm{C}$ lies in rather localized energy bands in ${ }^{12} \mathrm{C}$ which are a few MeV wide (70ME17). See also reactions 24 and 25 in (80AJ01) (85AJ01), (88HA01) in ${ }^{12} \mathrm{C}$ and (85CA32, 87GOZ0, 87GO37, 87VA15; theor.).
19. (a) ${ }^{12} \mathrm{C}\left(\pi^{ \pm}, \pi^{ \pm} \mathrm{n}\right)^{11} \mathrm{C}$

$$
Q_{\mathrm{m}}=-18.7215
$$

(b) ${ }^{12} \mathrm{C}(\mathrm{n}, 2 \mathrm{n})^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-18.7215$
(c) ${ }^{12} \mathrm{C}(\mathrm{p}, \mathrm{pn}){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-18.7215$
${ }^{11} \mathrm{C}^{*}(4.32)\left[\frac{5}{2}^{-}\right]$(and the analog state in ${ }^{11} \mathrm{~B}$ ) is surprisingly strongly populated for $E_{\pi^{+}}=$ 60 to 300 MeV : see (80AJ01, 85AJ01). For reaction (b) see ${ }^{13} \mathrm{C}$ in (86AJ01). In reaction (c) at 1 GeV the separation energy between 6 and 13 MeV broad $1 \mathrm{p}_{3 / 2}$ and $1 \mathrm{~s}_{1 / 2}$ groups is $\approx 17$ MeV (85BE1J, 85DO1B). See also ${ }^{12} \mathrm{C}$ and (84GO1F).
20. ${ }^{12} \mathrm{C}\left(\pi^{+}, \mathrm{p}\right){ }^{11} \mathrm{C} \quad Q_{\mathrm{m}}=121.629$

Angular distributions at $E_{\pi^{+}}=49.3,90$ and 180 MeV have been obtained to ${ }^{11} \mathrm{C}^{*}(0,2.0$, $4.3+4.8,6.5,8.5)$. At the same momentum transfer this reaction and the ( $\mathrm{p}, \mathrm{d}$ ) reaction give similar intensities to the low lying states of ${ }^{11} \mathrm{C} . T=\frac{3}{2}$ states have been suggested at $E_{\mathrm{x}}=12.5 \pm 0.3$ and 13.3 MeV : see (85AJ01). See also (82DO01).
21. ${ }^{12} \mathrm{C}(\mathrm{p}, \mathrm{d}){ }^{11} \mathrm{C} \quad Q_{\mathrm{m}}=-16.4972$

Angular distributions have been measured for $E_{\mathrm{p}}=19$ to 800 MeV [see (68AJ02, 75AJ02, 80AJ01, 85AJ01) for references], at $E_{\overrightarrow{\mathrm{p}}}=497 \mathrm{MeV}\left(84 \mathrm{OH} 06 ; \mathrm{p}_{0}\right.$; also $\left.A_{\mathrm{y}}\right)$ and at $E_{\mathrm{p}}=800$ MeV (84SM04; to ${ }^{11} \mathrm{C}^{*}(0,2.0,4.3,4.8,6.5,8.1,8.66+8.70,9.98 \pm 0.2,10.56 \pm 0.2)$ ). In the latter experiment ${ }^{11} \mathrm{C}^{*}(8.4)$ and a state at $13.22 \pm 0.25 \mathrm{MeV}(\Gamma \approx 2 \mathrm{MeV})$ are also reported (84SM04). Earlier observed states of ${ }^{11} \mathrm{C}$ are displayed in Table 11.24 of (80AJ01). See also ${ }^{13} \mathrm{~N}$ in (90AJ01), (87CA20) and (84RE1A).
22. ${ }^{12} \mathrm{C}(\mathrm{d}, \mathrm{t})^{11} \mathrm{C}$

$$
Q_{\mathrm{m}}=-12.4645
$$

Table 11.21: Levels of ${ }^{11} \mathrm{C}$ from ${ }^{12} \mathrm{C}\left({ }^{3} \mathrm{He}, \alpha\right){ }^{11} \mathrm{C}{ }^{\text {a }}$

| $E_{\mathrm{x}}$ <br> $(\mathrm{MeV} \pm \mathrm{keV})$ | $l$ | $S_{\text {rel }}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $E\left({ }^{3} \mathrm{He}\right)=16 \mathrm{MeV}$ | 24 MeV | 28 MeV | 35.6 MeV |
| 0 | 1 | 1 | 1 | 1 | 1.00 |
| $1.999 \pm 4$ | 1 | 0.10 | $\leq 0.6$ | $\leq 0.6$ | 0.19 |
| $4.3188 \pm 1.2$ | 3 | 0.057 | $(0.04)$ | $(0.06)$ | $(0.031)$ |
| $4.8042 \pm 1.2$ | 1 | 0.11 | 0.22 | 0.22 | 0.13 |
| $6.3392 \pm 1.4$ | 0 | $0.003^{\mathrm{b}}$ | $\leq 0.07$ | $\leq 0.07$ | $(\lesssim 0.2)$ |
| $6.4782 \pm 1.4$ | 3 | $0.11^{\mathrm{b}}$ | 0.06 | $(0.06)$ | $(0.21)$ |
| $6.9048 \pm 1.4$ | 2 | 0.018 | $(0.15)$ | $(0.17)$ | $(0.054)$ |
| $7.4997 \pm 1.5$ | 2 | $0.006^{\mathrm{b}}$ | $(0.07)$ | $(0.09)$ | $(0.046)$ |
| $8.1045 \pm 1.7$ | 1 | $0.017^{\mathrm{b}, \mathrm{c}}$ |  |  | $(0.035)$ |
| 8.42 | 3 | $0.034^{\mathrm{b}, \mathrm{d}}$ |  |  | $(0.041)$ |

${ }^{\text {a }}$ See Table 11.17 for $\gamma$-decay work. Higher excited states are also reported: see text. See Table 11.25 in (80AJ01) for references and for additional information.
${ }^{\mathrm{b}}$ At $E\left({ }^{3} \mathrm{He}\right)=18 \mathrm{MeV}$.
${ }^{\text {c }}$ Assuming $J^{\pi}=\frac{3}{2}{ }^{-}$.
${ }^{\mathrm{d}}$ Assuming $J^{\pi}=\frac{5}{2}^{-}$.

At $E_{\mathrm{d}}=28 \mathrm{MeV}$ the $\mathrm{t}_{0}$ angular distribution has been measured and a detailed comparison has been made with the results for the mirror reaction ${ }^{12} \mathrm{C}\left(\mathrm{d},{ }^{3} \mathrm{He}\right){ }^{11} \mathrm{~B}$. At $E_{\mathrm{d}}=29 \mathrm{MeV}$ the $\mathrm{t}_{0}$ angular distribution leads to spectroscopic factor $C^{2} S=2.82$ or 3.97 depending on different sets of parameters for ${ }^{11} \mathrm{C}_{\text {g.s. }} .{ }^{11} \mathrm{C}^{*}(2.0,4.32)$ are also populated. See also ${ }^{14} \mathrm{~N}$ in (86AJ01), (80AJ01) for references, and (84KO1M).
23. (a) ${ }^{12} \mathrm{C}\left({ }^{3} \mathrm{He}, \alpha\right){ }^{11} \mathrm{C}$

$$
\begin{aligned}
& Q_{\mathrm{m}}=1.8560 \\
& Q_{\mathrm{m}}=-17.9577
\end{aligned}
$$

(b) ${ }^{12} \mathrm{C}\left({ }^{3} \mathrm{He}\right.$, tp $){ }^{11} \mathrm{C}$

Angular distributions have been measured at many energies to $E\left({ }^{3} \mathrm{He}\right)=217 \mathrm{MeV}$ [see (68AJ02, 75AJ02, 80AJ01, 85AJ01) for references]. Observed states are displayed in Table 11.21. In addition the excitation of states at $E_{\mathrm{x}}=11.2,12.4,15.3,23$, and (28) MeV has also been suggested: see (80AJ01).

At $E\left({ }^{3} \mathrm{He}\right)=35.6 \mathrm{MeV}$ one finds good fits by DWBA for strong $l=1$ transitions, and reasonable agreement in the forward direction, as well as with $S_{\text {theor. }}$, for weak $l=1$ transitions. Transitions involving $l=0$ or 2 (and 3 ) are weak and the agreement with theory
is poor. It is suggested that ${ }^{11} \mathrm{C}^{*}(8.10)\left[\frac{3}{2}^{-}\right]$is predominantly a $\mathrm{p}_{3 / 2}$ hole state coupled to ${ }^{12} \mathrm{C}^{*}(7.65)\left[0^{+}\right]$: see (80AJ01).

Alpha- $\gamma$ correlations have been studied for $E\left({ }^{3} \mathrm{He}\right)=4.7$ to 12 MeV . Their results are summarized in Table 11.17 and are discussed in detail in reaction 22 of (68AJ02). A measurement of the linear polarization of the $2.00 \mathrm{MeV} \gamma$-ray (together with knowledge of the $\tau_{\mathrm{m}}$ ) fixes $J^{\pi}=\frac{1}{2}^{-}$for ${ }^{11} \mathrm{C}^{*}(2.00) . \tau_{\mathrm{m}}=10.3 \pm 0.7$ fs for ${ }^{11} \mathrm{C}^{*}(2.00)$. See also ${ }^{12} \mathrm{~N}$, and ${ }^{15} \mathrm{O}$ in (86AJ01).

Reaction (b) has been studied at $E\left({ }^{3} \mathrm{He}\right)=75 \mathrm{MeV}$ : transitions to ${ }^{11} \mathrm{C}^{*}(0,2.0,4.3,4.8$, 6.3 ) are observed by looking at $\mathrm{p}, \mathrm{t}$ angular correlations: see (85AJ01). See also (84BE1A; applied).

$$
\text { 24. }{ }^{12} \mathrm{C}\left({ }^{6} \mathrm{Li},{ }^{7} \mathrm{Li}\right)^{11} \mathrm{C} \quad Q_{\mathrm{m}}=-11.471
$$

The angular distributions involving ${ }^{7} \mathrm{Li}_{\mathrm{g} . \mathrm{s} .}+{ }^{11} \mathrm{C}_{\text {g.s. }}$ and ${ }^{7} \mathrm{Li}^{*}{ }_{0.48}+{ }^{11} \mathrm{C}^{*}{ }_{2.00}$ have been studied at $E\left({ }^{6} \mathrm{Li}\right)=36 \mathrm{MeV}$ : see (80AJ01). See also (86GL1E).
25. ${ }^{12} \mathrm{C}\left({ }^{10} \mathrm{~B},{ }^{11} \mathrm{~B}\right){ }^{11} \mathrm{C} \quad Q_{\mathrm{m}}=-7.2673$

At $E\left({ }^{10} \mathrm{~B}\right)=100 \mathrm{MeV}$, angular distributions have been measured involving ${ }^{11} \mathrm{~B}_{\text {g.s. }}+{ }^{11} \mathrm{C}_{\text {g.s. }}$, ${ }^{11} \mathrm{~B}_{\text {g.s. }}+{ }^{11} \mathrm{C}_{2.00}$ and ${ }^{11} \mathrm{C}_{\text {g.s. }} .{ }^{11} \mathrm{~B}_{2.12}$. Both ${ }^{12} \mathrm{C}\left({ }^{10} \mathrm{~B},{ }^{11} \mathrm{~B}\right.$ ) ${ }^{11} \mathrm{C}$ (with ${ }^{11} \mathrm{~B}$ detected in the forward direction) and ${ }^{12} \mathrm{C}\left({ }^{10} \mathrm{~B},{ }^{11} \mathrm{C}\right){ }^{11} \mathrm{~B}$ (with ${ }^{11} \mathrm{C}$ detected in the forward direction) were measured. In each case, ${ }^{11} \mathrm{~B}_{\text {g.s. }}+{ }^{11} \mathrm{C}_{2.00}$ and ${ }^{11} \mathrm{C}_{\text {g.s. }}+{ }^{11} \mathrm{~B}_{2.12}$ were not resolved, but the authors argues that the $\left({ }^{10} \mathrm{~B},{ }^{11} \mathrm{~B}\right)$ case would have little contribution from ${ }^{11} \mathrm{C}_{\text {g.s. }}+{ }^{11} \mathrm{~B}_{2.12}$ (because of the spins of ${ }^{10} \mathrm{~B}$ and ${ }^{11} \mathrm{~B}_{2.12}$ ), so that it essentially gives the ${ }^{11} \mathrm{~B}_{\text {g.s. }}+{ }^{11} \mathrm{C}_{2.00}$ angular distribution, and similarly for the other case. See (85AJ01) and (87OS1E; theor.)
26. ${ }^{12} \mathrm{C}\left({ }^{12} \mathrm{C},{ }^{13} \mathrm{C}\right){ }^{11} \mathrm{C} \quad Q_{\mathrm{m}}=-13.7751$

Angular distributions involving ${ }^{11} \mathrm{C}_{\text {g.s. }}$ have been studied at $E\left({ }^{12} \mathrm{C}\right)=93.8$ and 114 MeV [see (80AJ01, 85AJ01)], at $20 \mathrm{MeV} / A(85 \mathrm{BO} 39)$ and at 25,35 , and $50 \mathrm{MeV} / A$ (88WI09, 89WI07). The strongest peak observed is due to the unresolved ${ }^{13} \mathrm{C}^{*}(3.68+3.85)+{ }^{11} \mathrm{C}^{*}(4.32)$ (88WI09, 89WI07). The results are in agreement with the predictions of the exact FRDWBA. Above $\approx 30 \mathrm{MeV} / A$ the angle-integrated cross sections fall off with an approximately exponential shape (88WI09).
27. ${ }^{13} \mathrm{C}\left(\pi^{+}, \mathrm{d}\right){ }^{11} \mathrm{C}$

$$
Q_{\mathrm{m}}=118.908
$$

At $E_{\pi^{+}}=32 \mathrm{MeV}$ angular distributions have been obtained for the deuterons to ${ }^{11} \mathrm{C}^{*}(0$, 6.48): see (85AJ01).
28. ${ }^{13} \mathrm{C}(\mathrm{p}, \mathrm{t}){ }^{11} \mathrm{C} \quad Q_{\mathrm{m}}=-15.1863$

At $E_{\mathrm{p}}=43.7$ to 50.5 MeV angular distributions of the tritons have been studied to ${ }^{11} \mathrm{C}^{*}(0,2.00,4.32,4.80,6.48,6.90,7.50)$ and to a $T=\frac{3}{2}$ state at $E_{\mathrm{x}}=12.47 \mathrm{MeV}$ [see Table 11.18, whose $J^{\pi}$ is determined to be $\frac{1^{-}}{2}$ [it is thus the analog of ${ }^{11} \mathrm{Be}^{*}(0.32)$ ]. The state decays primarily by $\mathrm{p} \rightarrow{ }^{10} \mathrm{~B}^{*}(1.74)$. Alpha decay to ${ }^{7} \mathrm{Be}^{*}{ }_{\mathrm{g} . \mathrm{s} .+0.4}$ is also observed. Angular distributions have also been measured for $E_{\mathrm{p}}=26.9$ to 65 MeV [see (80AJ01, 85AJ01)]. At $E_{\mathrm{p}}=46.7 \mathrm{MeV}$ the $T=\frac{3}{2}$ state is also observed by (74BE20) who, in addition, report the population of states with $E_{\mathrm{x}}=11.03 \pm 0.03,13.33 \pm 0.06,13.90 \pm 0.04$ and $14.07 \pm 0.04 \mathrm{MeV}$ $[\Gamma=300 \pm 60,270 \pm 80,150 \pm 50$ and $135 \pm 50 \mathrm{keV}$, respectively]. See also (89AR1G).
29. (a) ${ }^{14} \mathrm{~N}(p, \alpha){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-2.9228$
(b) ${ }^{14} \mathrm{~N}(\mathrm{p}, \mathrm{pt}){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-22.737$

Angular distributions have been reported at a number of energies in the range $E_{\mathrm{p}}=5.0$ to 44.3 MeV for the $\alpha_{0}$ and $\alpha_{1}$ groups: see (75AJ02, 80AJ01). For reaction (b) see (86VD1C; $E_{\mathrm{p}}=50 \mathrm{MeV}$; prelim.). See also (84RE1A, 85HA1J), (86MA1P, 87HI1B; applied), (88CA26; astrophysics) and (86GO28; theor.).
30. ${ }^{14} \mathrm{~N}\left(\alpha,{ }^{7} \mathrm{Li}\right){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-20.269$

See (88SH1E; theor.).
31. ${ }^{14} \mathrm{~N}\left({ }^{10} \mathrm{~B},{ }^{13} \mathrm{C}\right){ }^{11} \mathrm{C} \quad Q_{\mathrm{m}}=1.139$

This reaction has been studied at $E\left({ }^{10} \mathrm{~B}\right)=100 \mathrm{MeV}$; see (80AJ01). See also (87OS1E; theor.).
32. ${ }^{16} \mathrm{O}\left(\alpha,{ }^{9} \mathrm{Be}\right){ }^{11} \mathrm{C}$
$Q_{\mathrm{m}}=-24.3099$

See (87KW01, 87KW03; theor.).

## ${ }^{11} \mathrm{~N}$

(Fig. 4)
The ${ }^{14} \mathrm{~N}\left({ }^{3} \mathrm{He},{ }^{6} \mathrm{He}\right){ }^{11} \mathrm{~N}$ reaction has been studied at $E\left({ }^{3} \mathrm{He}\right)=70 \mathrm{MeV}$. A ${ }^{6} \mathrm{He}$ group is observed which corresponds to a state in ${ }^{11} \mathrm{~N}$ with an atomic mass excess of $25.23 \pm 0.10$ MeV and $\Gamma=740 \pm 100 \mathrm{keV}$. The cross section for forming this state is $0.5 \mu \mathrm{~b} / \mathrm{sr}$ at $10^{\circ}$. The observed state is interpreted as being the $J^{\pi}=\frac{1}{2}^{-}$mirror of ${ }^{11} \mathrm{Be}^{*}(0.32)$ because of its width; the $\frac{1}{2}^{+}$mirror ${ }^{11} \mathrm{Be}$ (g.s.) would be expected to be much broader (74BE20). This ${ }^{11} \mathrm{~N}$ state is unbound with respect to decay into ${ }^{10} \mathrm{C}+\mathrm{p}$ by 2.24 MeV . (88WA18) adopt an atomic mass excess of $24.89 \pm 0.14 \mathrm{MeV}$ for ${ }^{11} \mathrm{~N}_{\text {g.s. }}$. TThis value assumes that the first excited state in ${ }^{11} \mathrm{~N}$ is at $E_{\mathrm{x}}=0.34 \mathrm{MeV}$.] We suggest an uncertainty of $\pm 0.2 \mathrm{MeV}$ because the $E_{\mathrm{x}}$ of the first excited state in ${ }^{11} \mathrm{~N}$ may be depressed relative to ${ }^{11} \mathrm{Be}^{*}$. The ground state is then unstable with respect to ${ }^{10} \mathrm{C}+\mathrm{p}$ by 1.90 MeV . See also (85AN28, 86AN07; theor.).

$$
\begin{aligned}
& { }^{11} 0,{ }^{11} \mathrm{~F},{ }^{11} \mathrm{Ne} \\
& \text { (Not illustrated) }
\end{aligned}
$$

These nuclei have not been observed: see (80AJ01, 85AJ01) and (86AN07, 87SA15; theor.).

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[^0]:    a See discussion in (MI82C). See also Table 11.4 in (80AJ01) and Tables 11.5 and 11.13 here.
    ${ }^{\text {b }}$ See also (65OL03).
    c $\delta=-0.19 \pm 0.03$.
    d $\delta=0.03 \pm 0.05$.
    e $\delta=-0.05 \pm 0.02$.
    ${ }^{\text {f }} \delta=-0.45 \pm 0.18$. This value leads to too large a value of $\Gamma_{\gamma}$ for an M3 transition (P.M. Endt, private communication).
    ${ }^{\mathrm{g}}$ This is probably the ${ }^{11} \mathrm{~B}$ analog of ${ }^{11} \mathrm{C}^{*}(8.10)$. If so $J^{\pi}=\frac{3}{2}^{-}$.
    h $\delta=-0.11 \pm 0.04$.
    ${ }^{\text {i }}$ Weighted mean of branching ratios and $\Gamma_{\gamma}$ (84HA13). Earlier work is also included: see (84HA13).

[^1]:    ${ }^{\text {a }}$ Mostly from (66CU02). For other parameters see Table 11.9 in (75AJ02). See also Table 11.8 in (85AJ01).
    ${ }^{\text {b }}{ }^{7} \mathrm{Li}\left(\alpha, \alpha^{\prime} \gamma\right)^{7} \mathrm{Li}: \sigma$ (total).
    ${ }^{\text {c }}{ }^{7} \mathrm{Li}\left(\alpha, \alpha_{0}\right)^{7} \mathrm{Li}$.
    ${ }^{d}{ }^{7} \operatorname{Li}(\alpha, \mathrm{n}){ }^{10} \mathrm{~B}$ threshold.
    ${ }^{e}$ Anomaly in angular distribution.
    ${ }^{\mathrm{f}}$ Observed at $\theta=60^{\circ}$.

[^2]:    ${ }^{\text {a }}$ See also Table 11.4 here, and Table 11.16 in (80AJJ01). For references see Table 11.14 in (85AJ01).
    ${ }^{\mathrm{b}}$ Mean of values shown in Table 11.14 (85AJ01).

[^3]:    ${ }^{\text {a }}$ See also Table 11.17 here, and Tables 11.23 and 11.24 in (75AJ02). Table 11.23 displays some other reported resonances; Table 11.24 gives detailed parameters for ${ }^{11} \mathrm{C}^{*}(9.73,10.08,10.68,12.65)$. For references see Table 11.22 in (80AJ01). For unpublished work and other references see Table 11.20 in (85AJ01). (88ABZW) [in (p, $\mathrm{p}^{\prime} \gamma$ ) and ( $\mathrm{p}, \alpha \gamma$ ); $E_{\mathrm{p}}=2$ to 5 MeV prelim.] report 5 states with energies $11.84,11.37(?), 12.63,12.75$, and 13.1 MeV .
    b (83WI09).
    ${ }^{c} \Gamma_{\gamma} / \Gamma_{\text {tot }}=(2.6 \pm 0.15) \times 10^{-4}:$ see (83WI09) $. \Gamma_{\gamma} / \Gamma_{\text {tot }}=0.20 \pm 0.05$ and $<0.06$, respectively for ${ }^{11} \mathrm{C}^{*}(8.42,8.66)$, respectively: $\Gamma_{\text {tot }} \leq 5 \mathrm{keV}$ for both states (83WI09).
    ${ }^{\mathrm{d}} \Gamma_{\mathrm{p}} \Gamma_{\gamma} / \Gamma \approx 20 \mathrm{eV}$.
    ${ }^{e}$ Probably part of the E1 giant resonance.

